

Covariant and observer-independent approach to geometric optics in GR

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in collaboration with

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Michele Grasso (CFT PAN, Warsaw)

Julius Serbenta (CFT PAN, Warsaw)

IX International Meeting on Lorentzian Geometry

17th-24th June 2018

Banach Center, Warsaw



Papers

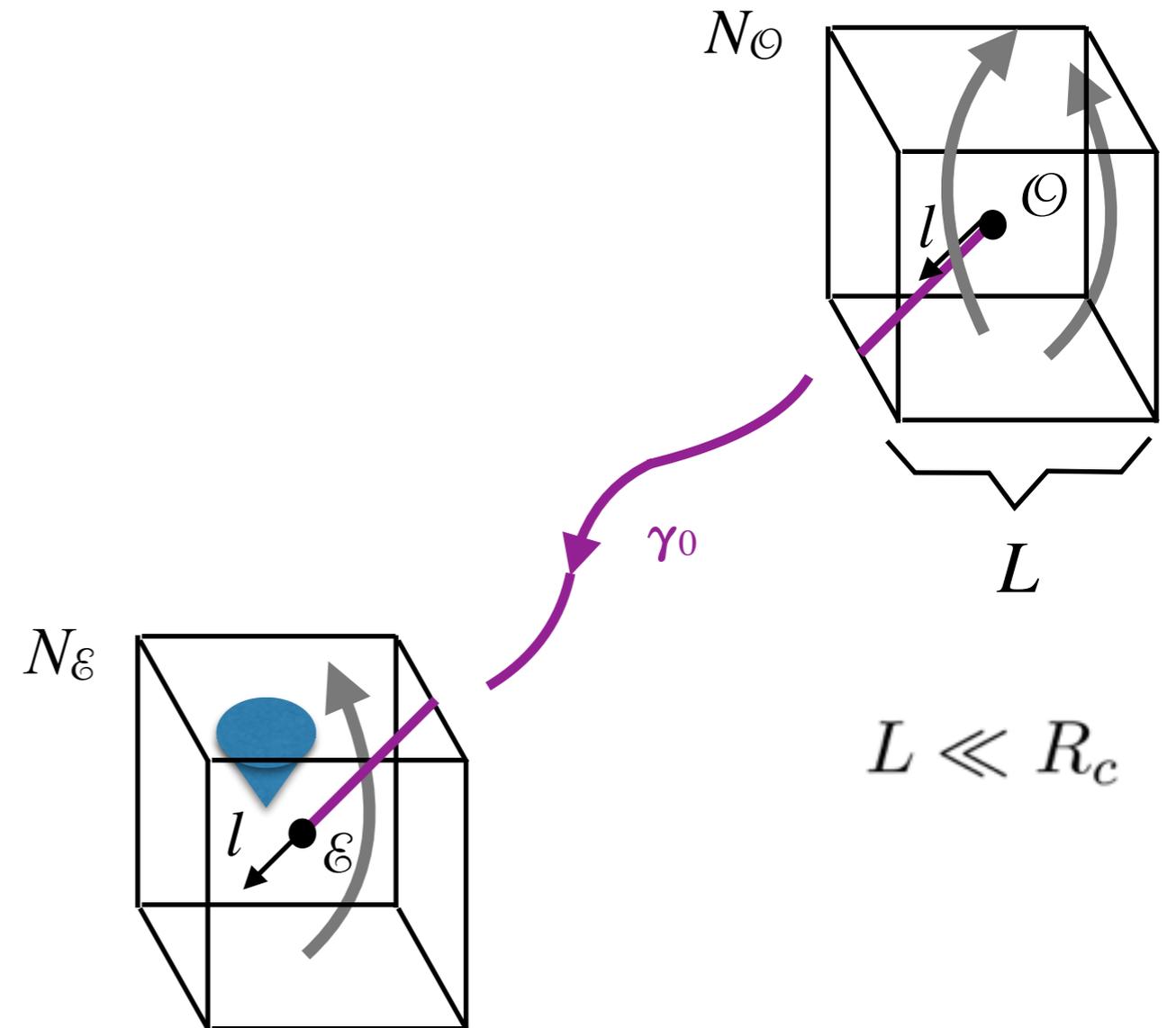
Based on:

- M.K., J. Kopiński “*Optical drift effects in general relativity*”, JCAP 03 (2018) 012, e-print: 1711.00584 [gr-qc]
- M.K., M. Grasso, J. Serbenta “*Geometric optics in general relativity using bi-local operators*”, in preparation

NCN project SONATA BIS No 2016/22/E/ST9/00578

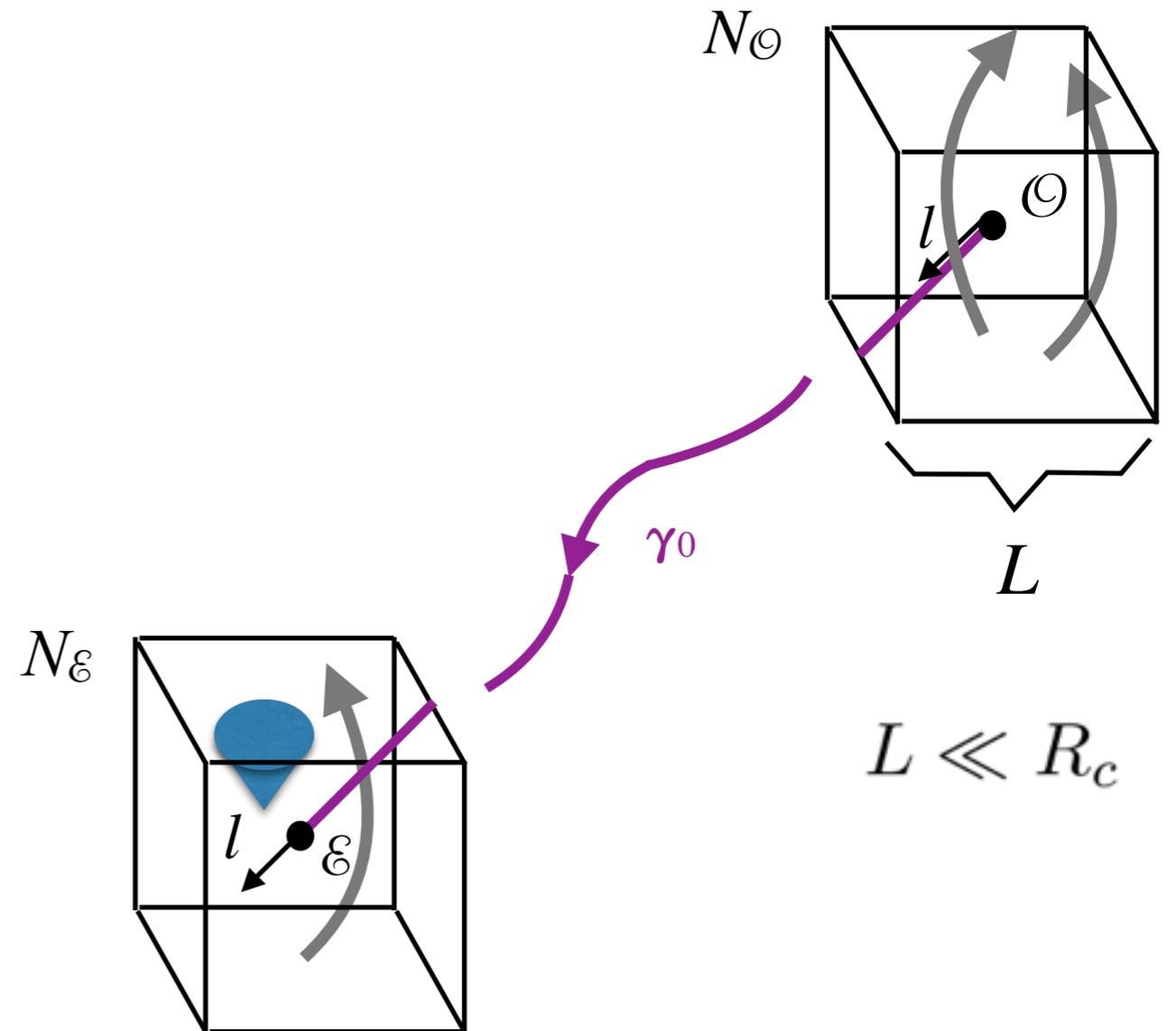
“*Local relativistic perturbative framework in hydrodynamics and general relativity and its application to cosmology*”

Motivation



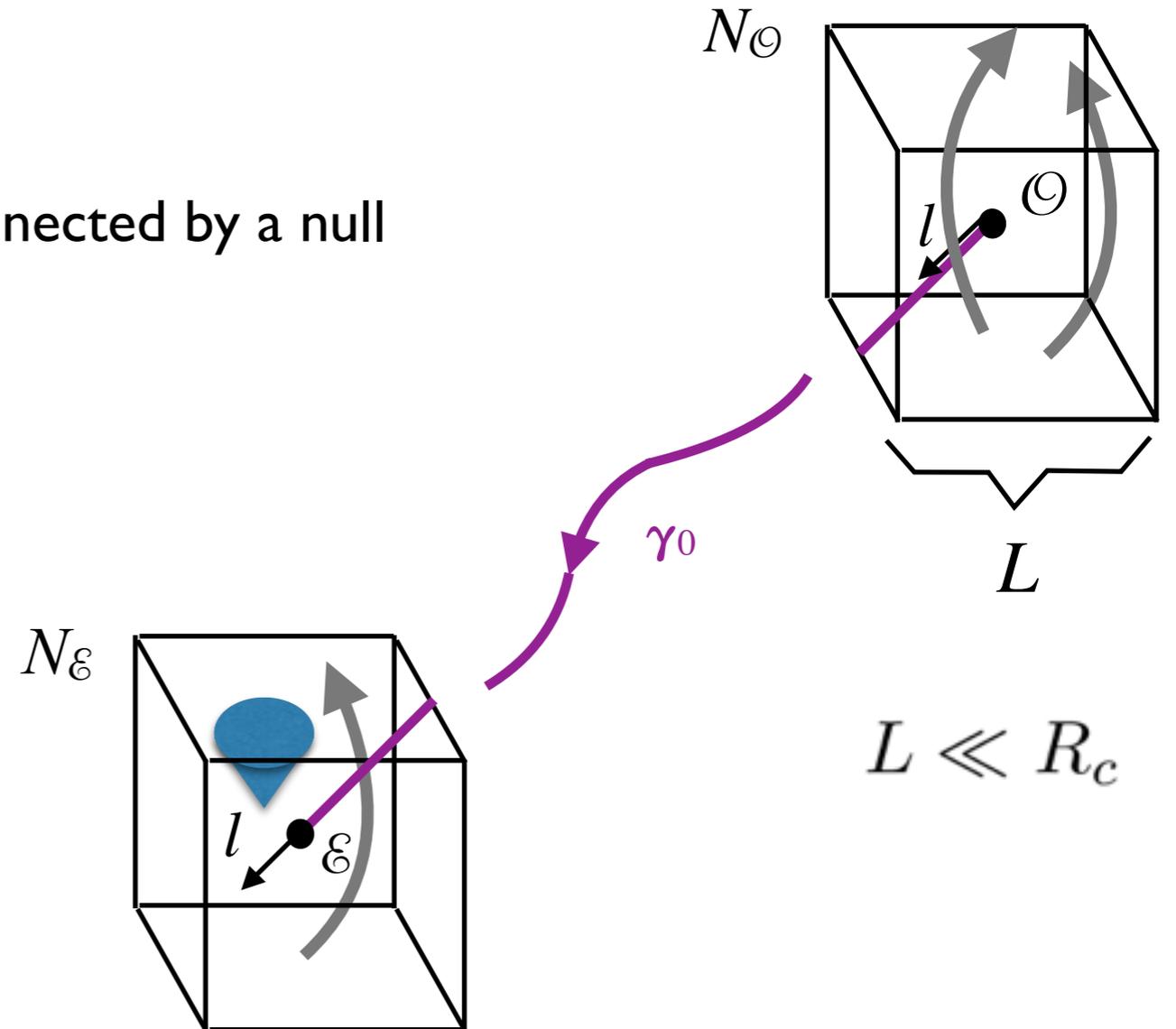
Motivation

- Spacetime with *any* Lorentzian metric



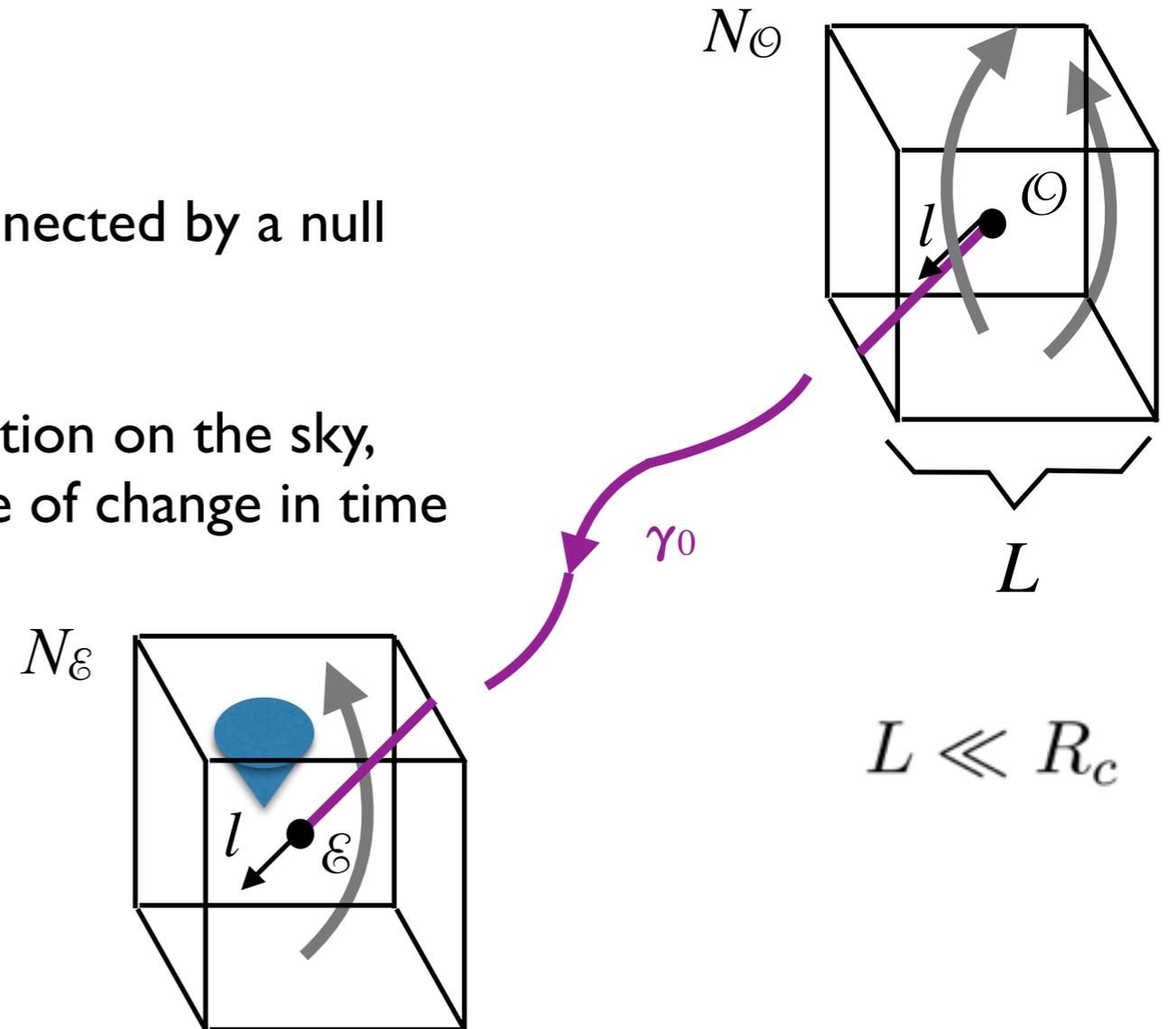
Motivation

- Spacetime with *any* Lorentzian metric
- Two distant, small regions of spacetime, connected by a null geodesics



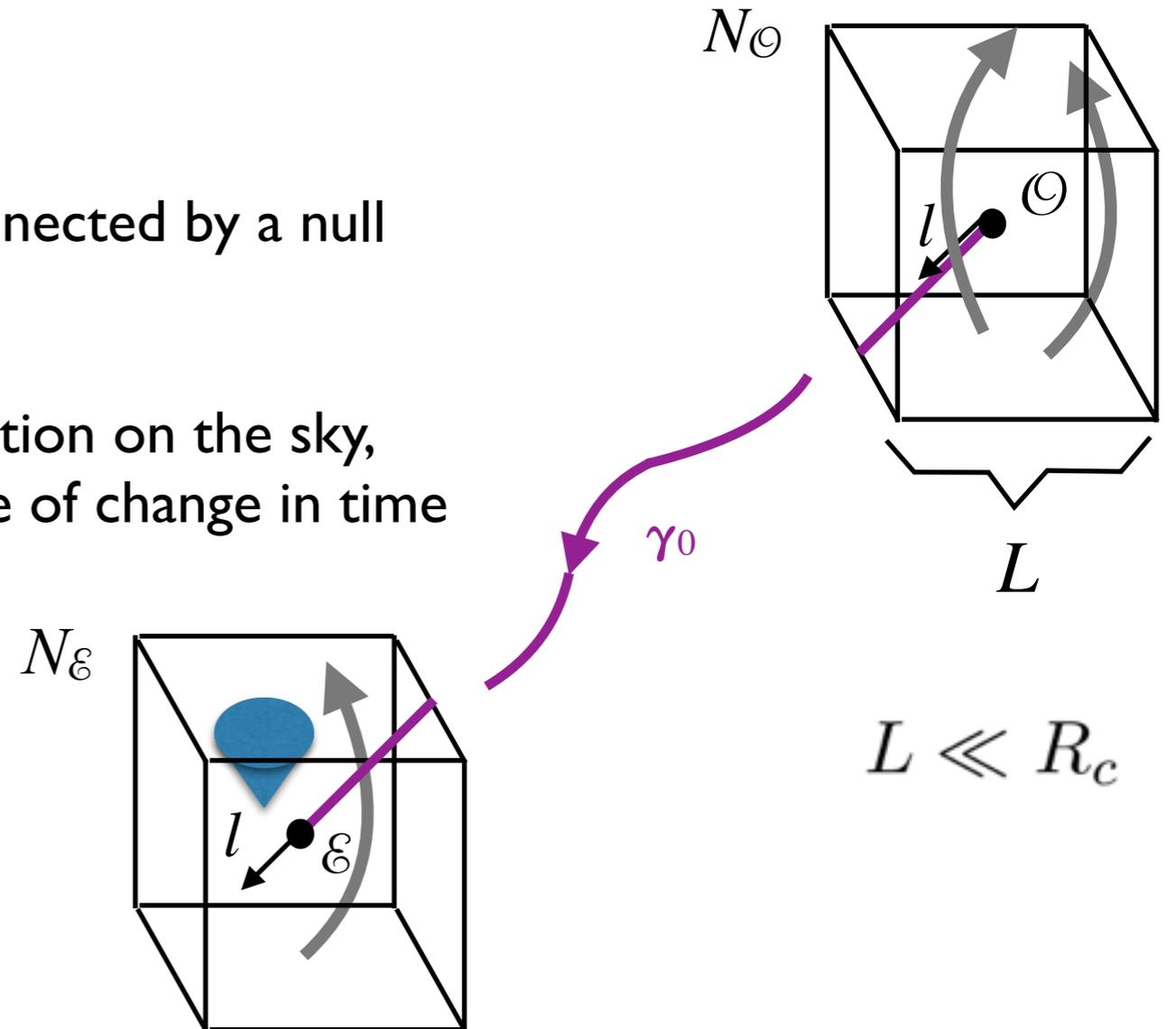
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- Spacetime with *any* Lorentzian metric
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- Observers measure the time of arrival, position on the sky, redshift, image size and distortion, their rate of change in time (drift) ...



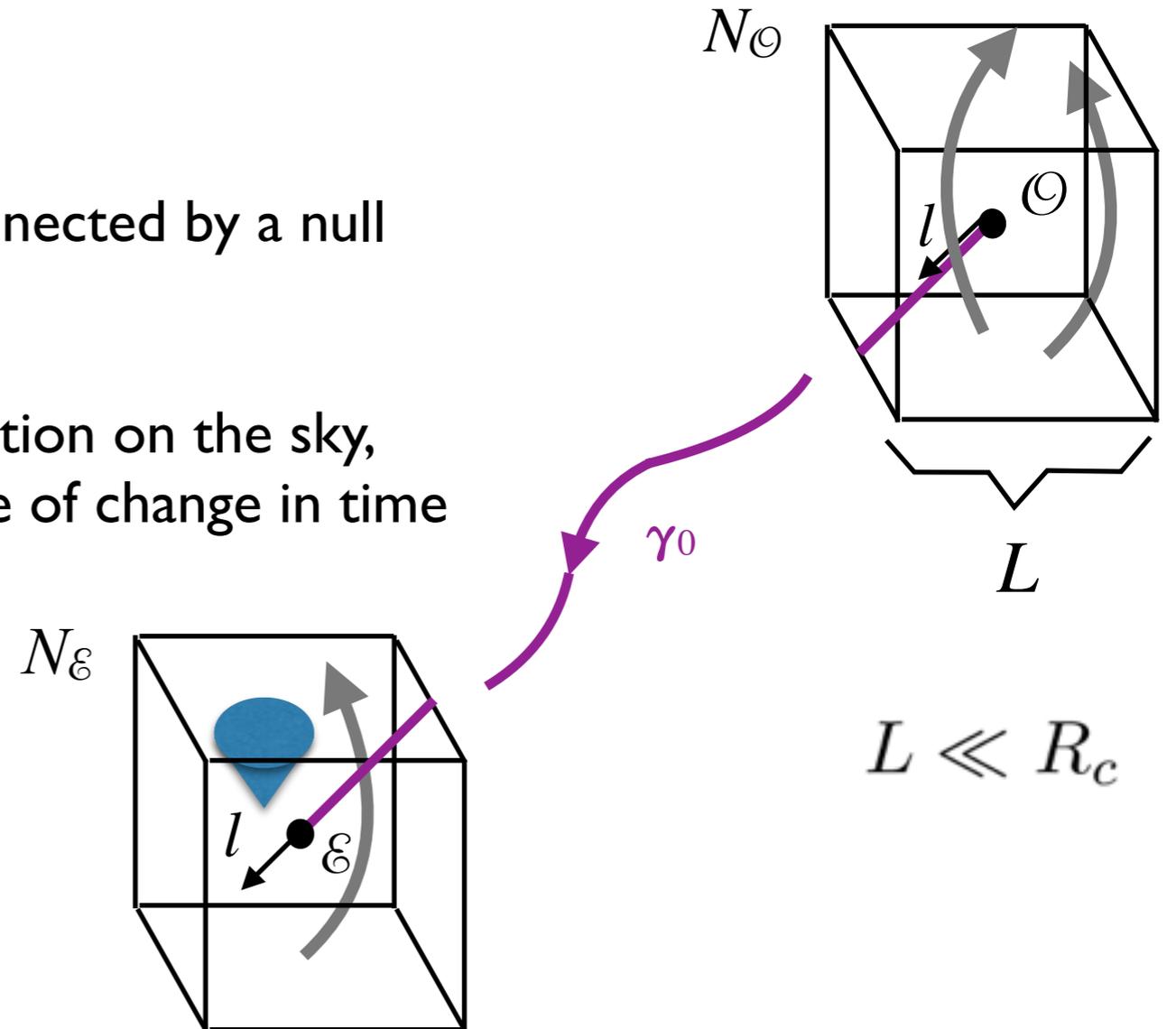
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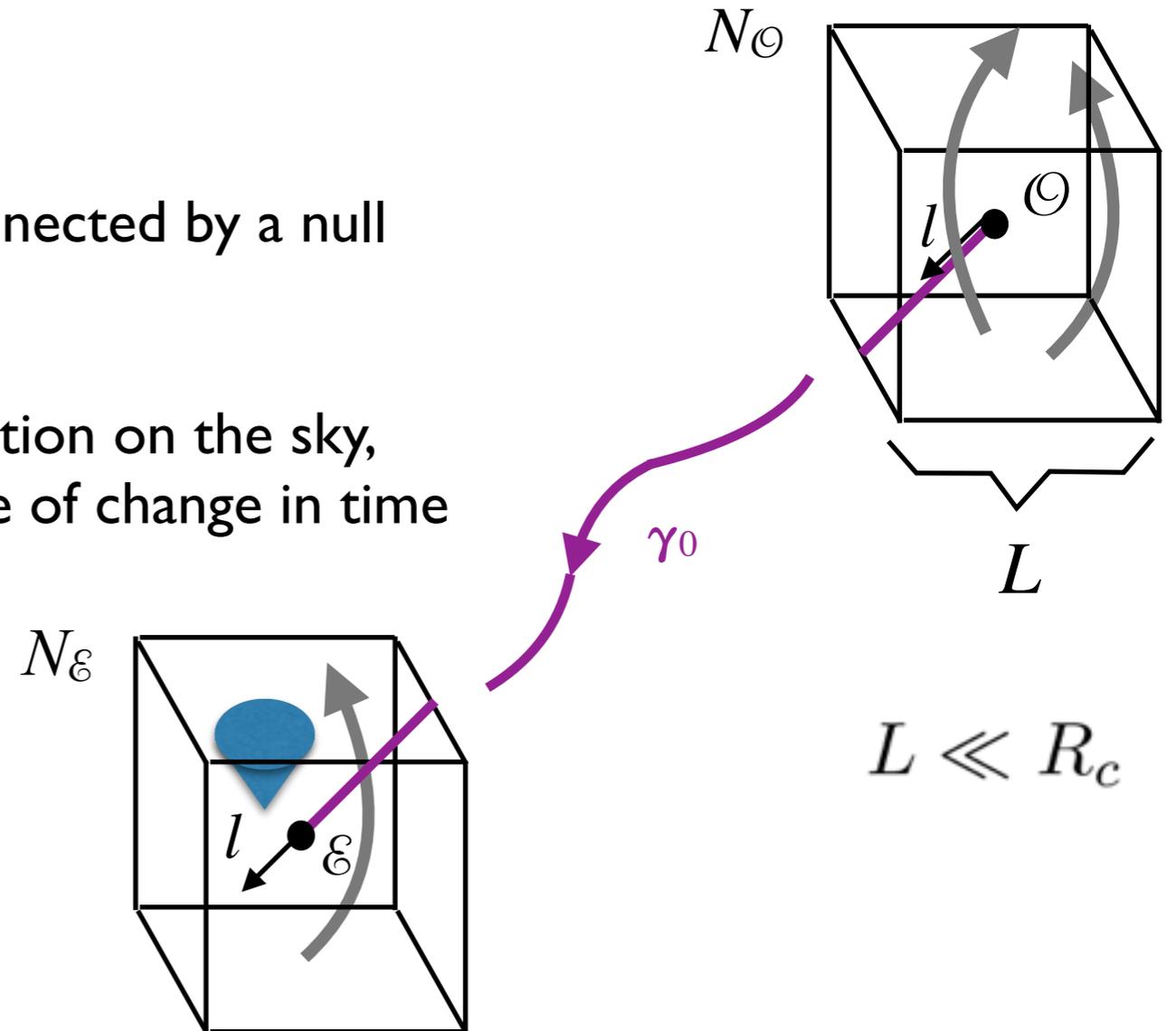
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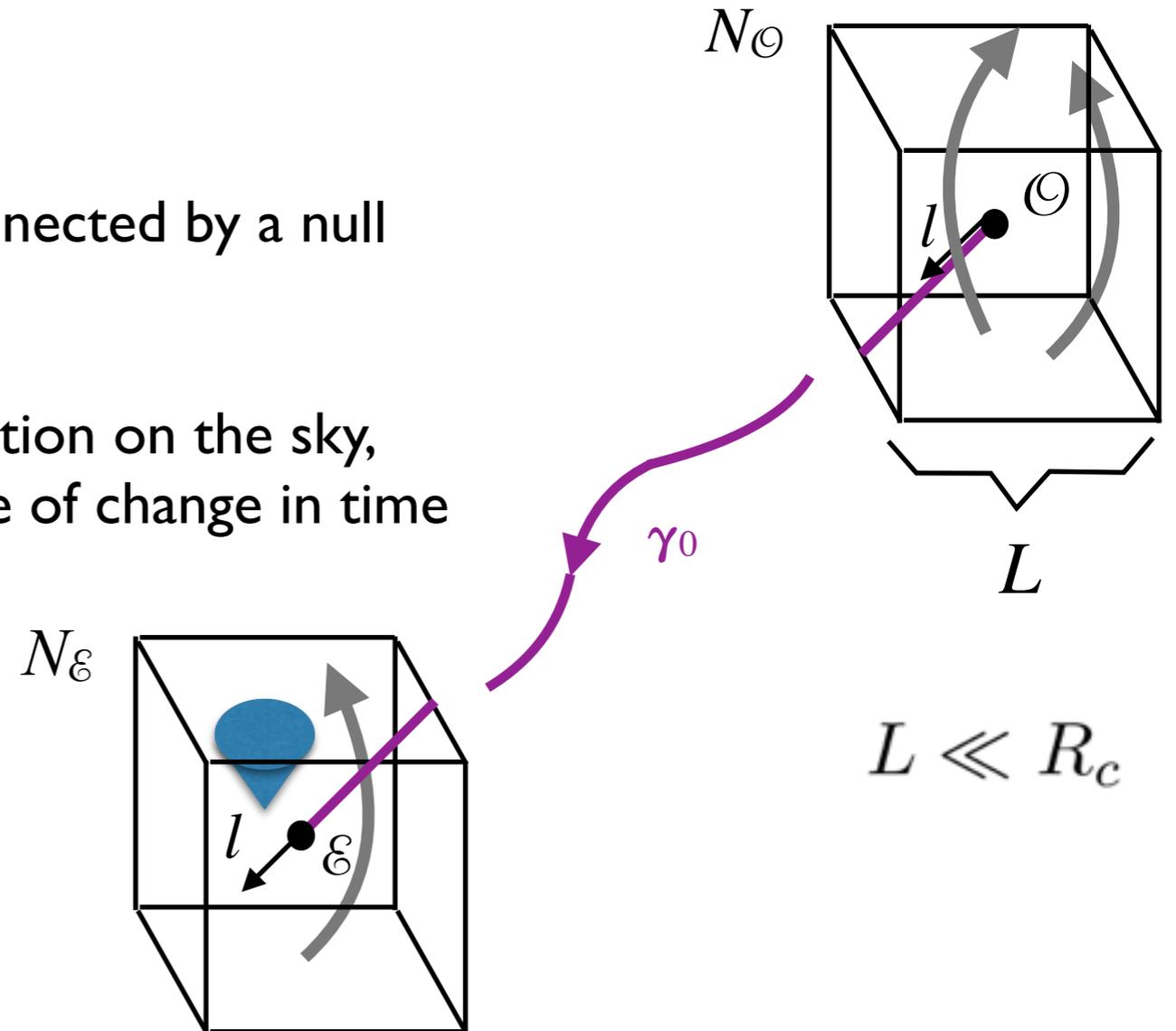
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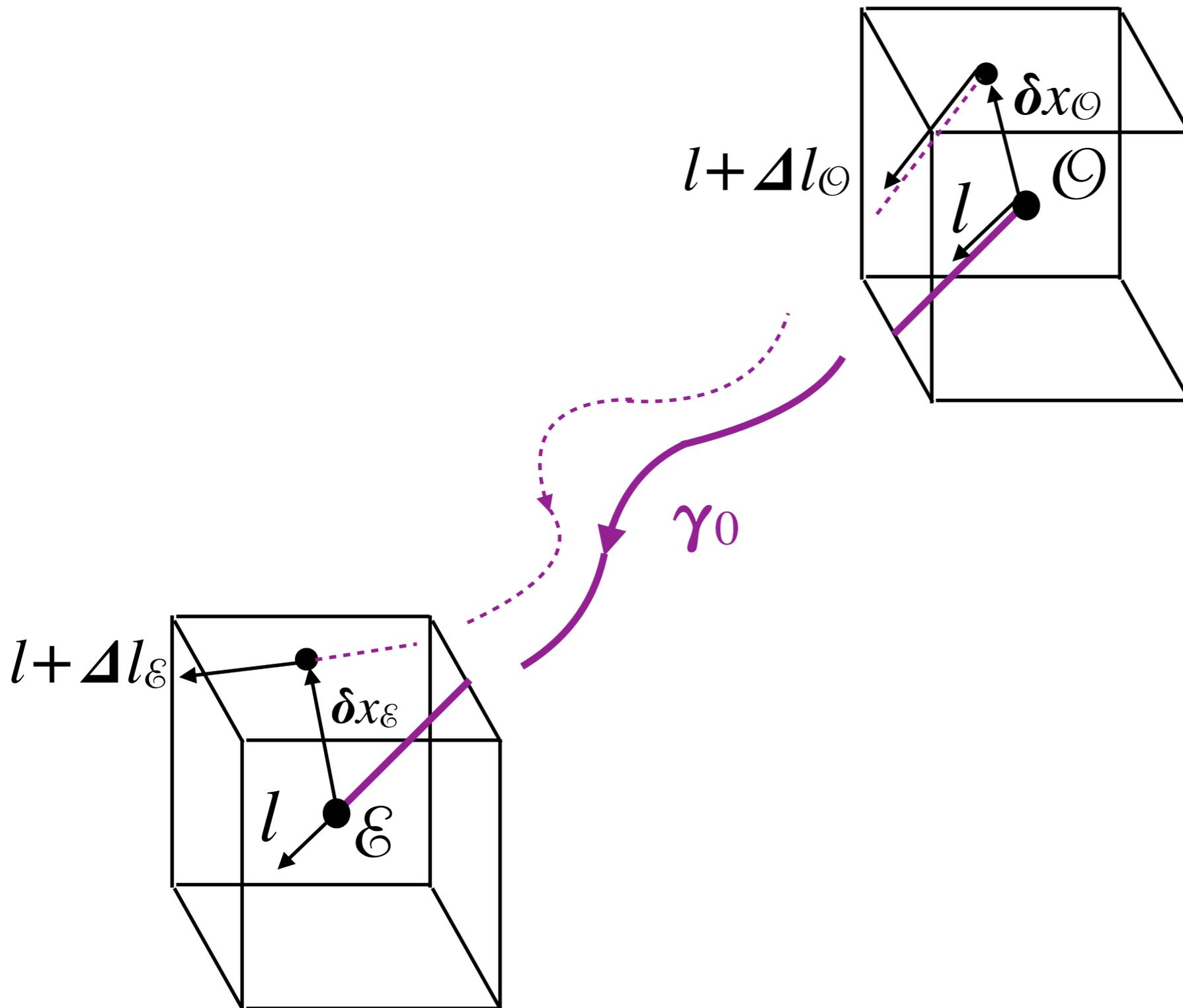


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- Idea: geometric formulation of the problem (coordinate- and frame-independent)



Geometry

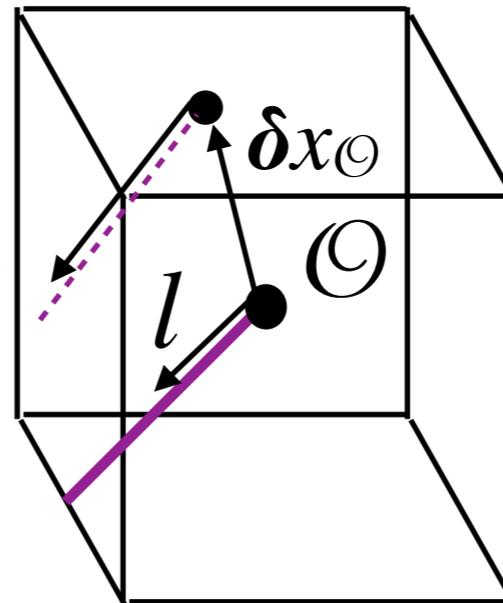


Geometry

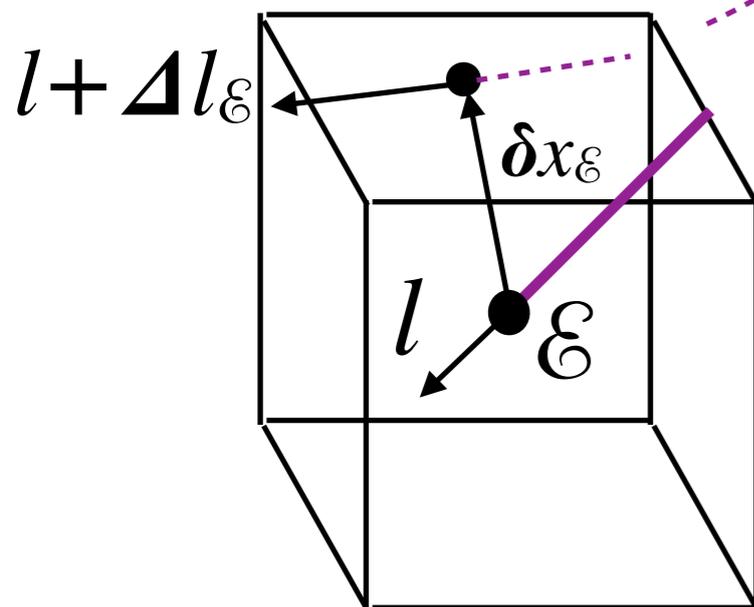
$$\delta x_{\mathcal{O}}^{\mu}$$

$$\Delta l_{\mathcal{O}}^{\mu} = \delta l_{\mathcal{O}}^{\mu} + \Gamma^{\mu}_{\nu\sigma}(\mathcal{O}) l^{\nu} \delta x_{\mathcal{O}}^{\sigma}$$

$l + \Delta l_{\mathcal{O}}$



γ_0

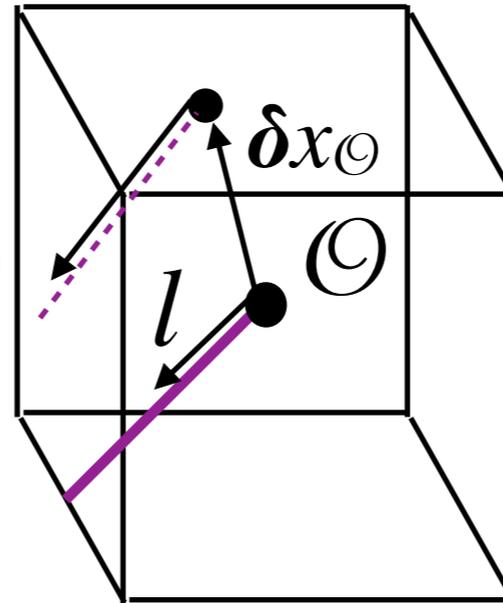


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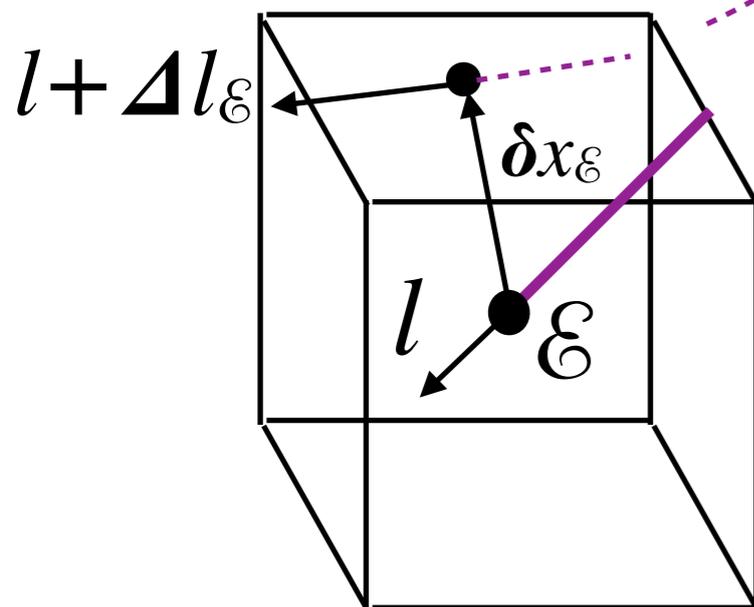


$$\nabla_l \nabla_l \xi^{\mu} - R^{\mu}_{\alpha\beta\nu} l^{\alpha} l^{\beta} \xi^{\nu} = 0$$

$$\xi^{\mu}(\lambda_{\mathcal{O}}) = \delta x_{\mathcal{O}}^{\mu}$$

$$\nabla_l \xi^{\mu}(\lambda_{\mathcal{O}}) = \Delta l_{\mathcal{O}}^{\mu}$$

γ_0



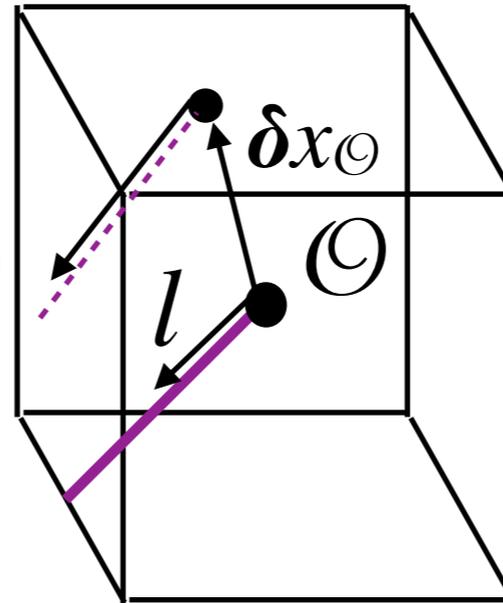
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Geometry

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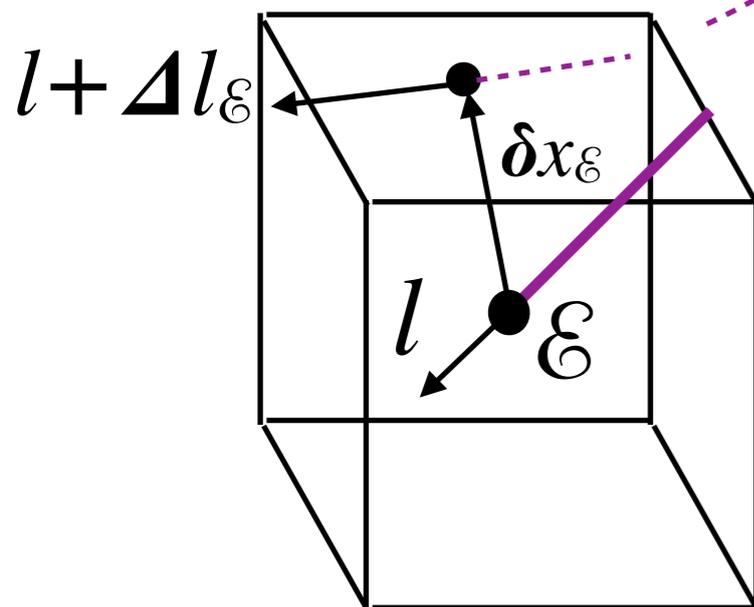
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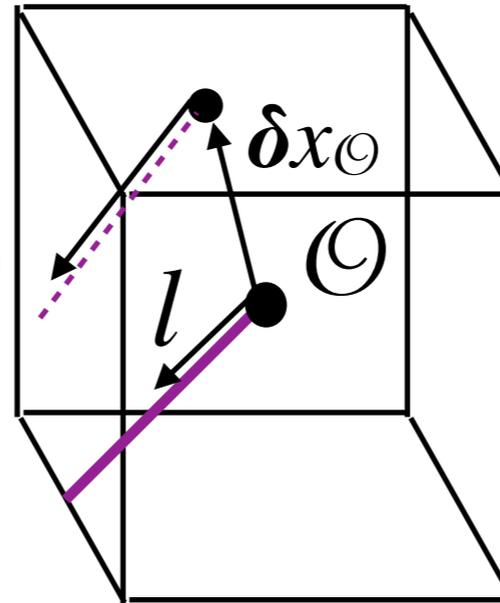


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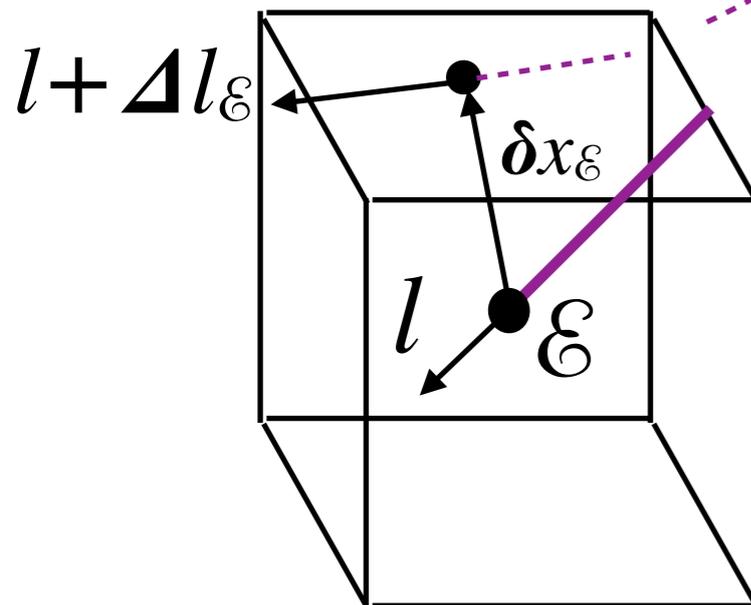
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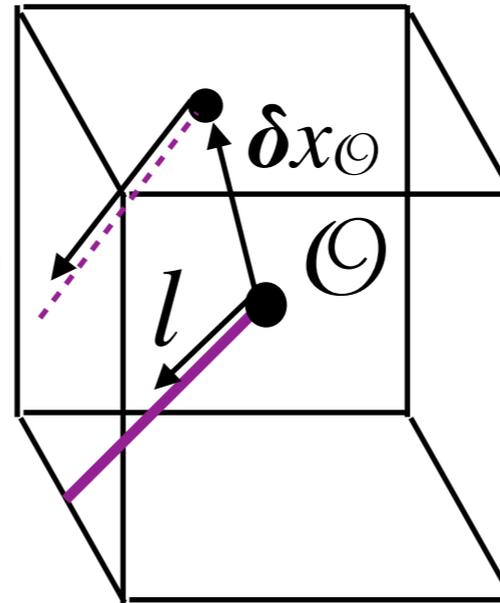
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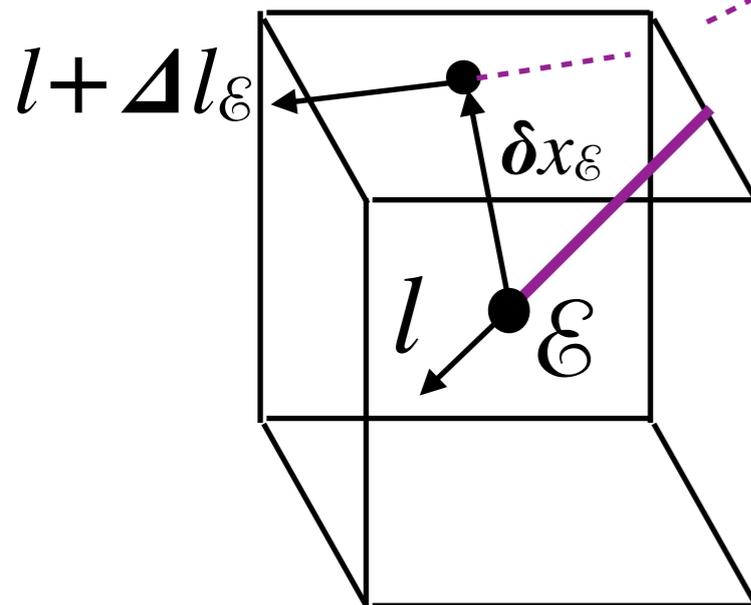
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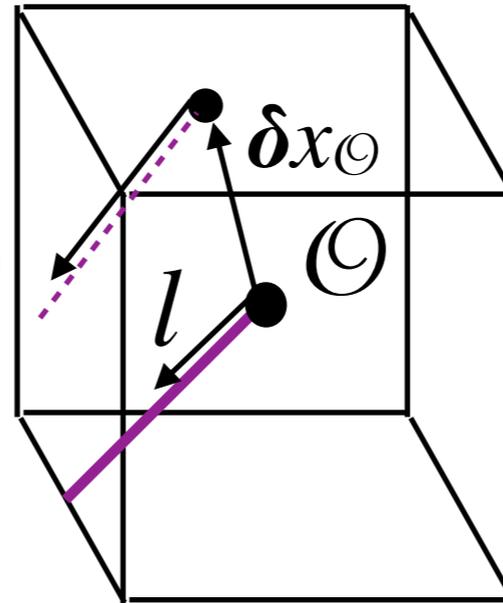
$$W_{**} : T_{\mathcal{O}}M \mapsto T_{\mathcal{E}}M$$

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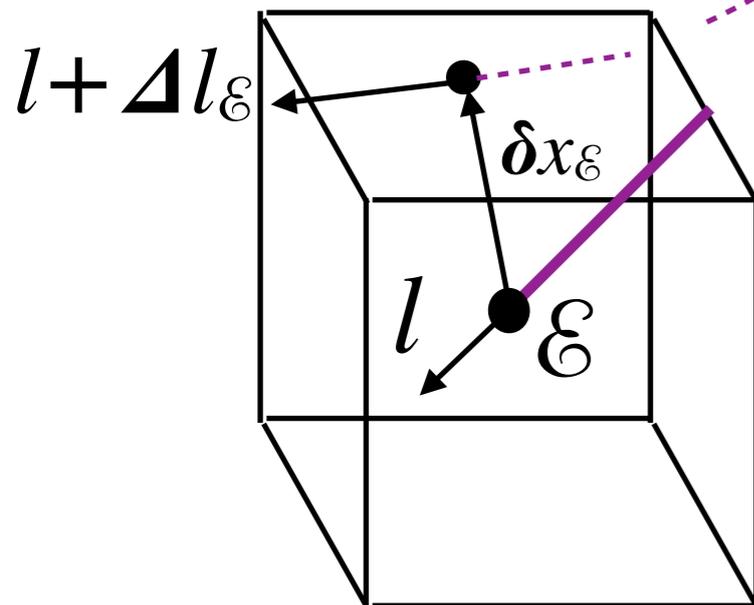
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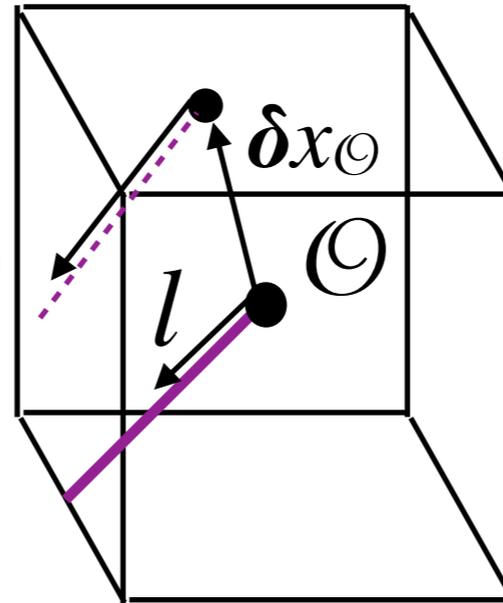
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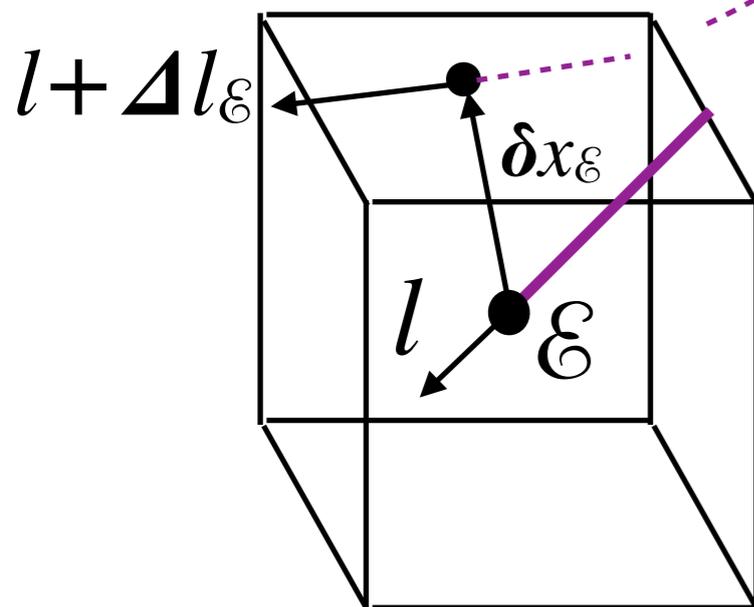
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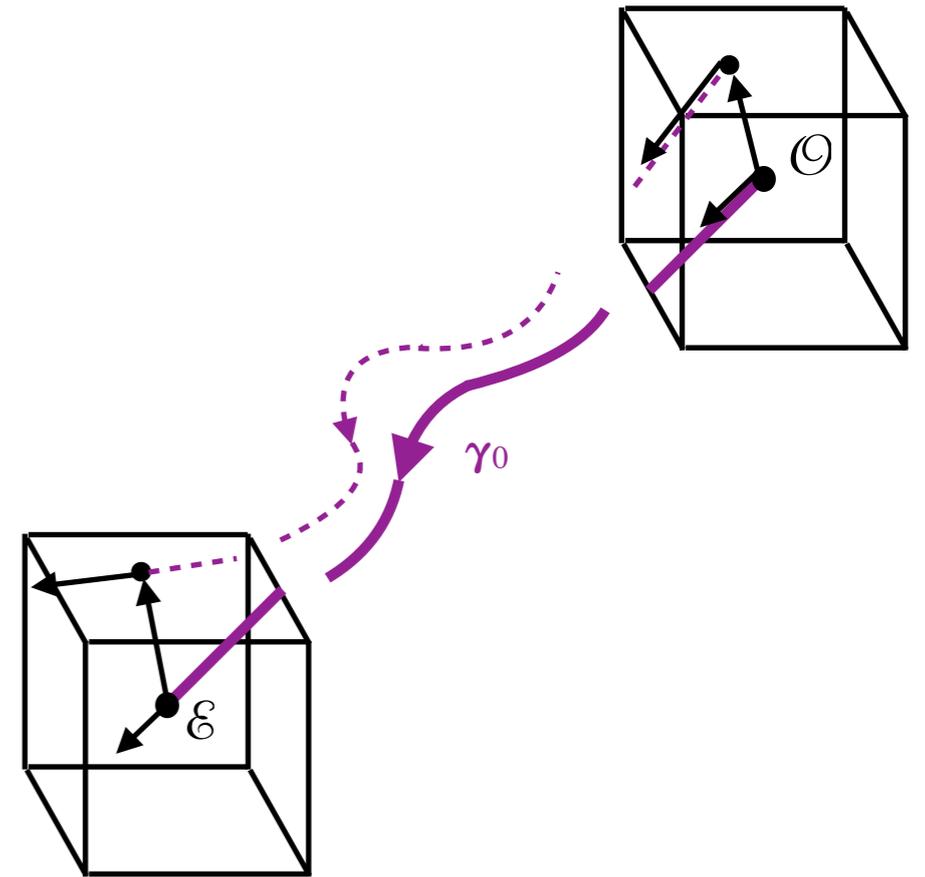


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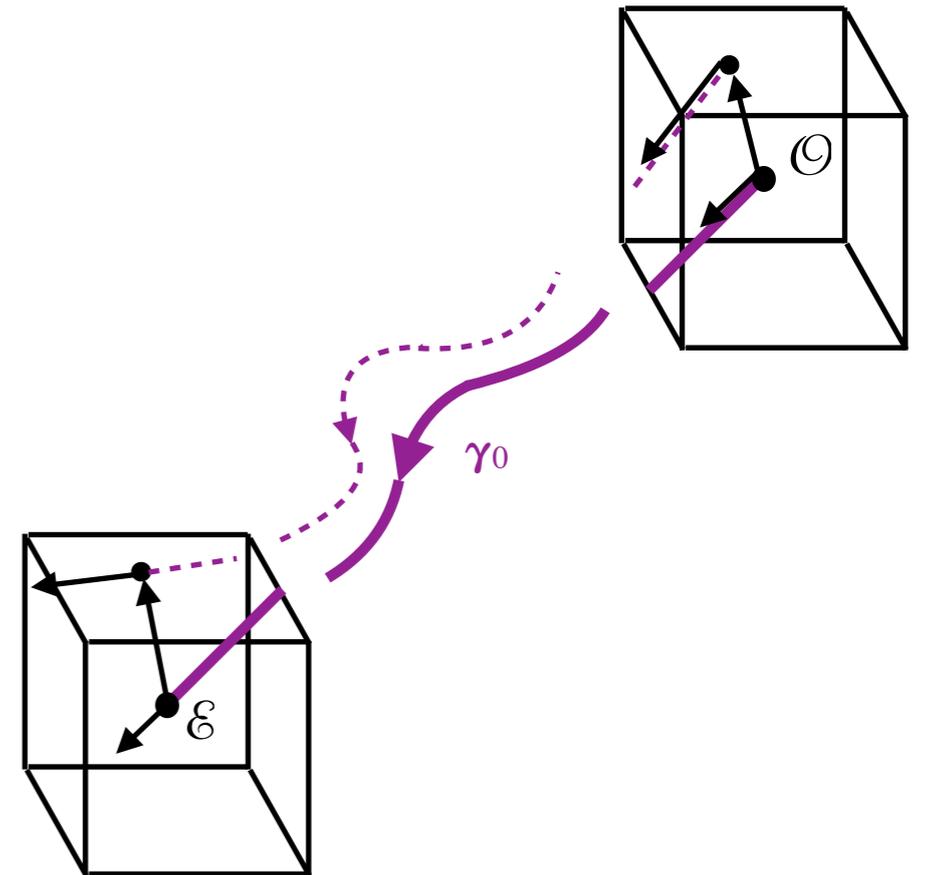
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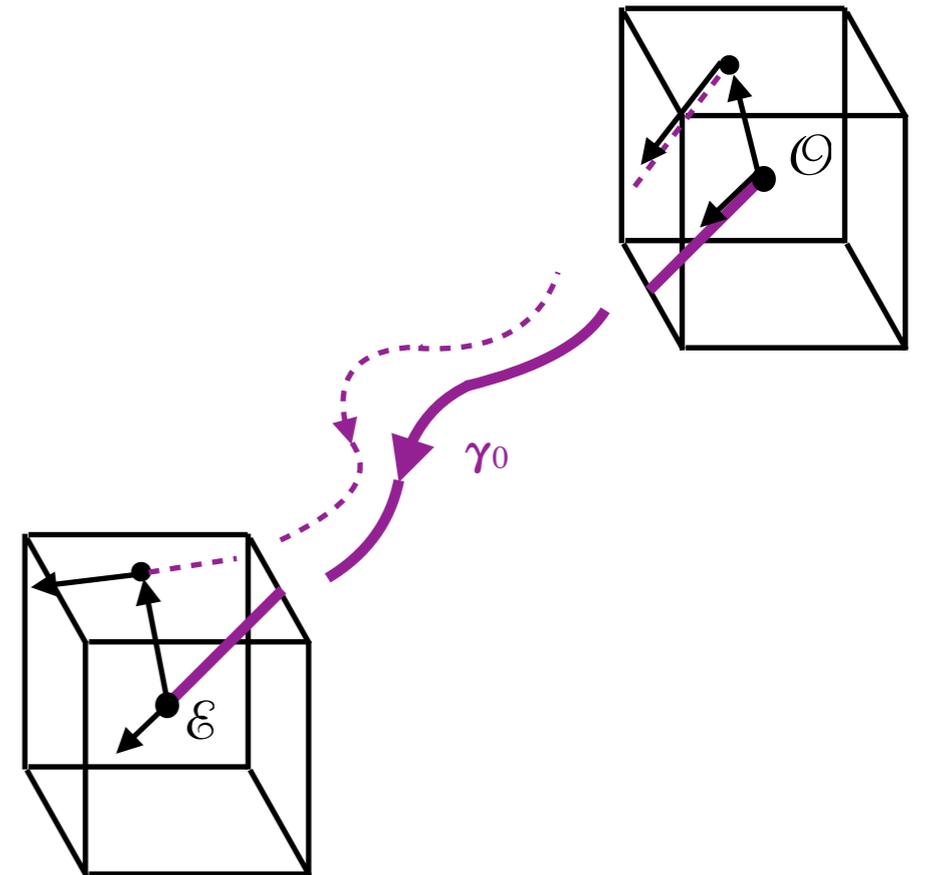
Geometry

- W_{**} contain information about *all* geodesics from N_{Θ} to $N_{\mathcal{E}}$



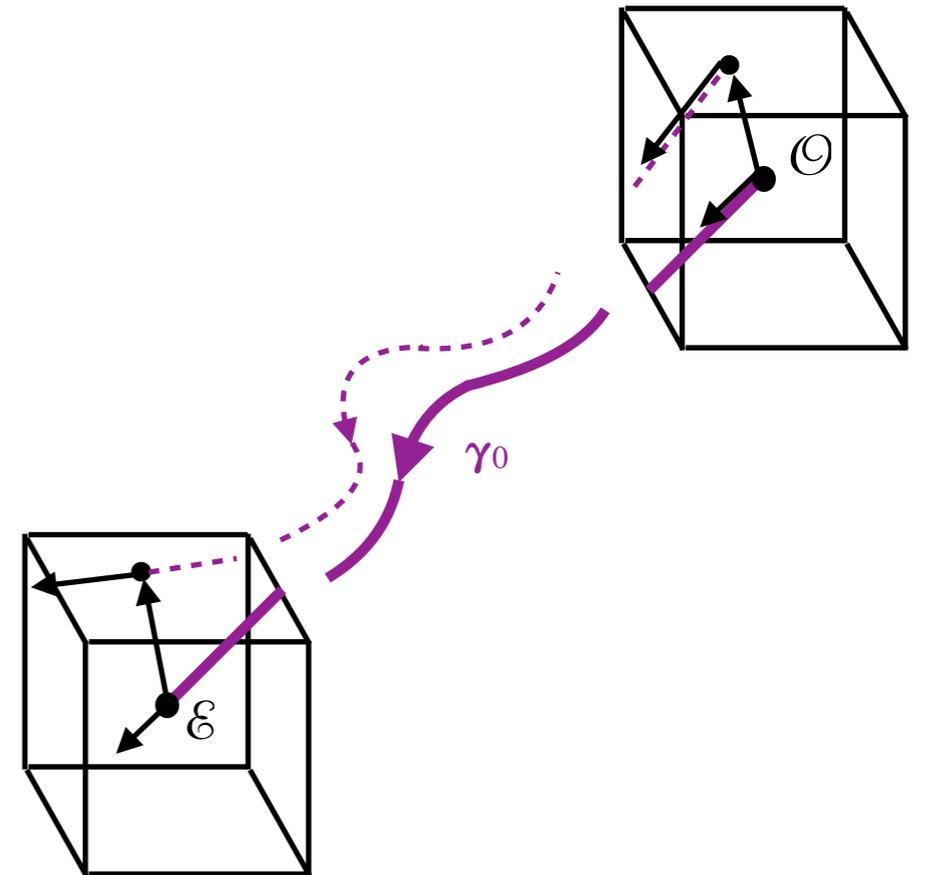
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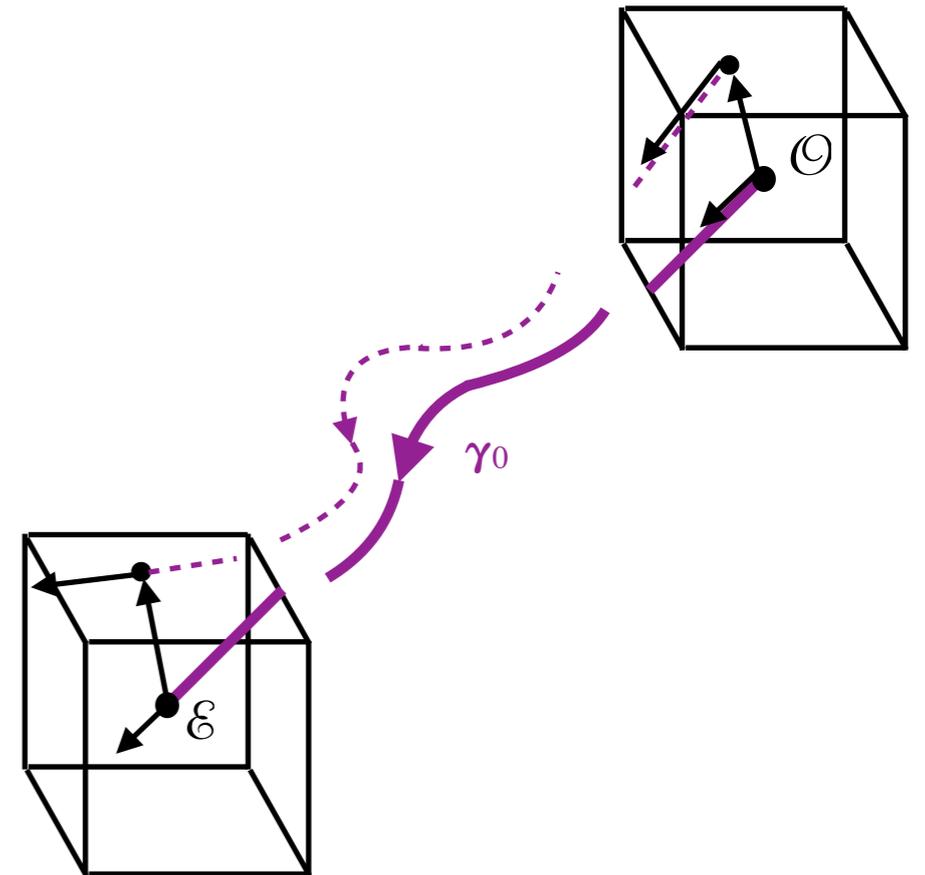
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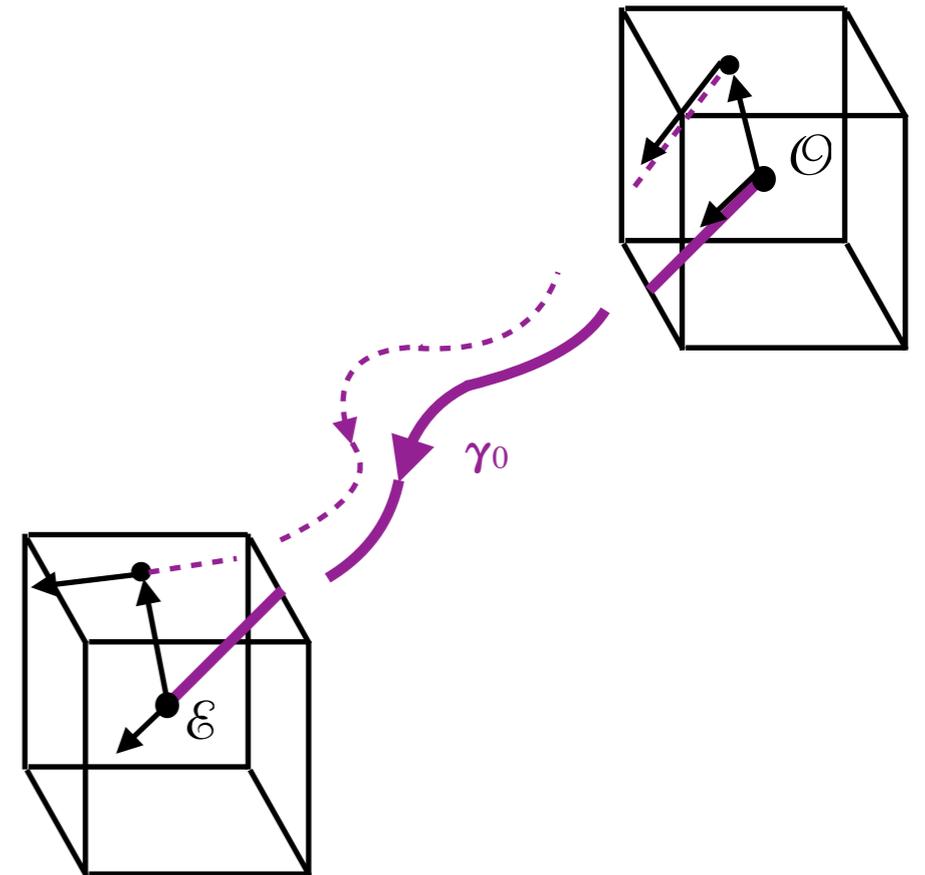


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- geometric optics - only the geodesic's path matters. We can identify null geodesics differing by affine reparametrization



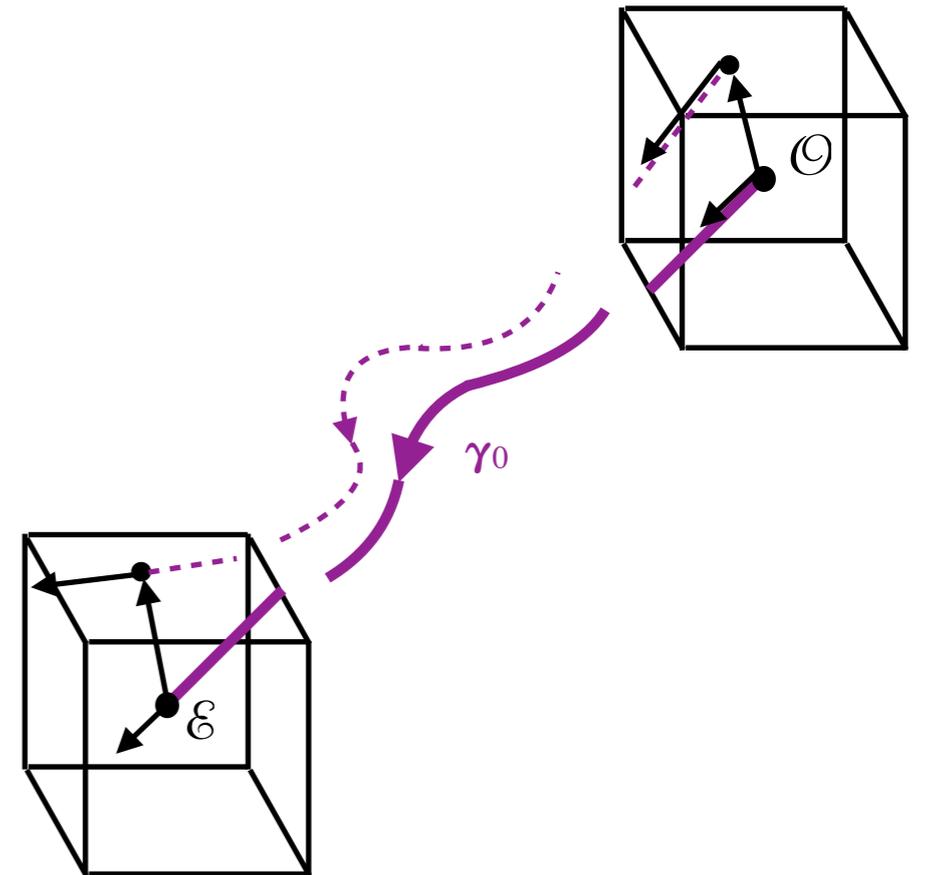
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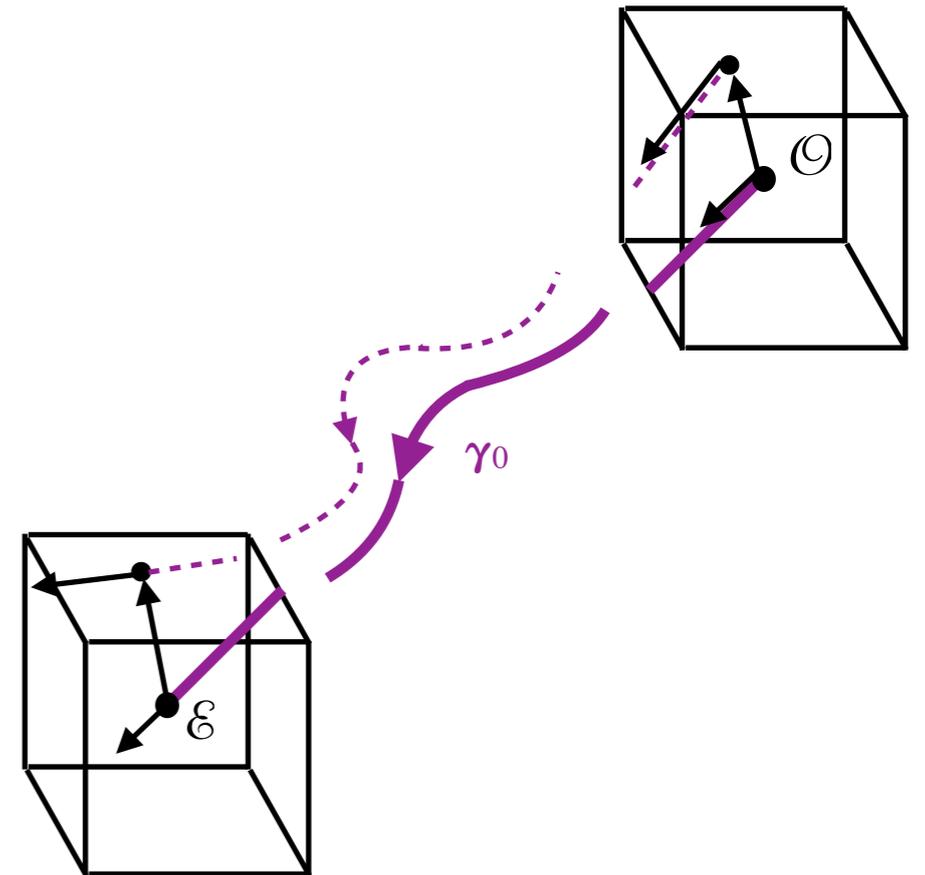
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$$Dim = 8 - 1 - 2 = 5$$

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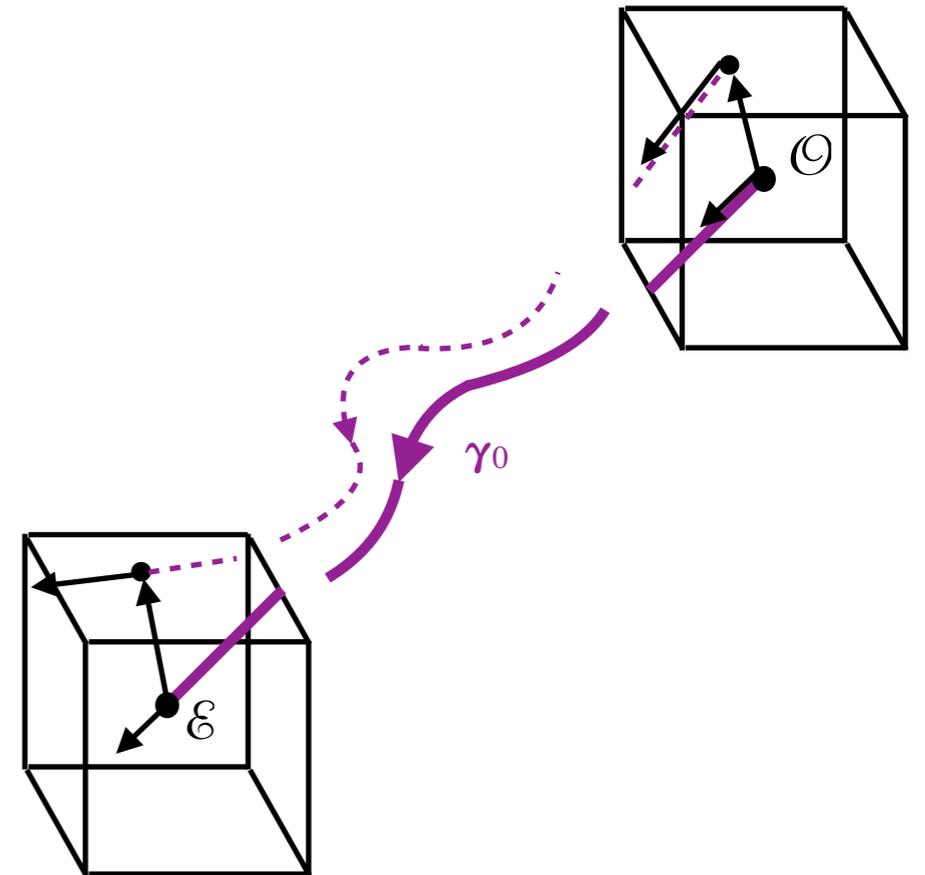
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$$\Delta l_{\mathcal{O}}^\mu \ll l_{\mathcal{O}}^\mu$$



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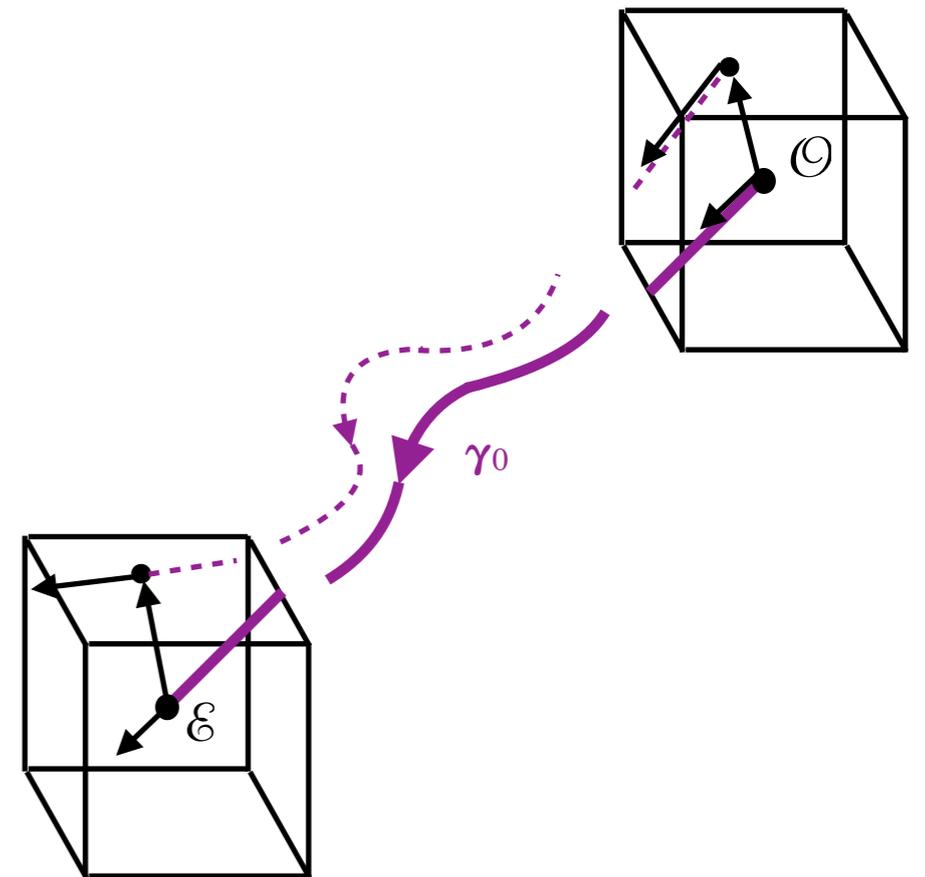
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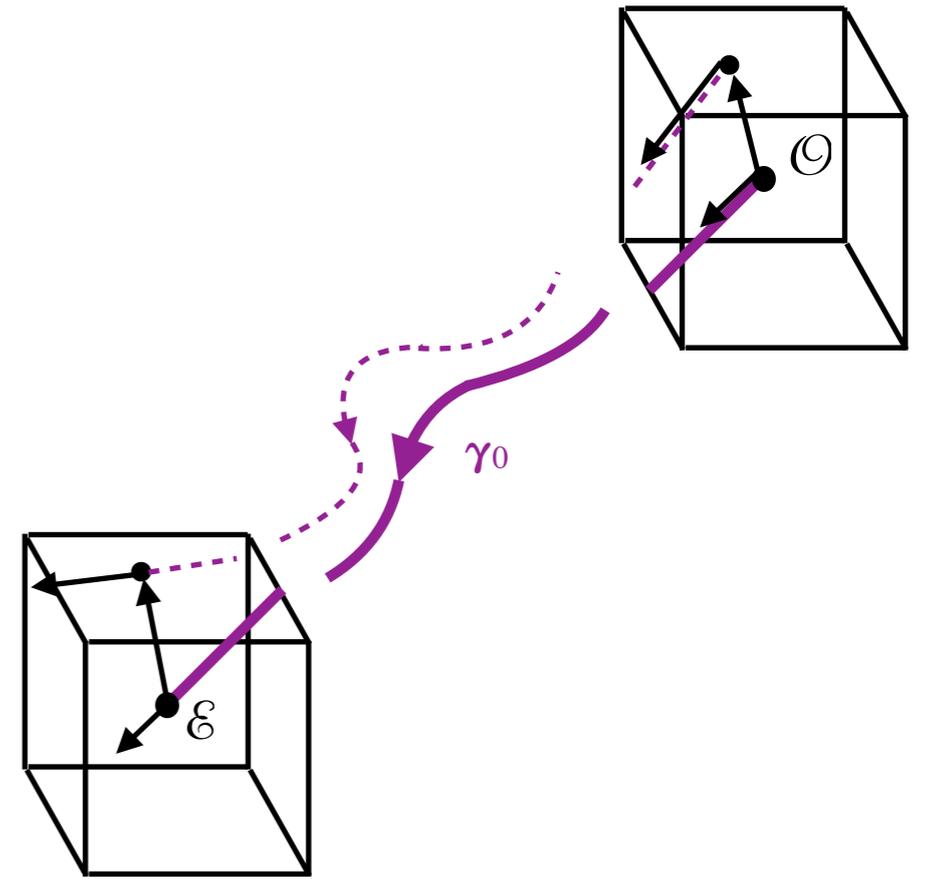
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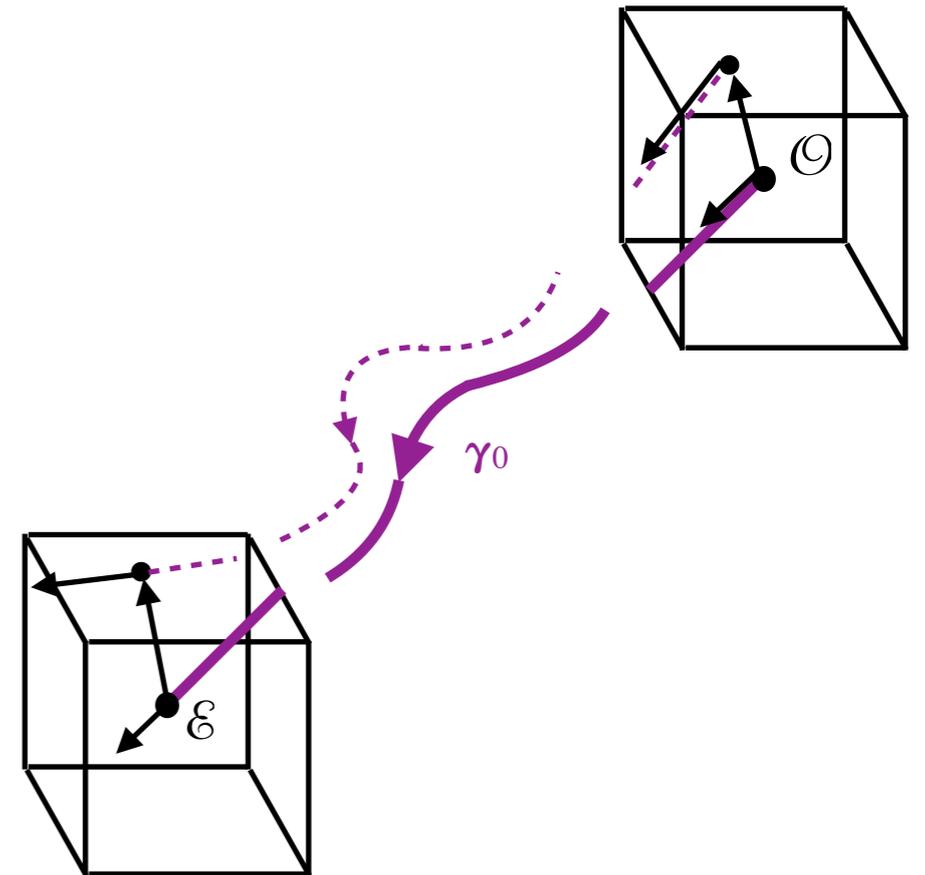
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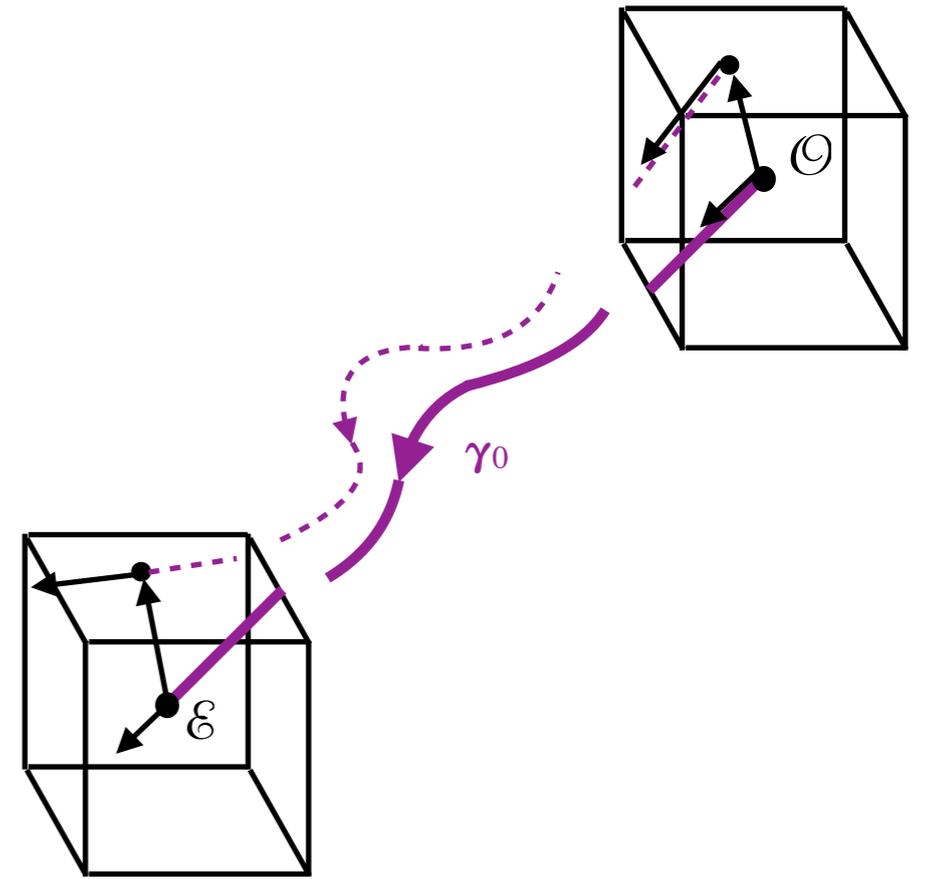
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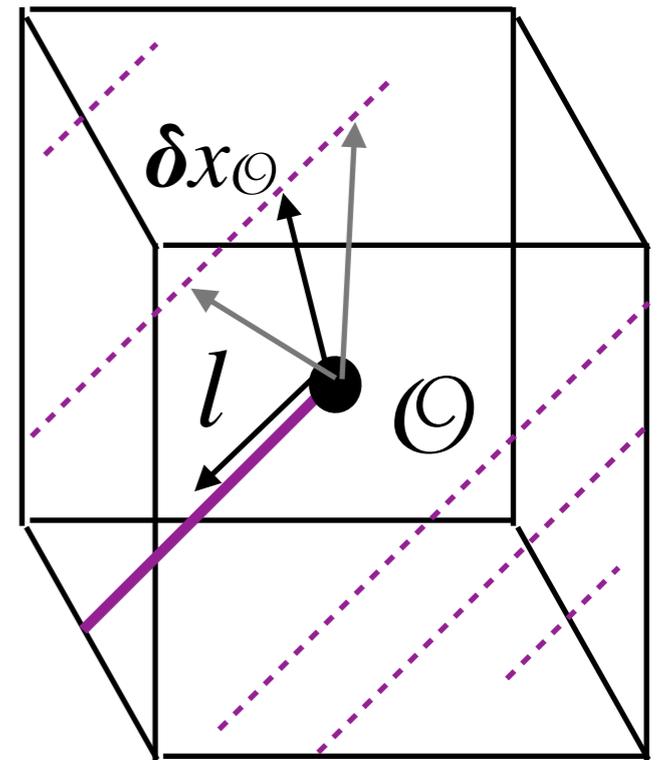
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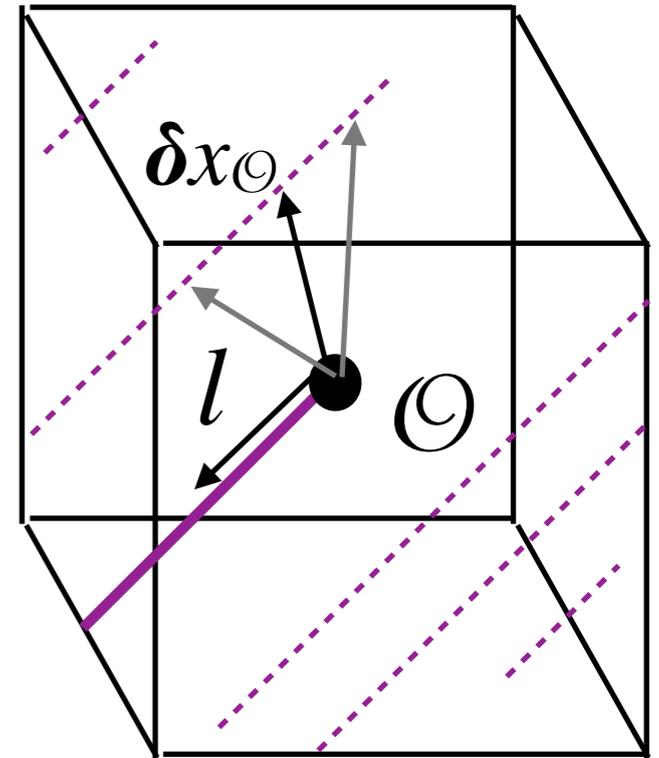


Geometry

$$[\delta x_{\mathcal{O}}] \in T_{\mathcal{O}}M / l$$

$$[\delta x_{\mathcal{O}}] = \{ X^{\mu} \in T_{\mathcal{O}}M \mid X^{\mu} = \delta x_{\mathcal{O}}^{\mu} + C l^{\mu} \}$$

$$\dim(T_{\mathcal{O}}M / l) = 3$$



Geometry

$$[\delta x_{\mathcal{O}}] \in T_{\mathcal{O}}M / l$$

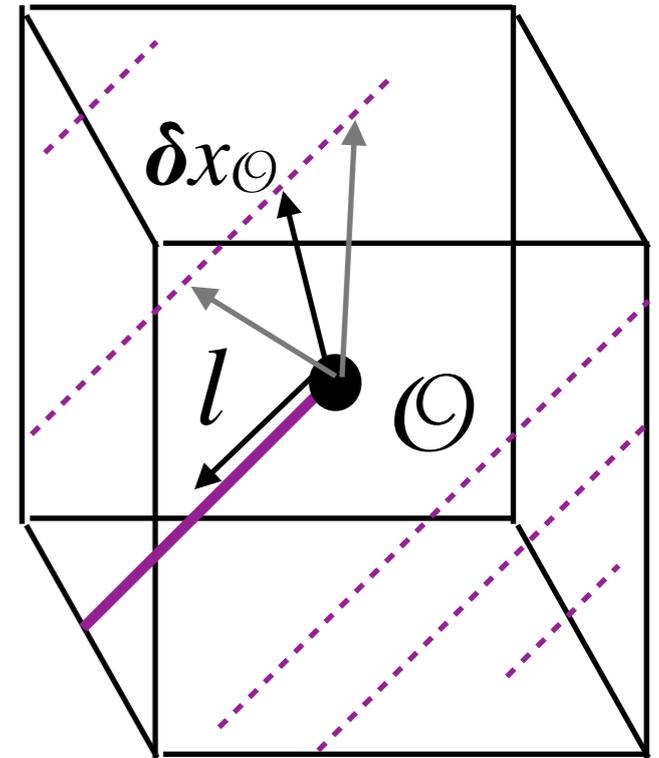
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$$\dim(T_{\mathcal{O}}M / l) = 3$$

$$[\Delta l_{\mathcal{O}}] \in l^{\perp} / l = \mathcal{P}_{\mathcal{O}}$$

$$[\Delta l_{\mathcal{O}}] = \{ X^{\mu} \in l^{\perp} \mid X^{\mu} = \Delta l_{\mathcal{O}}^{\mu} + C l^{\mu} \}$$

$$\dim \mathcal{P}_{\mathcal{O}} = 2$$



Geometry

$$[\delta x_{\mathcal{O}}] \in T_{\mathcal{O}}M / l$$

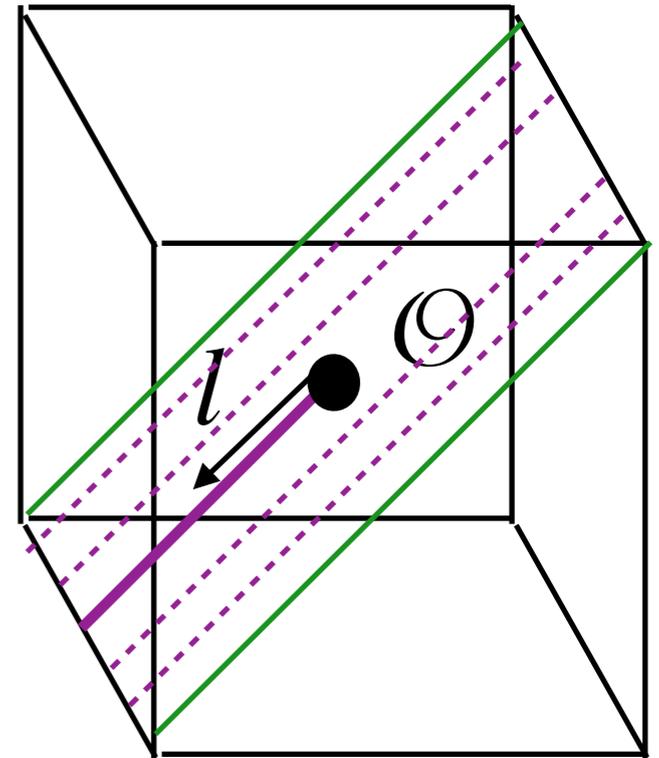
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Geometry

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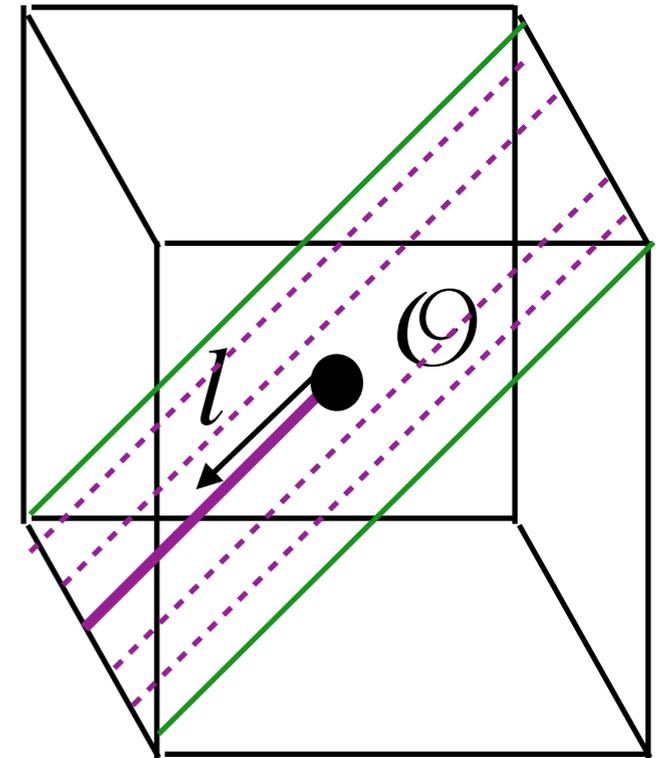
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- Perpendicular space - corresponds to direction deviations of null geodesics

Geometry

$$[\delta x_{\mathcal{O}}] \in T_{\mathcal{O}}M / l$$

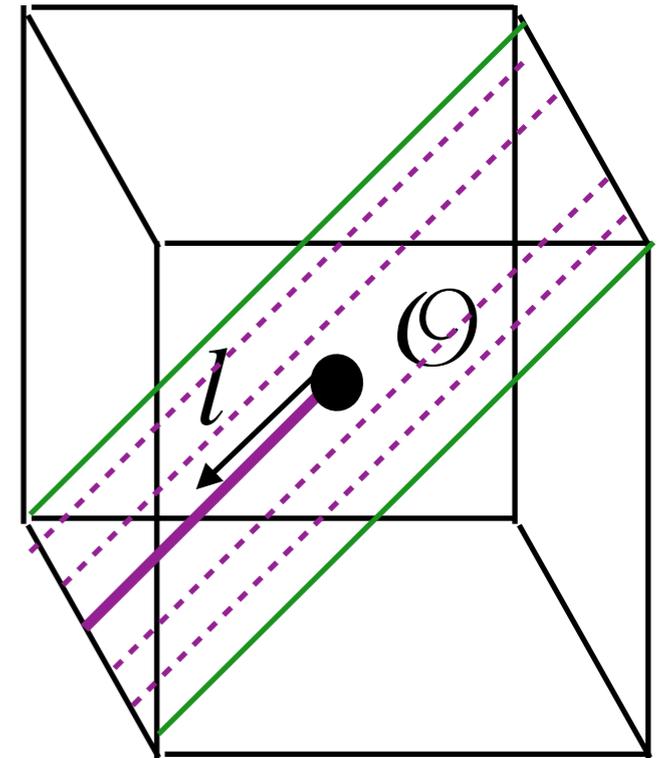
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- Perpendicular space - corresponds to direction deviations of null geodesics
- ...or to displacement of null geodesics on the same light front

Geometry

$$[\delta x_{\mathcal{O}}] \in T_{\mathcal{O}}M / l$$

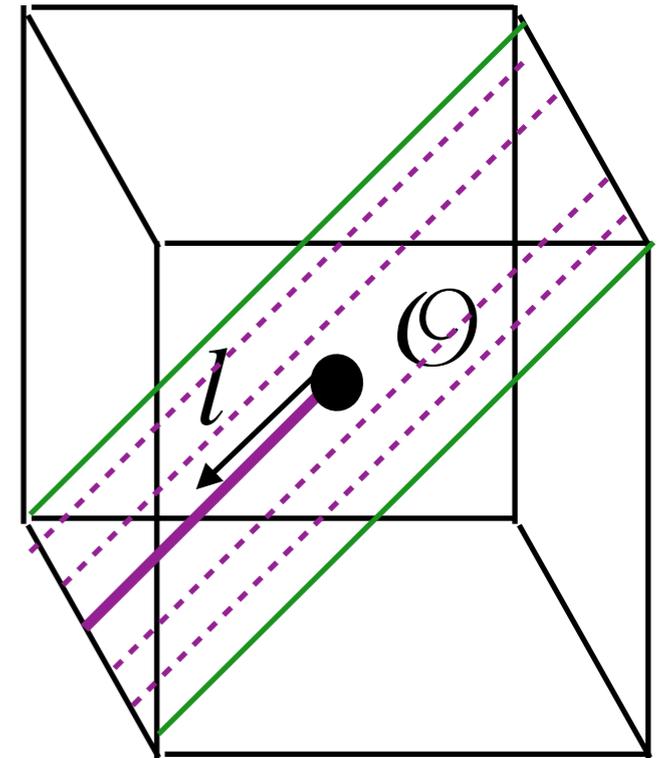
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- Perpendicular space - corresponds to direction deviations of null geodesics
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- inherits a positive definite metric $[g]$ from the spacetime metric

Geometry

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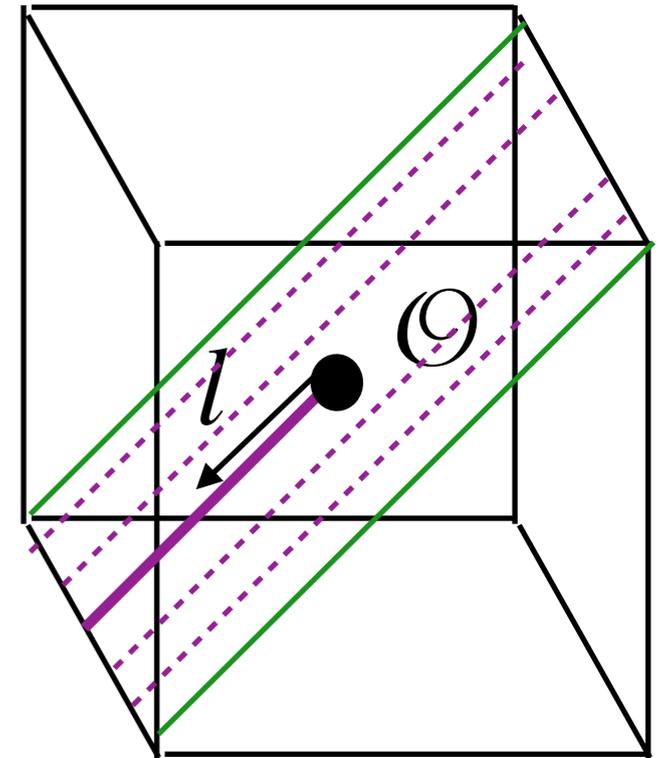
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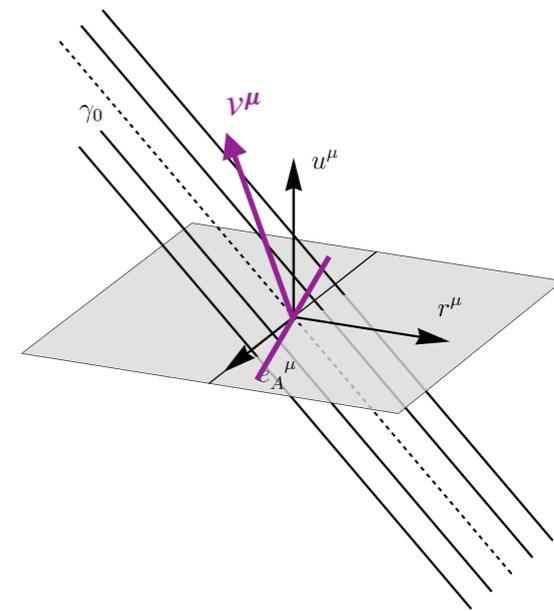
- Perpendicular space - corresponds to direction deviations of null geodesics
- ...or to displacement of null geodesics on the same light front
- inherits a positive definite metric $[g]$ from the spacetime metric
- can be identified with the screen space of any observer

Geometry

Perpendicular space \mathcal{P}

- orthogonally displaced null geodesics

$$\delta x_{\mathcal{O}}^{\mu} l_{\mu} = 0$$



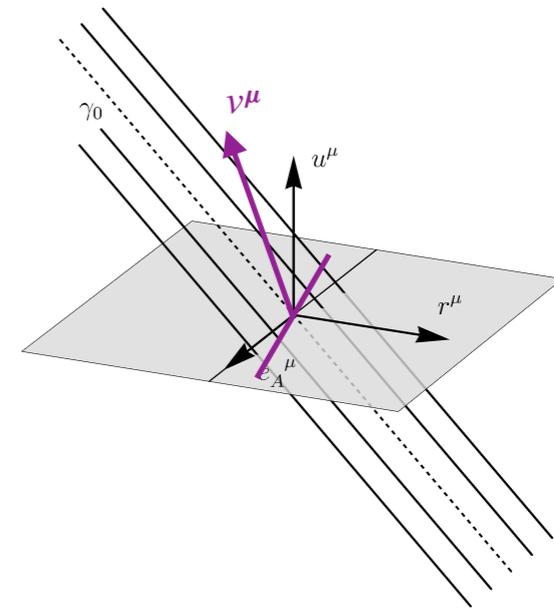
Geometry

Perpendicular space \mathcal{P}

Different observers u and $v \Rightarrow$ different notions of simultaneity

- orthogonally displaced null geodesics

$$\delta x_{\mathcal{O}}^{\mu} l_{\mu} = 0$$



Geometry

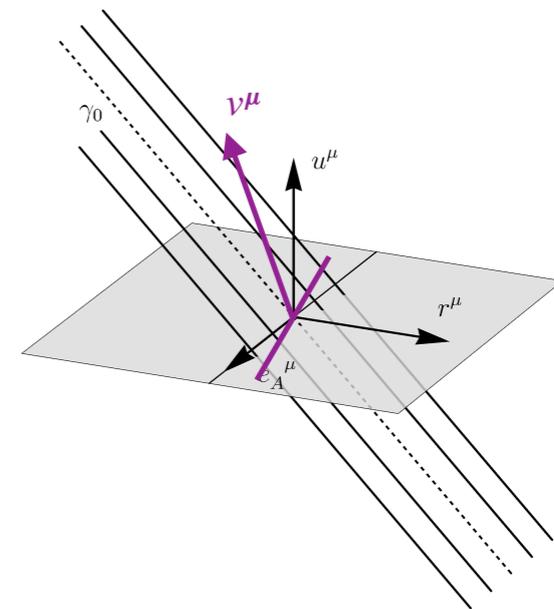
Perpendicular space \mathcal{P}

Different observers u and $v \Rightarrow$ different notions of simultaneity

\Rightarrow different screen spaces

- orthogonally displaced null geodesics

$$\delta x_{\mathcal{O}}^{\mu} l_{\mu} = 0$$



Geometry

Perpendicular space \mathcal{P}

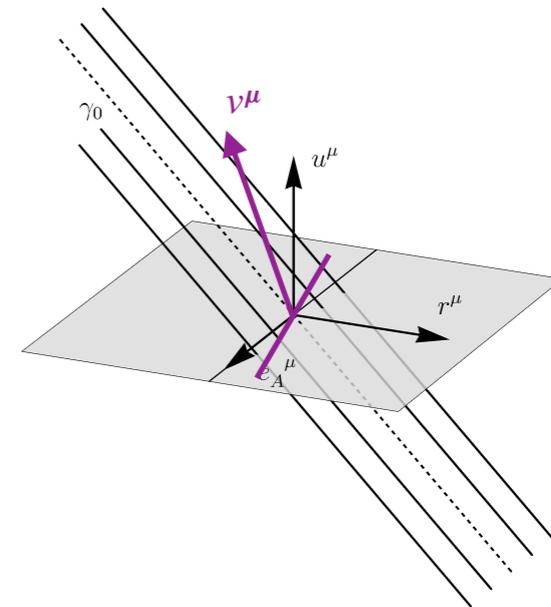
Different observers u and $v \Rightarrow$ different notions of simultaneity

\Rightarrow different screen spaces

\Rightarrow screen spaces punctured by light rays at different points

- orthogonally displaced null geodesics

$$\delta x_{\mathcal{O}}^{\mu} l_{\mu} = 0$$



Geometry

Perpendicular space \mathcal{P}

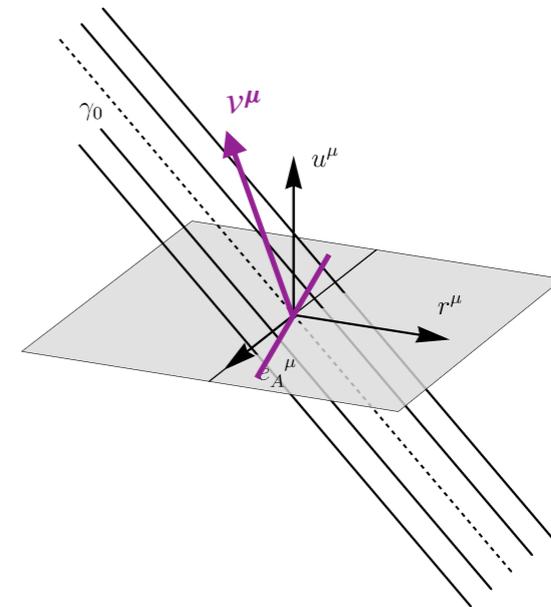
Different observers u and $v \Rightarrow$ different notions of simultaneity

\Rightarrow different screen spaces

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- orthogonally displaced null geodesics

$$\delta x_{\mathcal{O}}^{\mu} l_{\mu} = 0$$



Yet, the distances measured to a given light rays (+ angles) are *the same*

Geometry

Perpendicular space \mathcal{P}

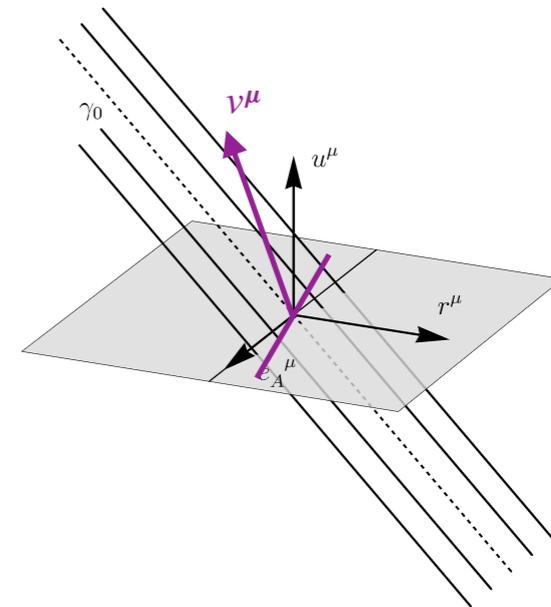
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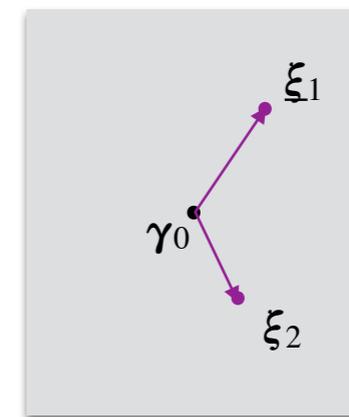
\Rightarrow screen spaces punctured by light rays at different points

- orthogonally displaced null geodesics

$$\delta x_{\mathcal{O}}^{\mu} l_{\mu} = 0$$



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Geometry

Perpendicular space \mathcal{P}

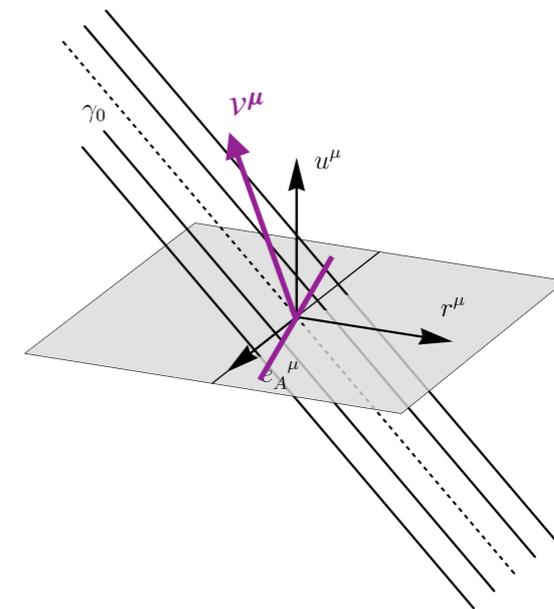
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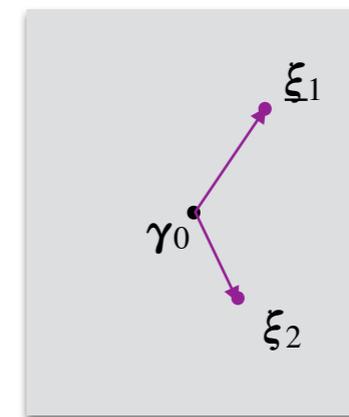
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Yet, the distances measured to a given light rays (+ angles) are *the same*

$$g_{\mu\nu} \xi_1^{\mu} \xi_2^{\nu} = g_{AB} \xi_1^A \xi_2^B$$

\nearrow \nwarrow \nwarrow
 \mathcal{P} \mathcal{P} \mathcal{P}



Geometry

Perpendicular space \mathcal{P}

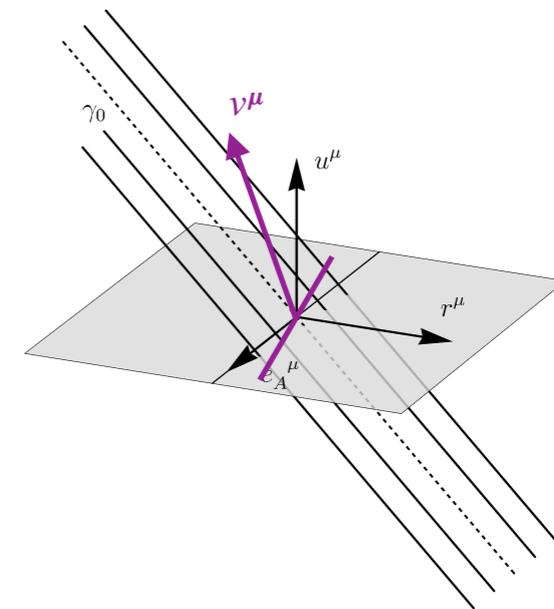
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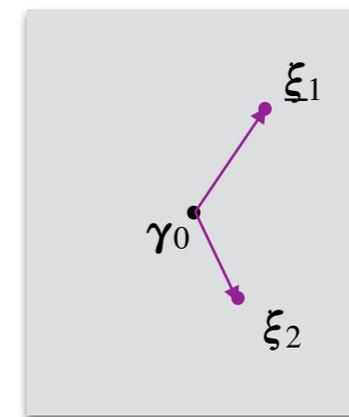
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Yet, the distances measured to a given light rays (+ angles) are *the same*

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\mathcal{P} \mathcal{P} \mathcal{P}

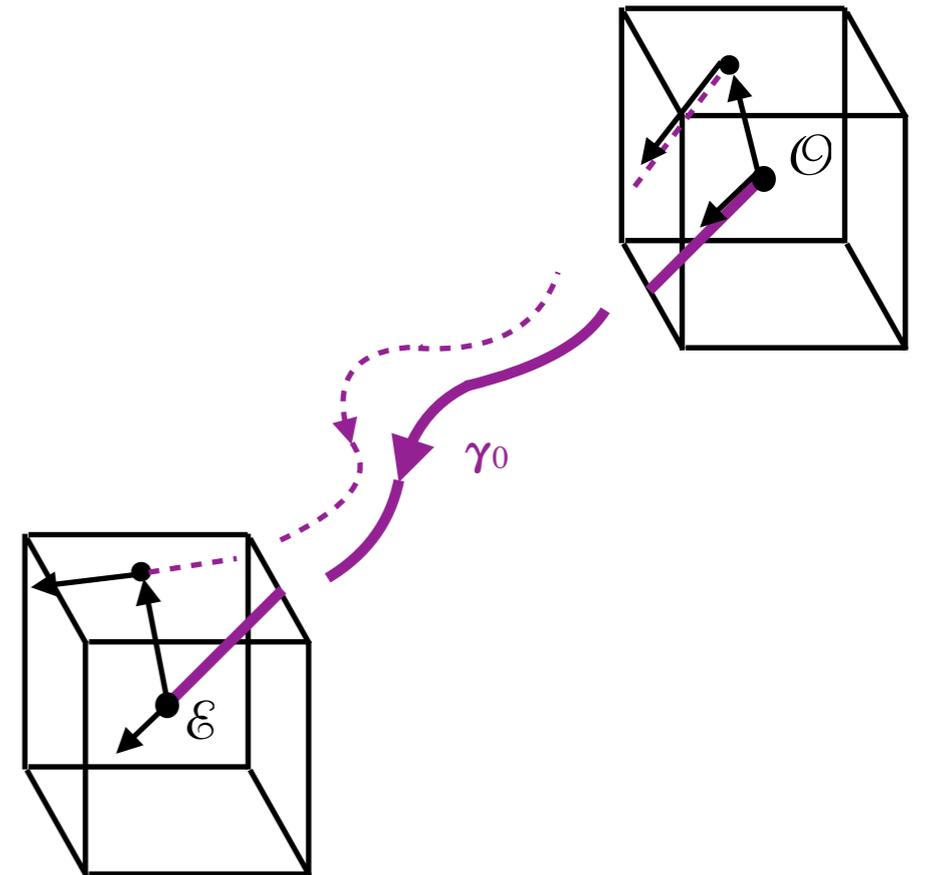


\mathcal{P} = identification of screen spaces of all observers

Displacement formula

$$\delta x_{\mathcal{O}}^{\mu} l_{\mu} = \delta x_{\mathcal{E}}^{\mu} l_{\mu}$$

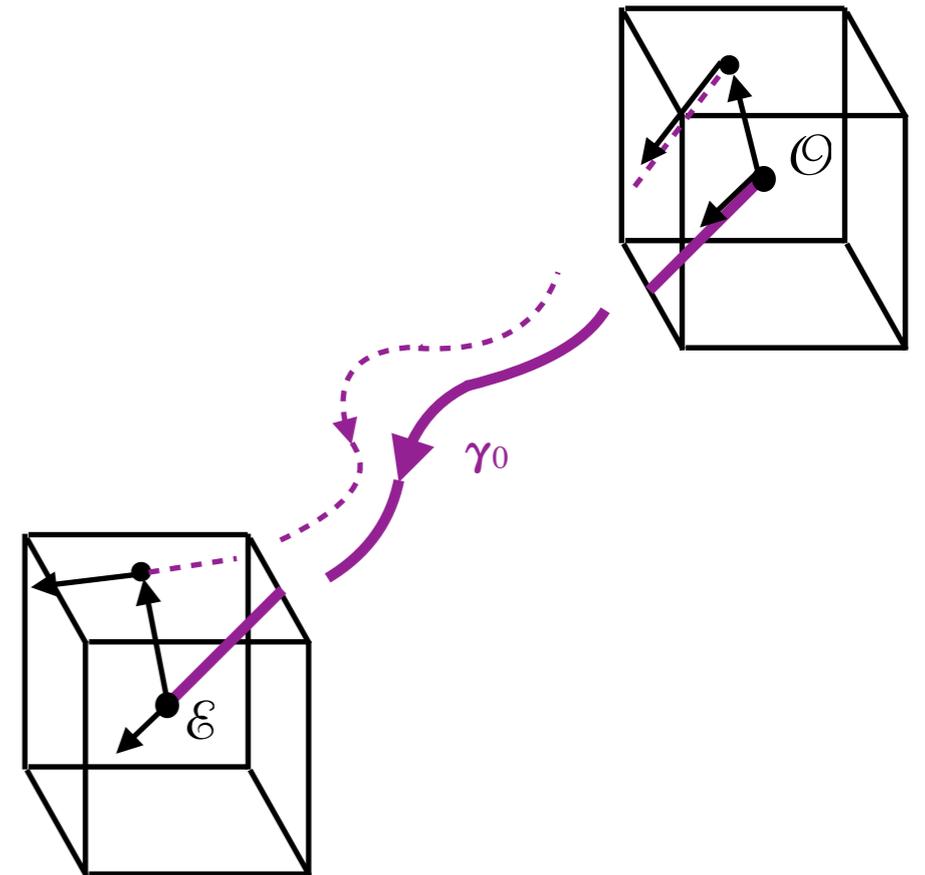
$$\delta x_{\mathcal{E}}^{\mu} = W_{xx}^{\mu\nu} \delta x_{\mathcal{O}}^{\nu} + W_{xl}^{\mu\nu} \Delta l_{\mathcal{O}}^{\nu}$$



Displacement formula

$$g([\delta x_{\mathcal{O}}], l) = g([\delta x_{\mathcal{E}}], l)$$

$$\mathcal{D}([\Delta l_{\mathcal{O}}]) = [\delta x_{\mathcal{E}} - \delta \hat{x}_{\mathcal{O}}] - m([\delta x_{\mathcal{O}}])$$



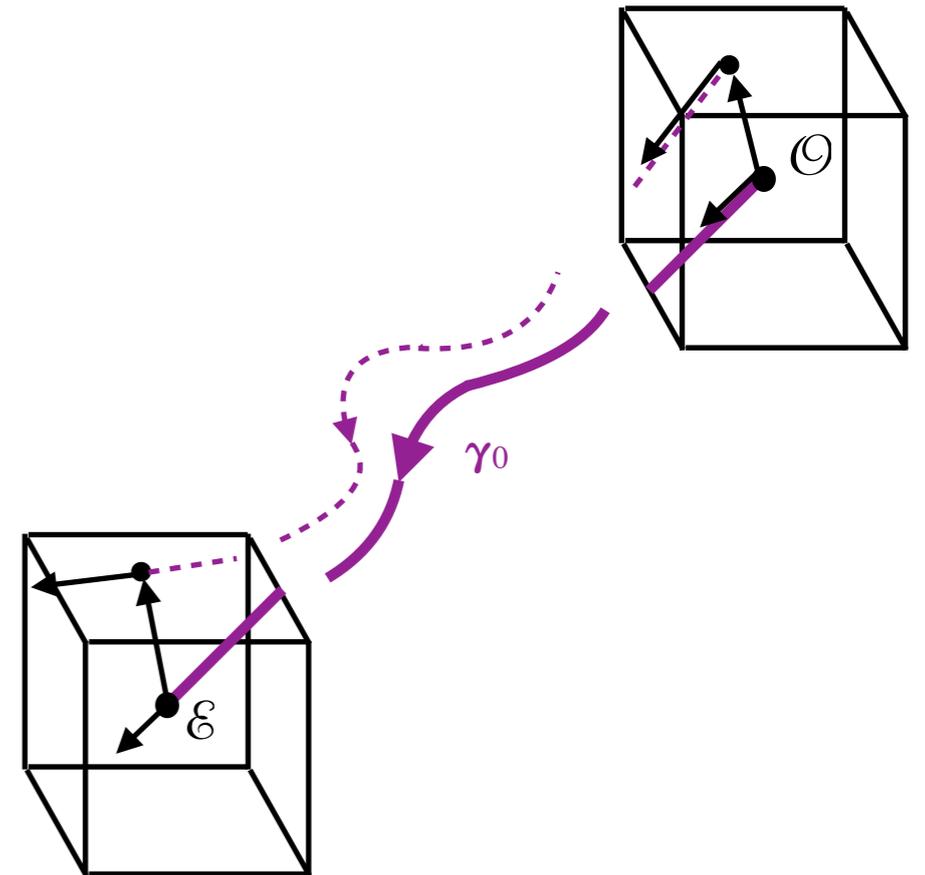
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where

- $\hat{\ }^{\wedge}$ - parallel transport from \mathcal{O} to \mathcal{E}



Displacement formula

$$g([\delta x_{\mathcal{O}}], l) = g([\delta x_{\mathcal{E}}], l)$$

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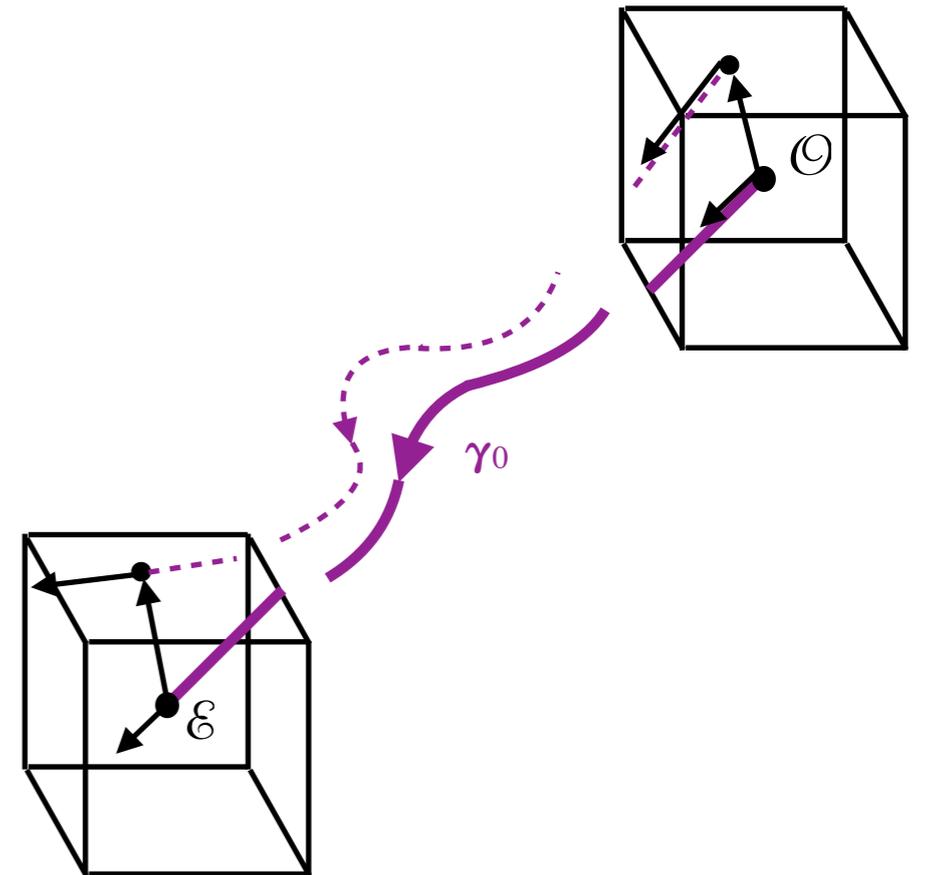
where

- $\hat{\cdot}$ - parallel transport from \mathcal{O} to \mathcal{E}
- Jacobi operator $\mathcal{D} : \mathcal{P}_{\mathcal{O}} \rightarrow \mathcal{P}_{\mathcal{E}}$

$$\ddot{\mathcal{D}}^A_B - R^A_{\mu\nu C} l^\mu l^\nu \mathcal{D}^C_B = 0$$

$$\mathcal{D}^A_B(\lambda_{\mathcal{O}}) = 0$$

$$\dot{\mathcal{D}}^A_B(\lambda_{\mathcal{O}}) = \delta^A_B$$



Displacement formula

$$g([\delta x_{\mathcal{O}}], l) = g([\delta x_{\mathcal{E}}], l)$$

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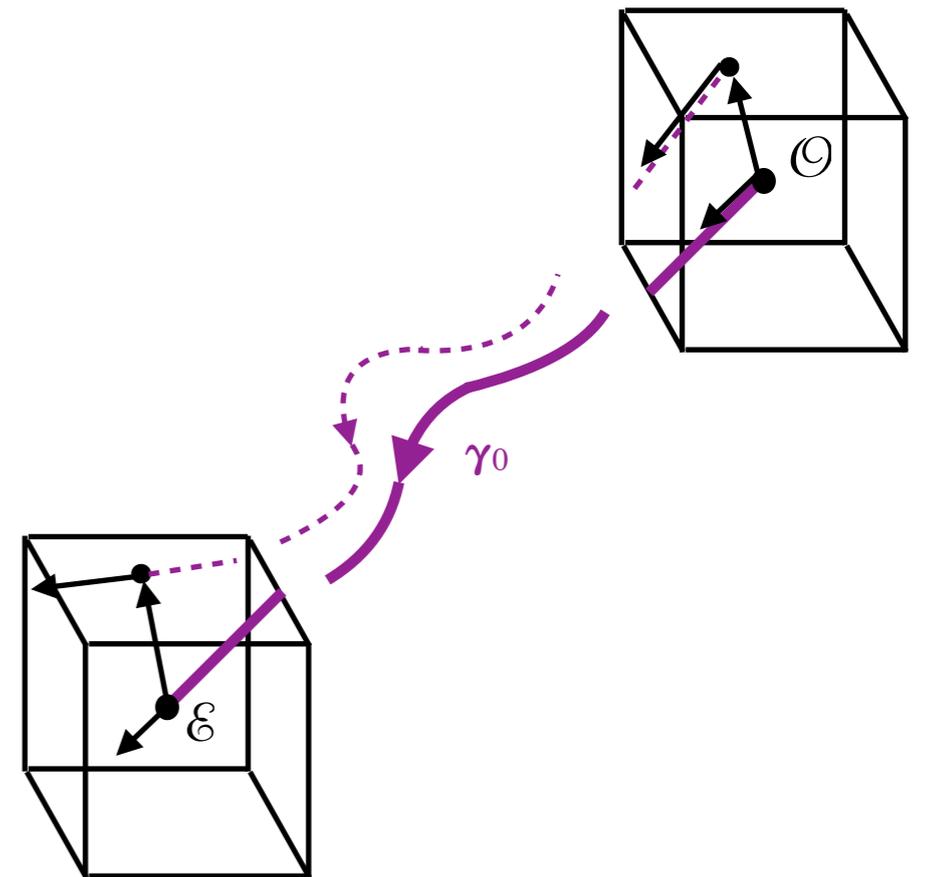
$$\dot{\mathcal{D}}^A_B(\lambda_{\mathcal{O}}) = \delta^A_B$$

- \mathcal{O}/\mathcal{E} asymmetry operator $m : T_{\mathcal{O}}M / l \rightarrow \mathcal{P}_{\mathcal{E}}$

$$\ddot{m}^A_\sigma - R^A_{\mu\nu C} l^\mu l^\nu m^C_\sigma = R^A_{\mu\nu\sigma} l^\mu l^\nu$$

$$m^A_\mu(\lambda_{\mathcal{O}}) = 0$$

$$\dot{m}^A_\mu(\lambda_{\mathcal{O}}) = 0$$



Displacement formula

$$g([\delta x_{\mathcal{O}}], l) = g([\delta x_{\mathcal{E}}], l)$$

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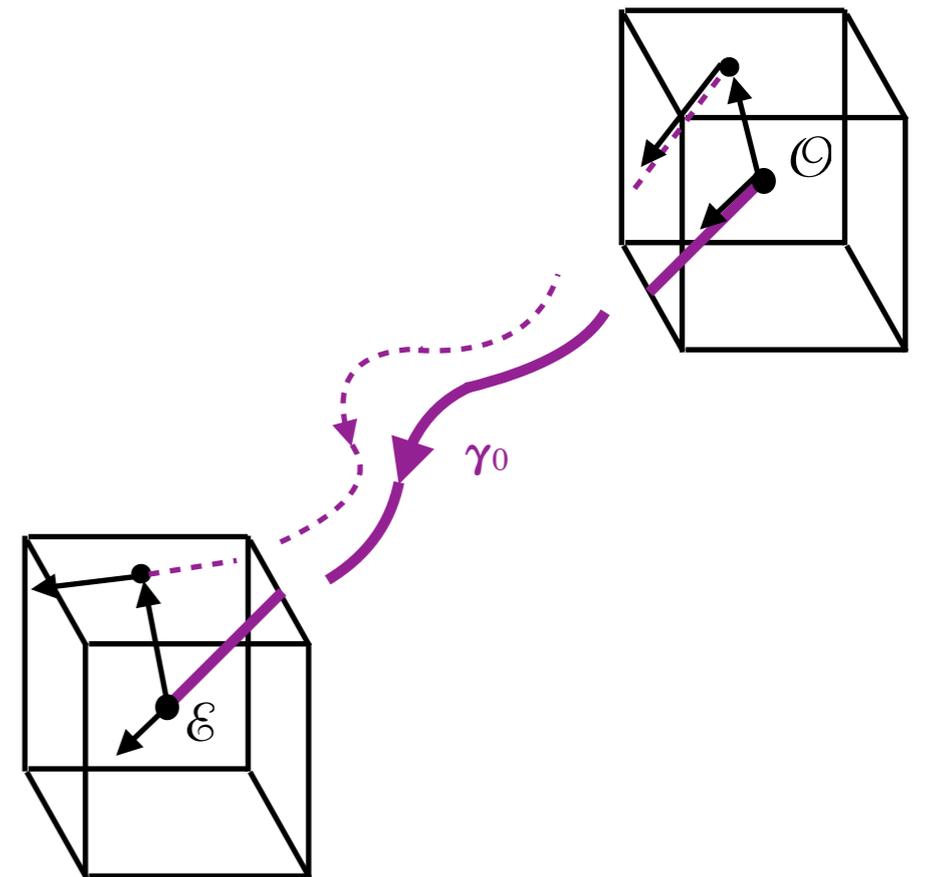
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$$m^A_\mu(\lambda_{\mathcal{O}}) = 0$$

$$\dot{m}^A_\mu(\lambda_{\mathcal{O}}) = 0$$

- vanishes in a flat space!



Displacement formula

$$g([\delta x_{\mathcal{O}}], l) = g([\delta x_{\mathcal{E}}], l)$$

$$\mathcal{D}([\Delta l_{\mathcal{O}}]) = [\delta x_{\mathcal{E}} - \delta \hat{x}_{\mathcal{O}}] - m([\delta x_{\mathcal{O}}])$$

where

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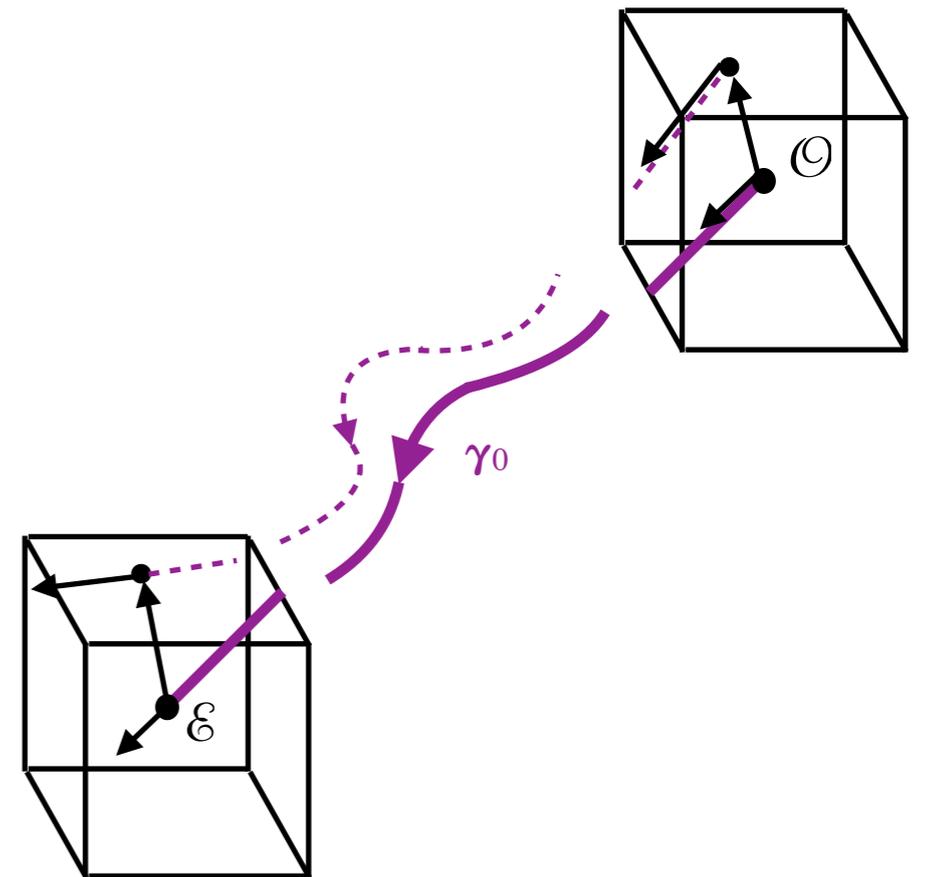
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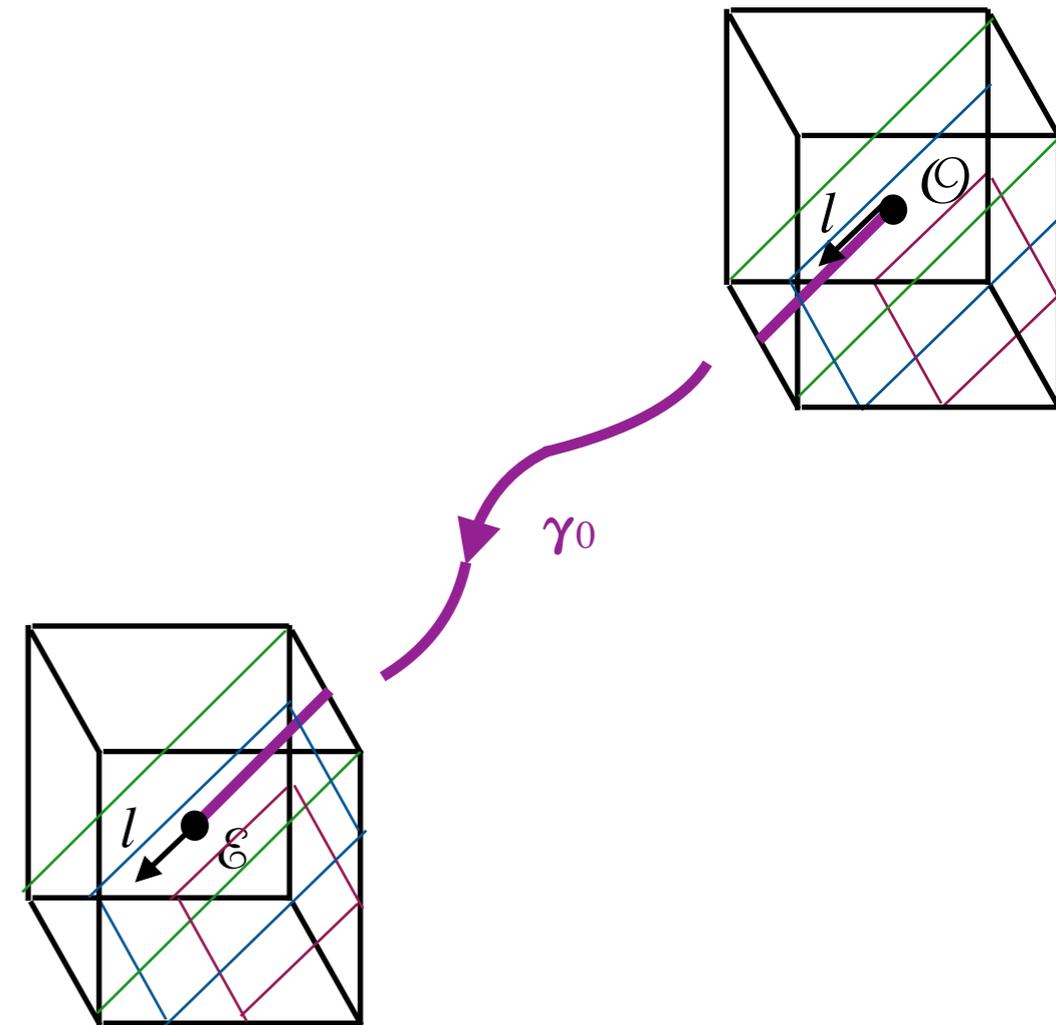


covariant, frame- and coordinate system-independent description of how observers in $N_{\mathcal{O}}$ see what is happening in $N_{\mathcal{E}}$

Displacement formula

$$g([\delta x_{\mathcal{O}}], l) = g([\delta x_{\mathcal{E}}], l)$$

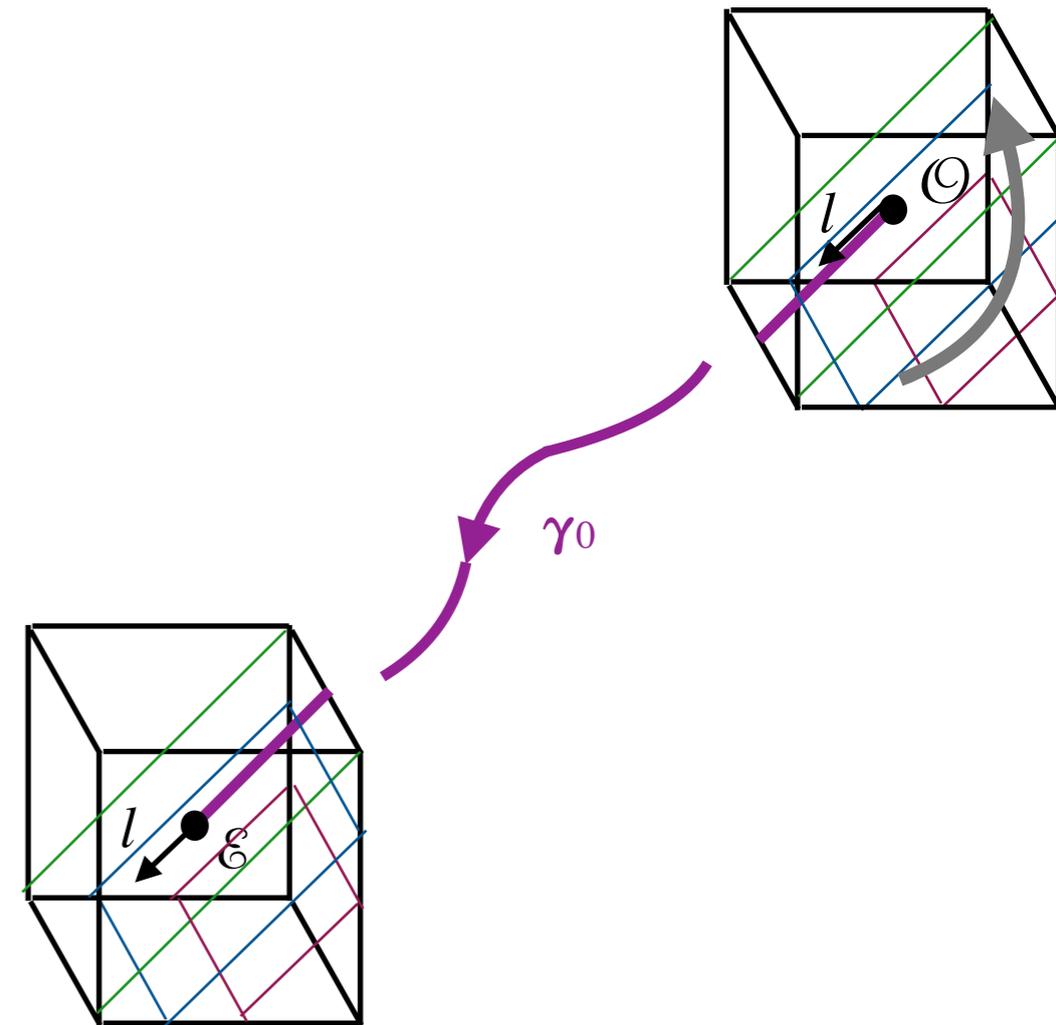
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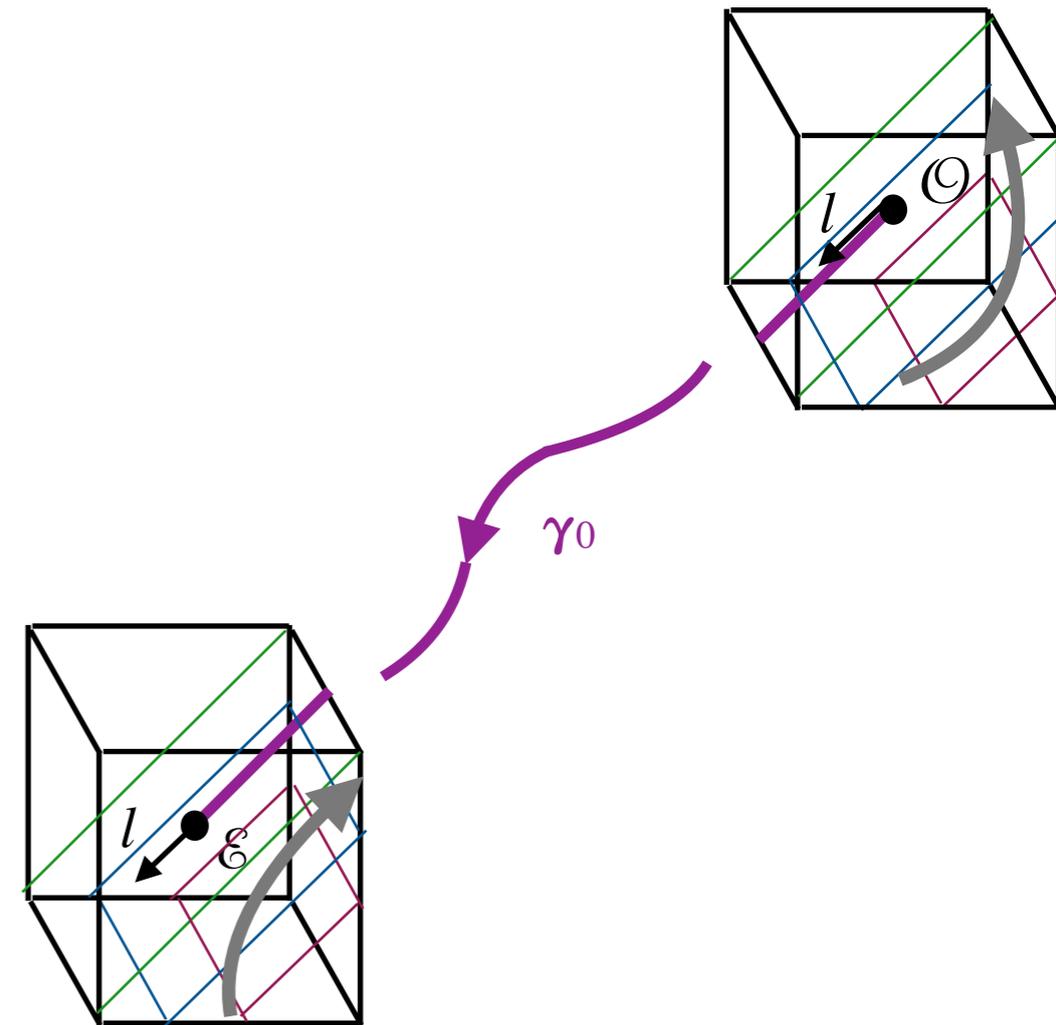
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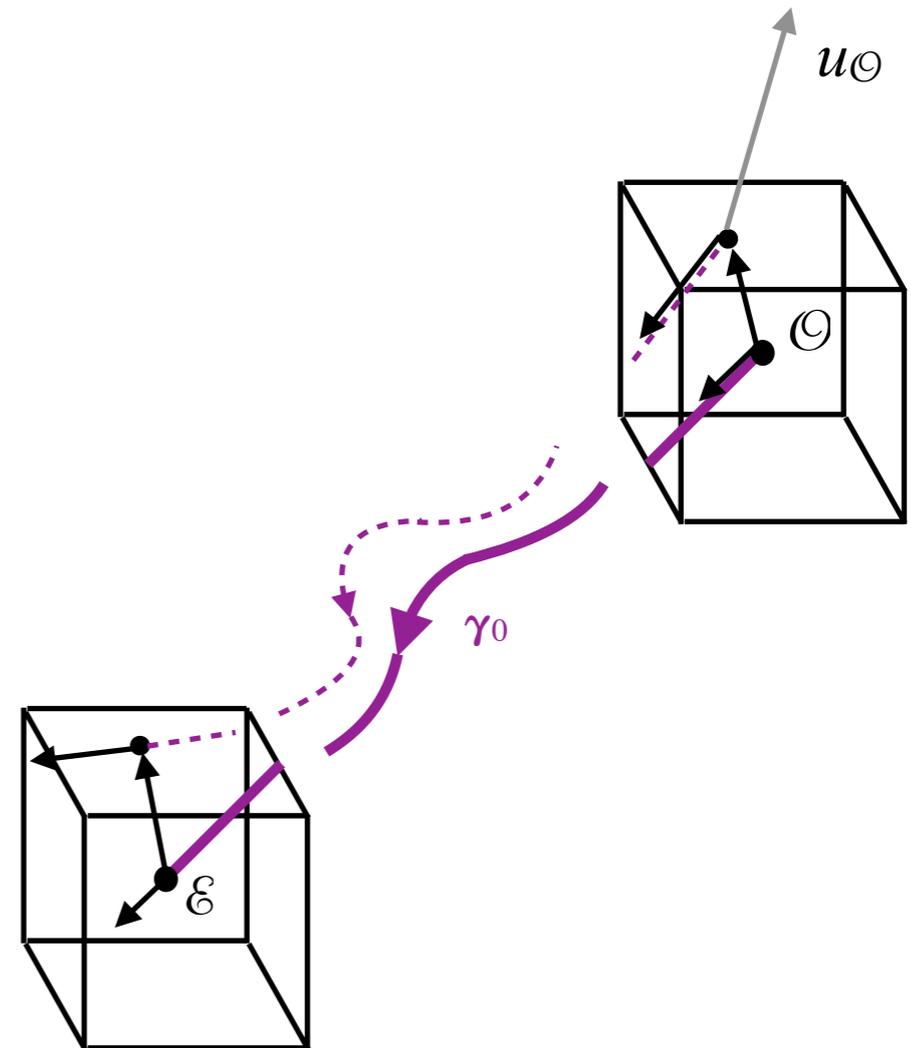
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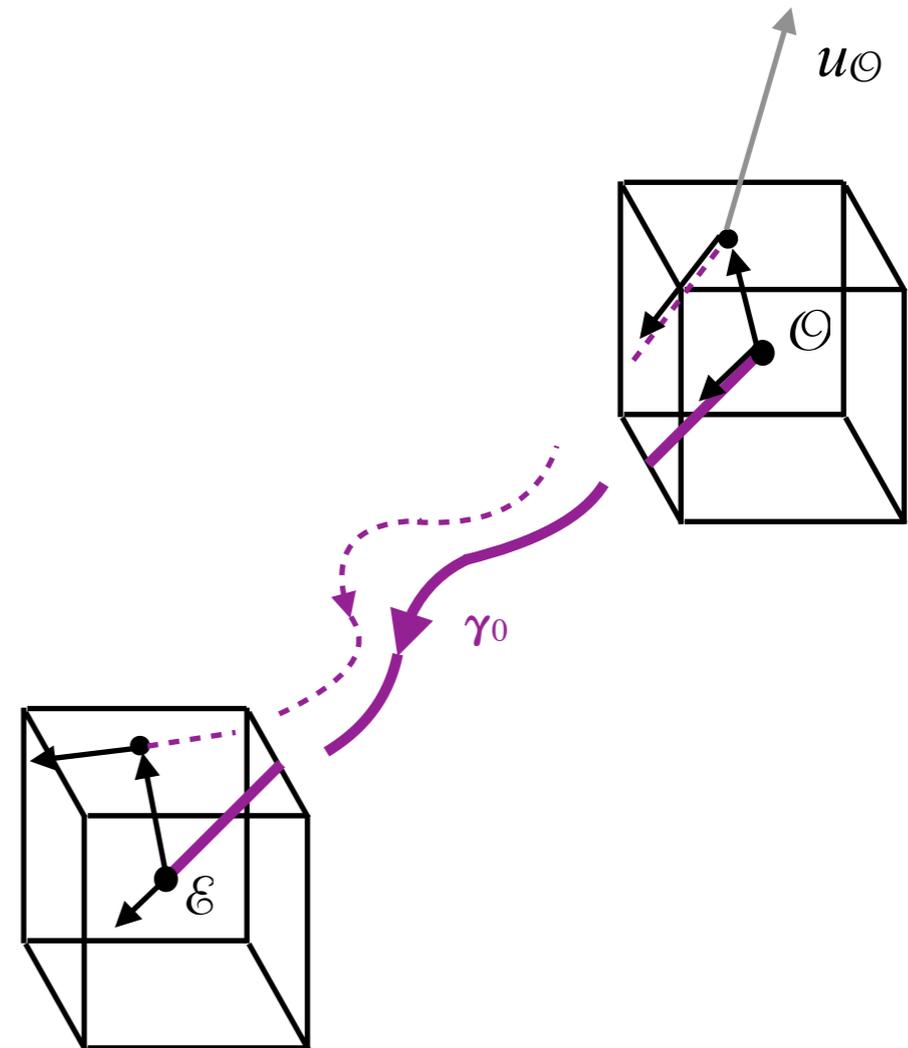


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propagation through curved spacetime (GR)



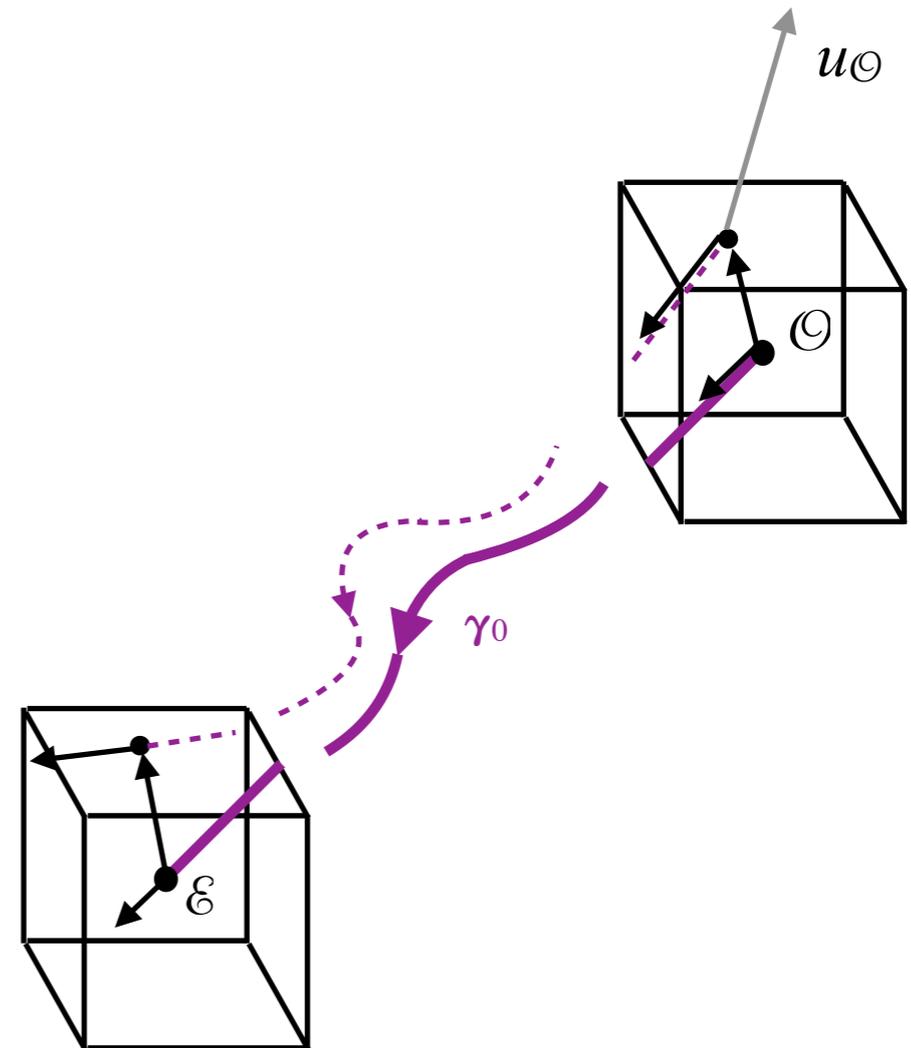
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propagation through curved spacetime (GR)

angles on the celestial sphere of observer $u_{\mathcal{O}}$



Displacement formula

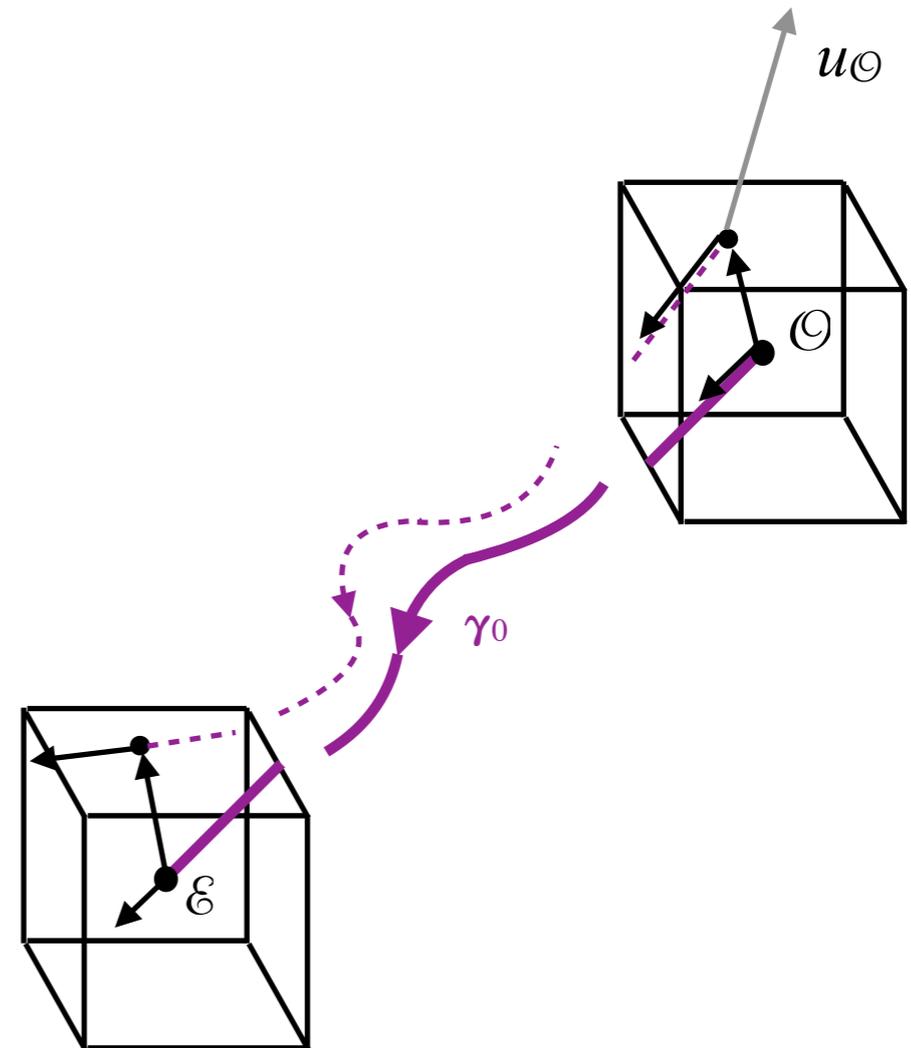
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propagation through curved spacetime (GR)

angles on the celestial sphere of observer $u_{\mathcal{O}}$

$$\delta \theta^A \approx \frac{1}{l_{\sigma} u_{\mathcal{O}}^{\sigma}} \Delta l^A$$



Displacement formula

$$g([\delta x_{\mathcal{O}}], l) = g([\delta x_{\mathcal{E}}], l)$$

$$\mathcal{D}([\Delta l_{\mathcal{O}}]) = [\delta x_{\mathcal{E}} - \delta \hat{x}_{\mathcal{O}}] - m([\delta x_{\mathcal{O}}])$$

propagation through curved spacetime (GR)

angles on the celestial sphere of observer $u_{\mathcal{O}}$

$$\delta \theta^A \approx \frac{1}{l_{\sigma} u_{\mathcal{O}}^{\sigma}} \Delta l^A$$

aberration effects (SR)

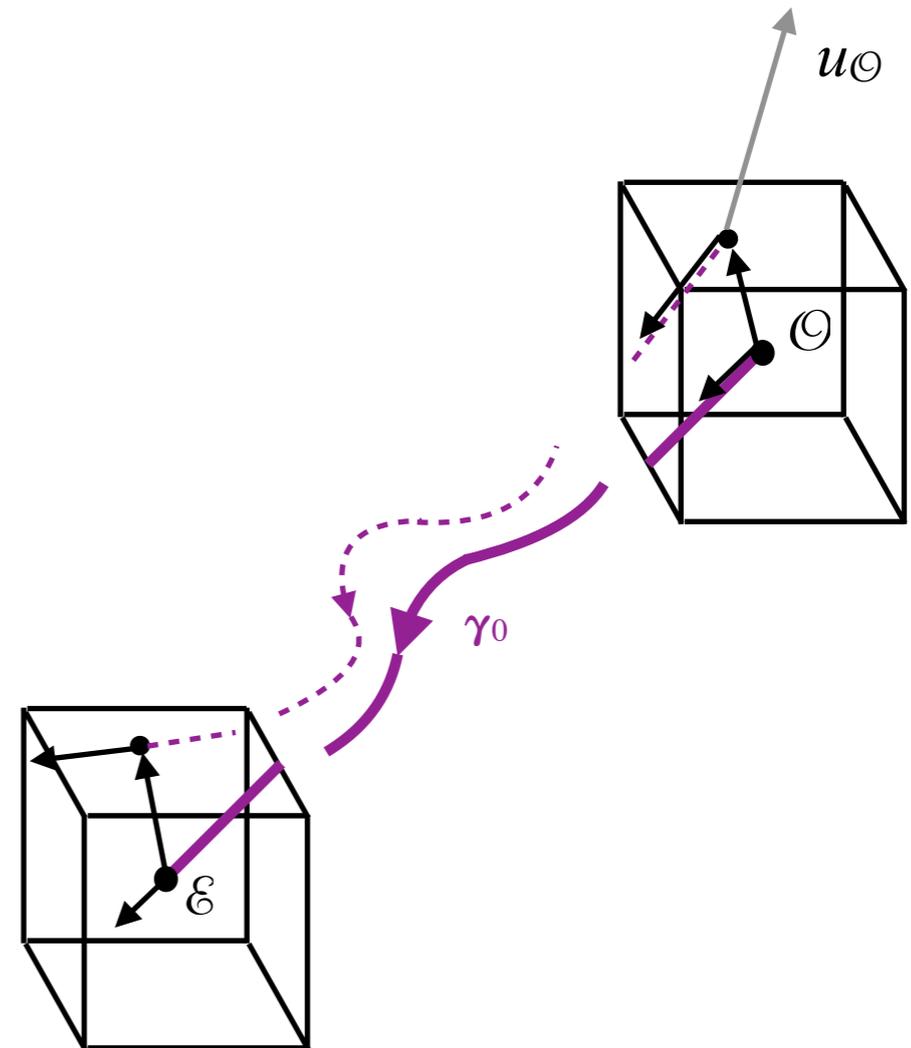
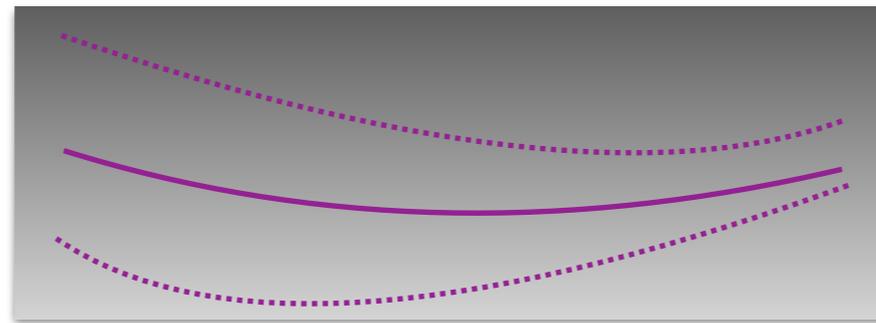
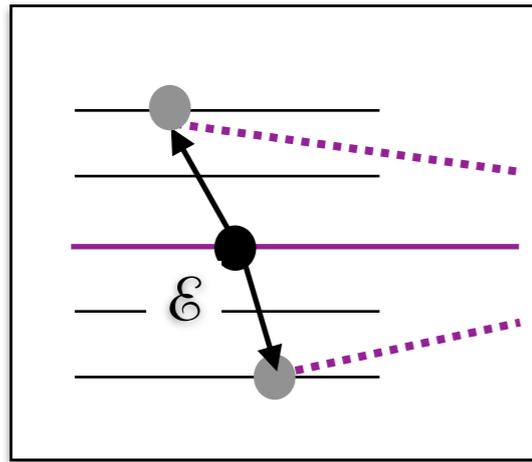


Image distortion



$$\mathcal{D}^A_B \Delta l^B_O = \delta x^A_\xi$$

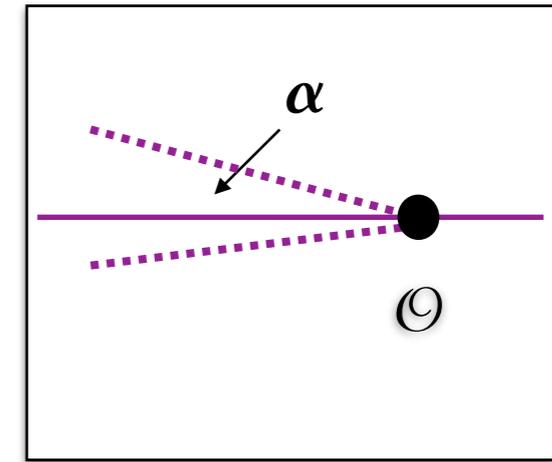
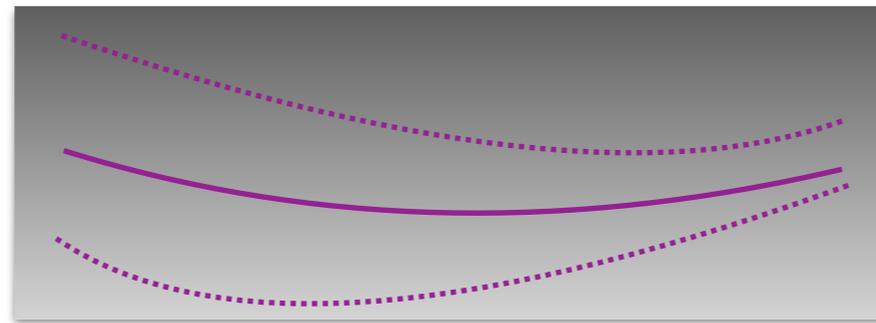
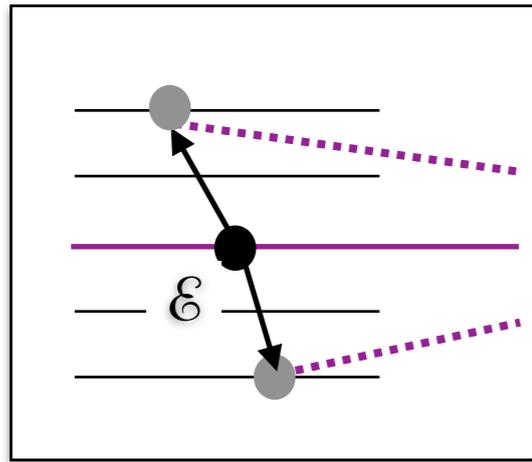
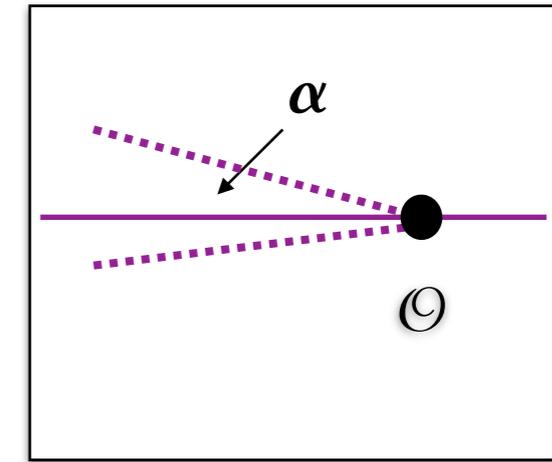


Image distortion

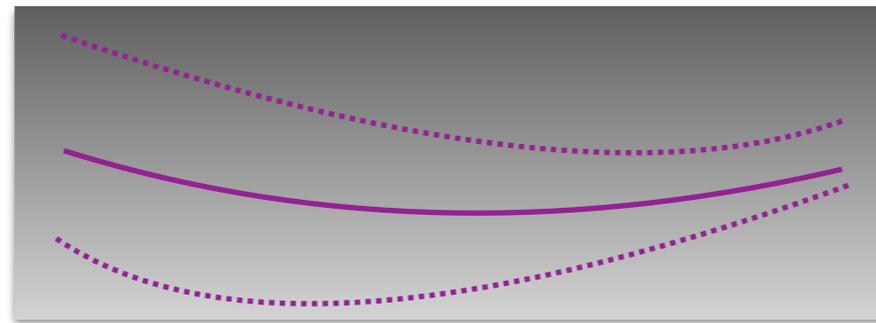
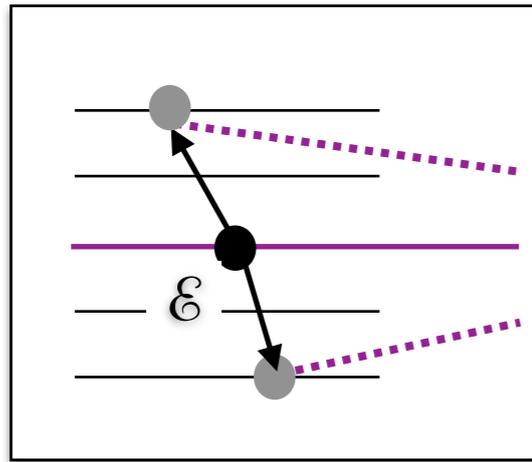


$$\mathcal{D}^A_B \Delta l^B_{\mathcal{O}} = \delta x^A_{\mathcal{E}}$$

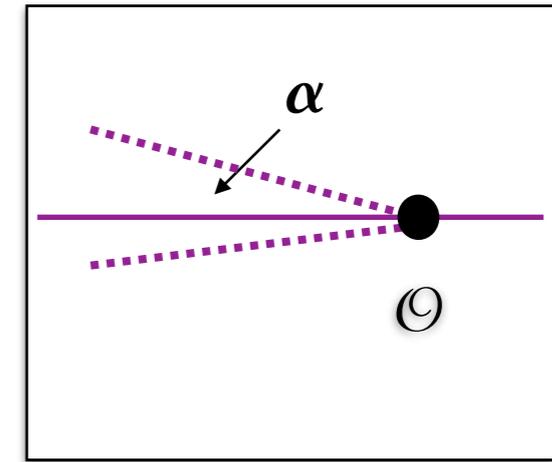


$$\delta\theta^A = \frac{1}{l_{\sigma} u^{\sigma}_{\mathcal{O}}} \mathcal{D}^{-1A}_B \delta x^B_{\mathcal{E}}$$

Image distortion



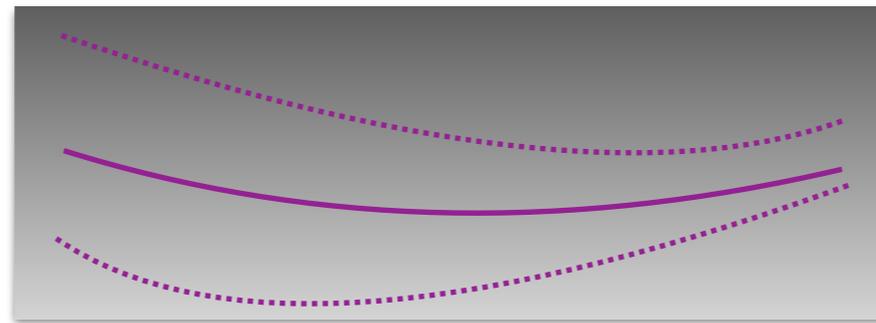
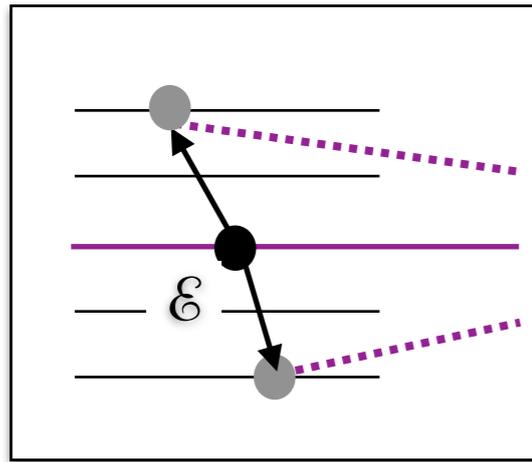
$$\mathcal{D}^A_B \Delta l^B_{\mathcal{O}} = \delta x^A_{\mathcal{E}}$$



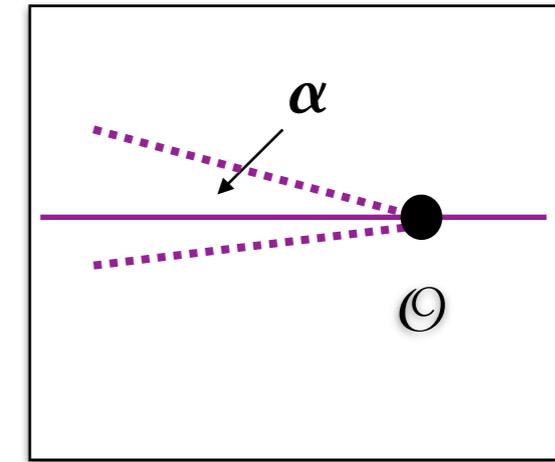
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M^A_B magnification matrix

Image distortion



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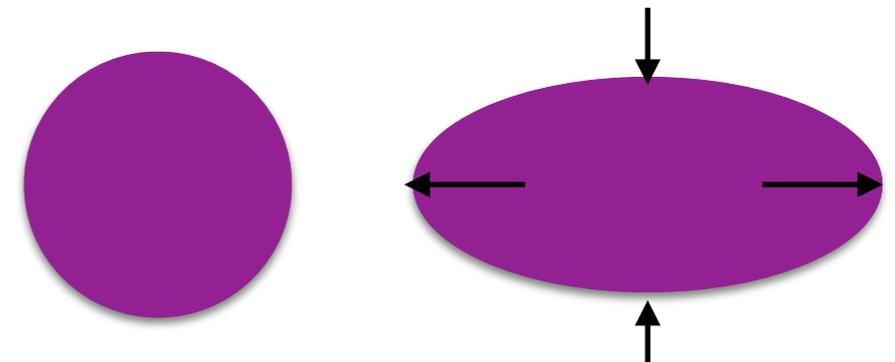
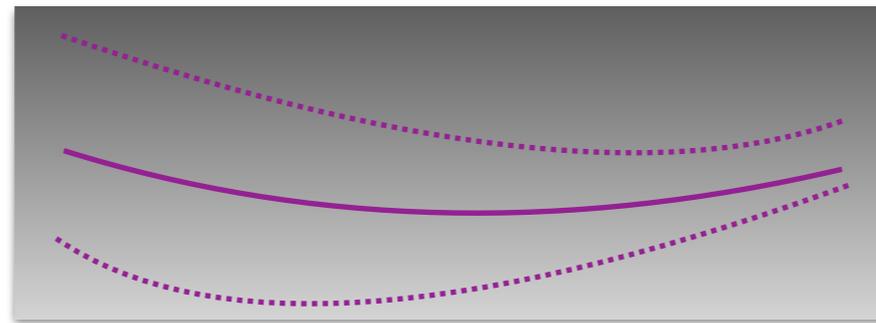
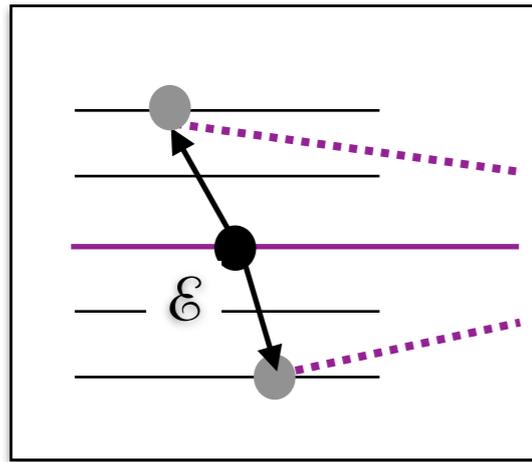
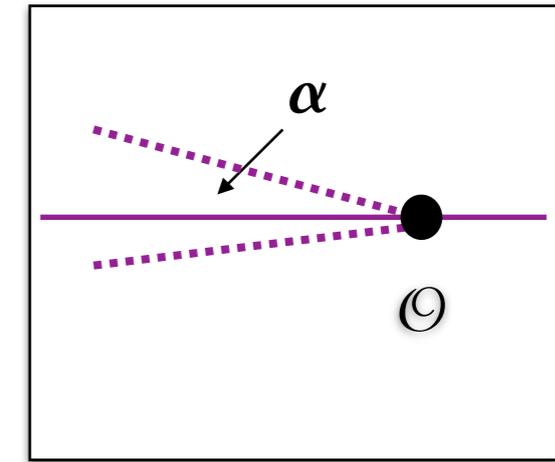


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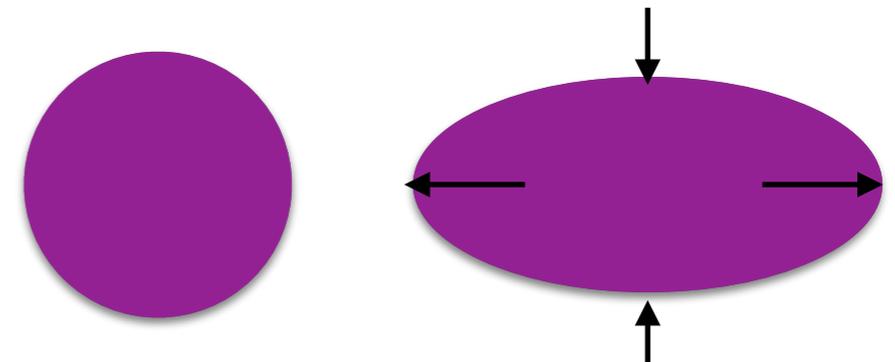


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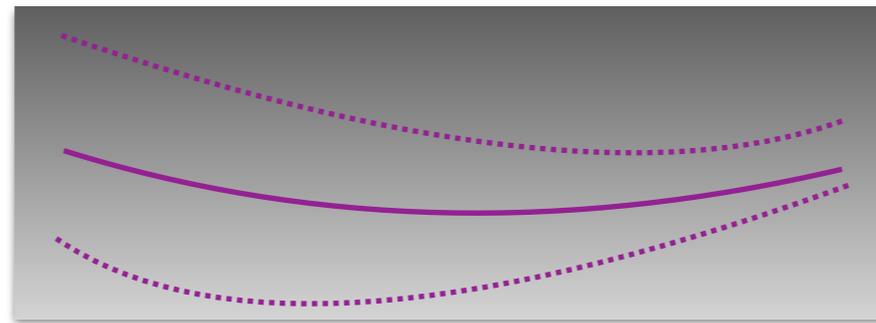
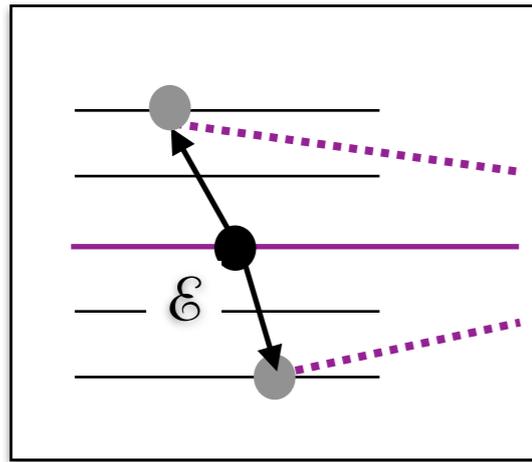
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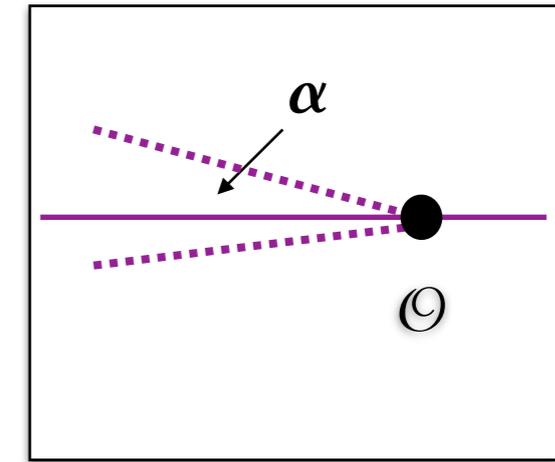


$$D_{ang} = (l_{\sigma} u^{\sigma}_{\mathcal{O}}) |\det \mathcal{D}^A_B|^{1/2}$$

Image distortion

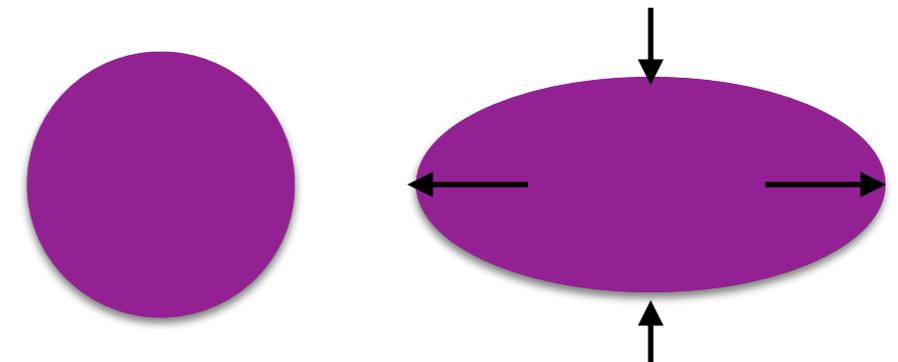


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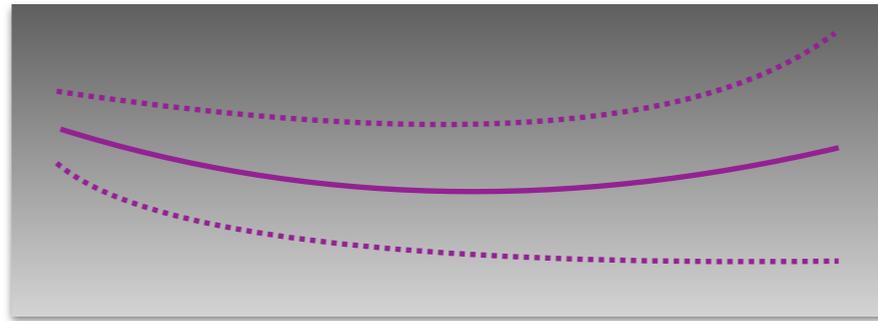
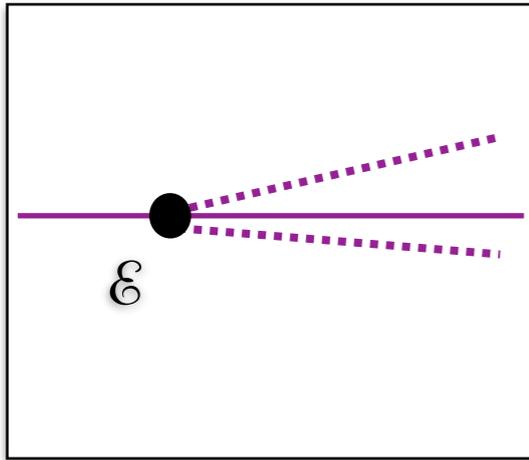
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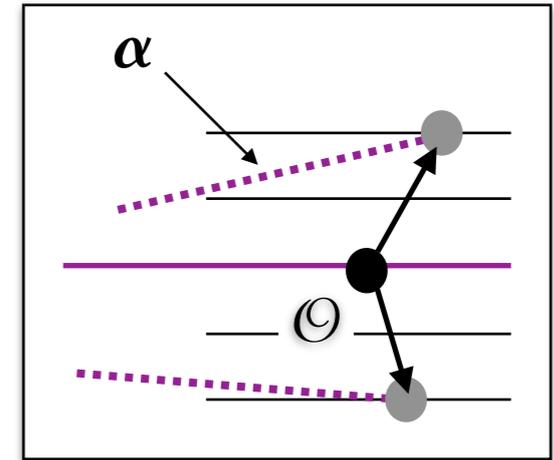
$$D_{ang} = (l_{\sigma} u^{\sigma}_{\mathcal{O}}) |\det \mathcal{D}^A_B|^{1/2}$$

depend on $u_{\mathcal{O}}$, but NOT on $u_{\mathcal{E}}$!

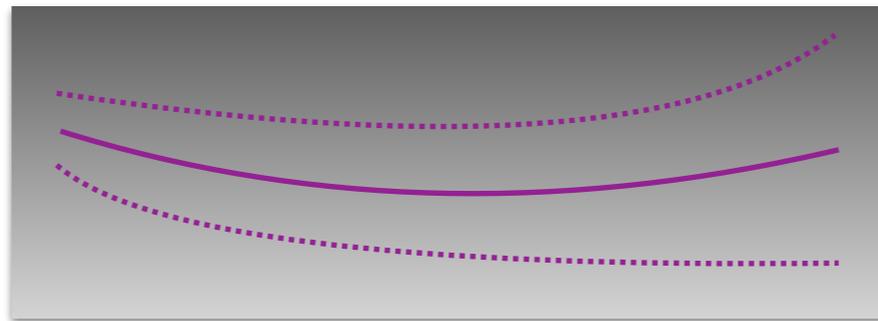
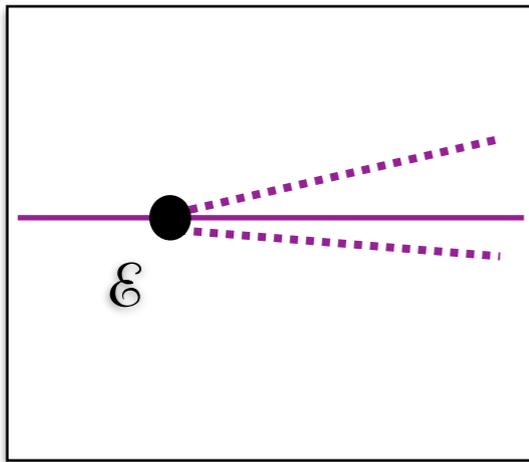
Stereoscopic parallax



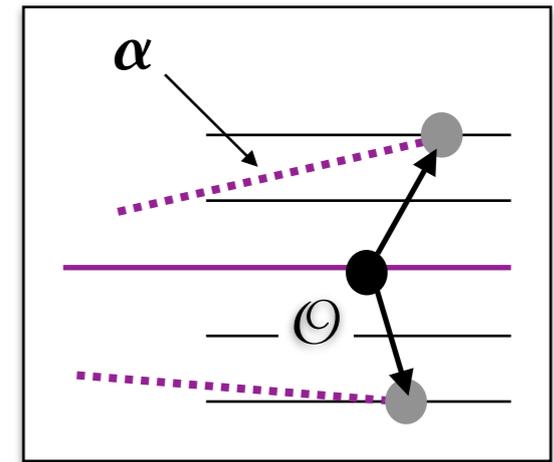
$$\mathcal{D}^A_B \Delta l_O^B = -\delta \hat{x}_O^A - m^A_B \delta x_O^B$$



Stereoscopic parallax

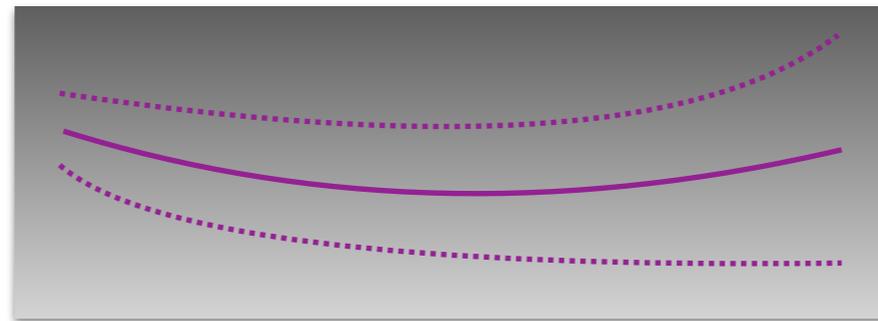
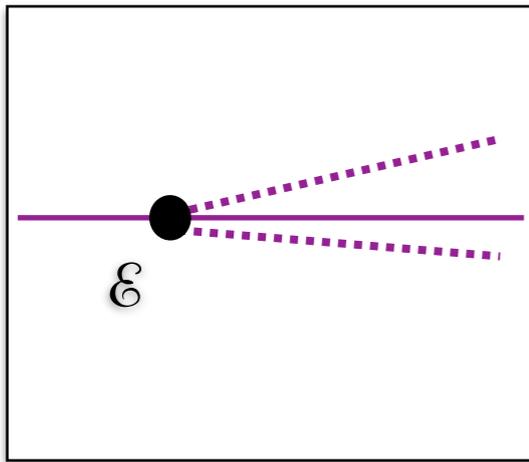


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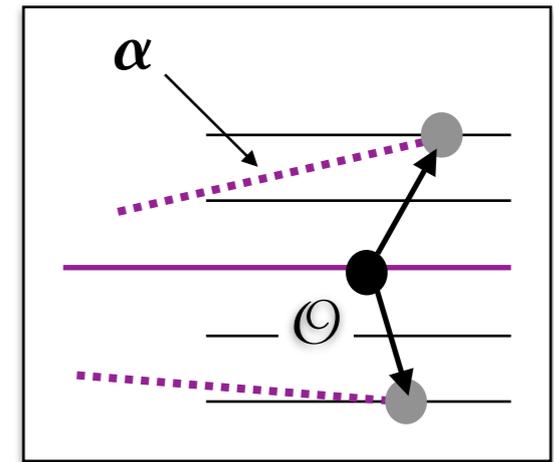


$$\delta \theta^A = -\frac{1}{l_{\sigma} u^{\sigma}_{\mathcal{O}}} \mathcal{D}^{-1A}_C (\delta^C_B + m^C_B) \delta x^B_{\mathcal{O}}$$

Stereoscopic parallax



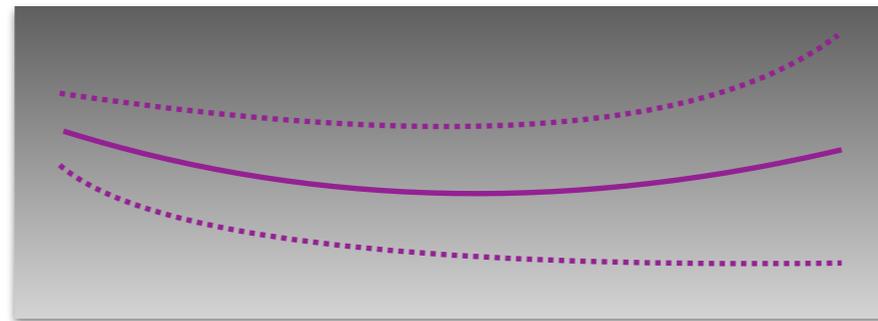
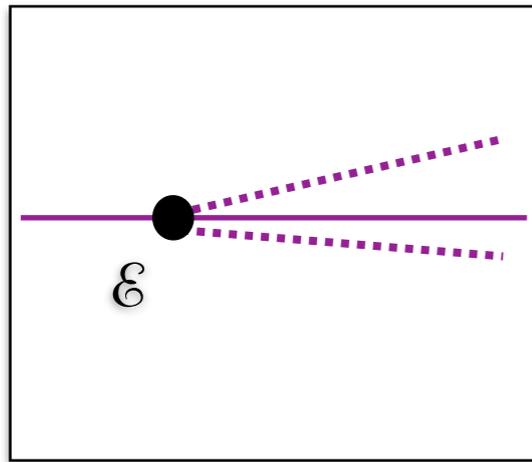
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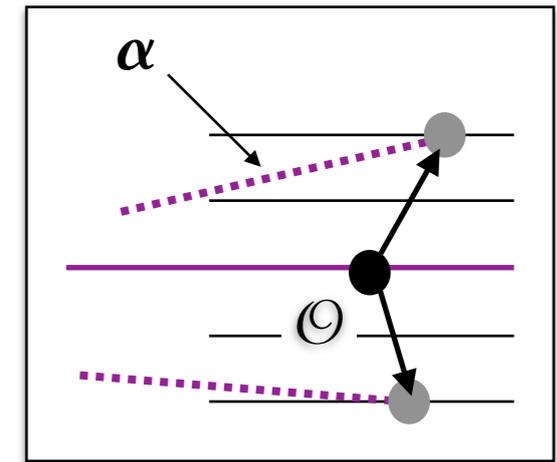
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Π^A_B parallax matrix

Stereoscopic parallax



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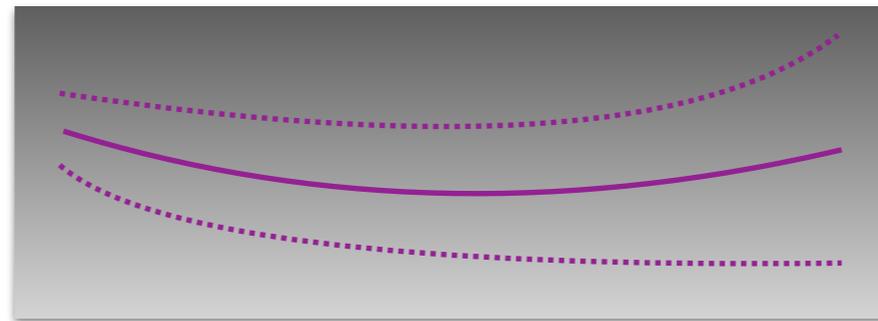
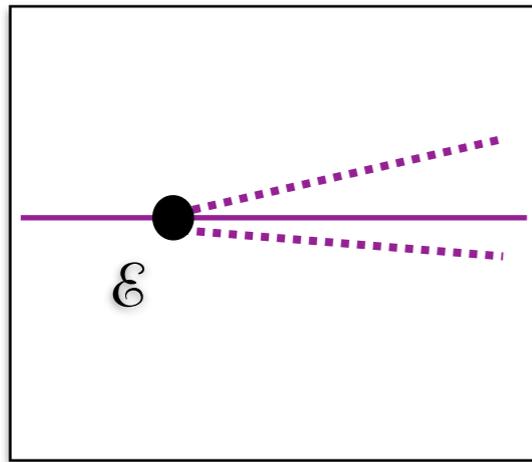


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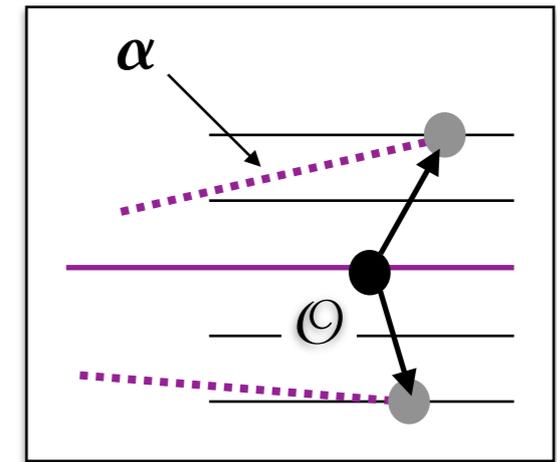
Π^A_B parallax matrix

$$D_{par} = D_{ang} |\det (\delta^A_B + m^A_B)|^{-1/2}$$

Stereoscopic parallax



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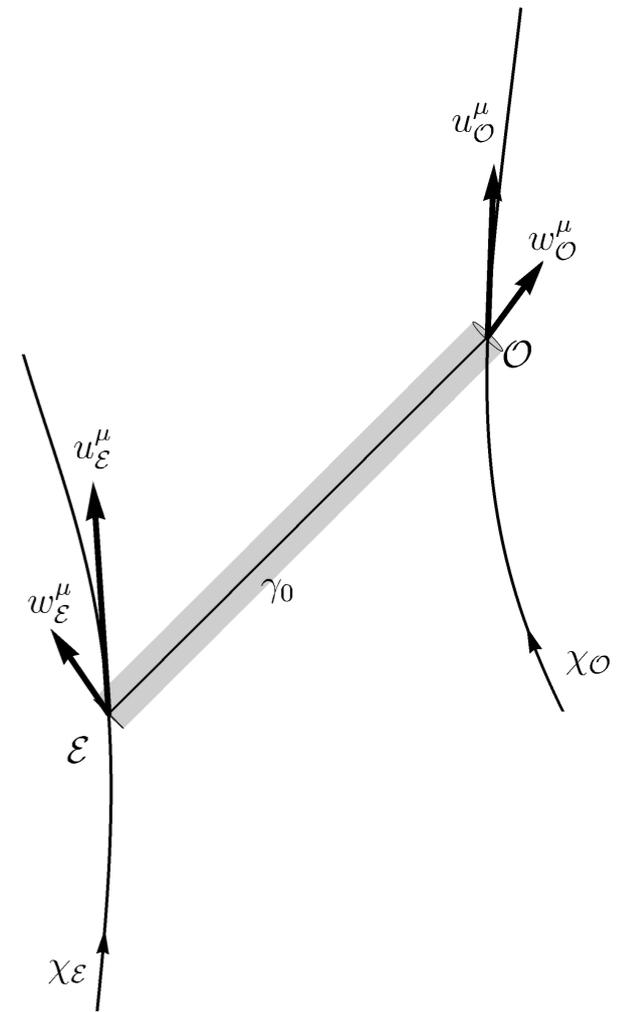
$$\delta \theta^A = -\frac{1}{l_{\sigma} u_{\mathcal{O}}^{\sigma}} \mathcal{D}^{-1A}_C (\delta^C_B + m^C_B) \delta x_{\mathcal{O}}^B$$

Π^A_B parallax matrix

$$D_{par} = D_{ang} \left| \det (\delta^A_B + m^A_B) \right|^{-1/2}$$

depend on $u_{\mathcal{O}}$, but NOT on $u_{\mathcal{E}}$!

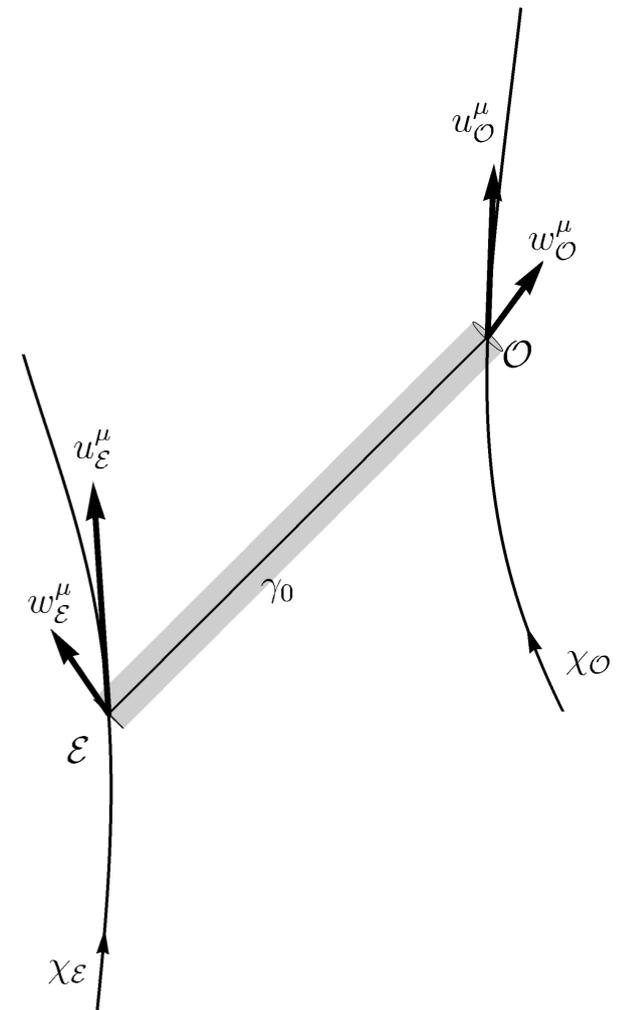
Position drift



Position drift

$$g([\delta x_{\mathcal{O}}], l) = g([\delta x_{\mathcal{E}}], l)$$

$$\mathcal{D}([\Delta l_{\mathcal{O}}]) = [\delta x_{\mathcal{E}} - \delta \hat{x}_{\mathcal{O}}] - m([\delta x_{\mathcal{O}}])$$

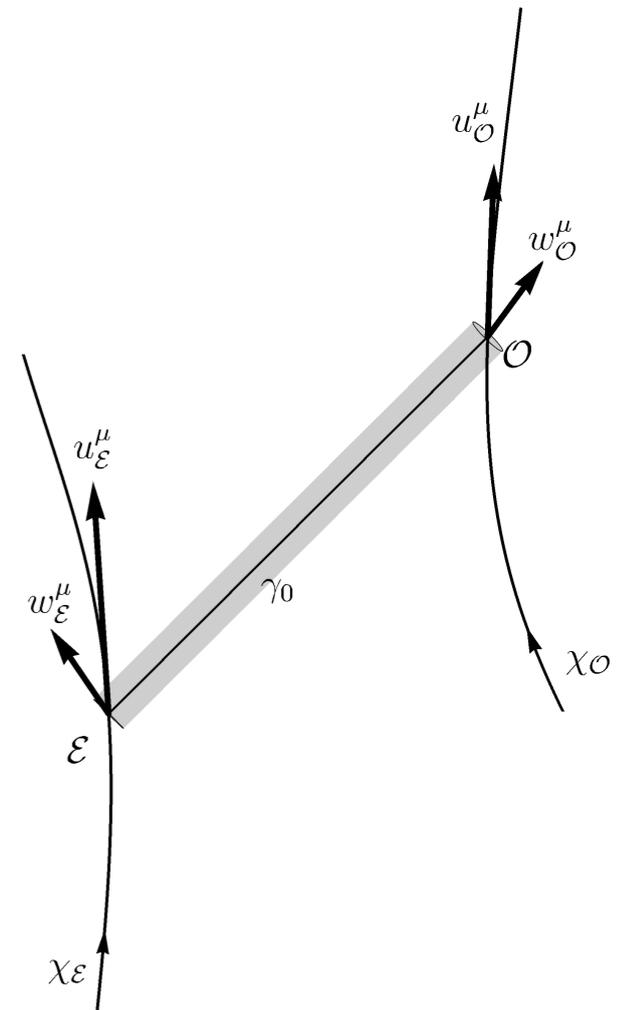


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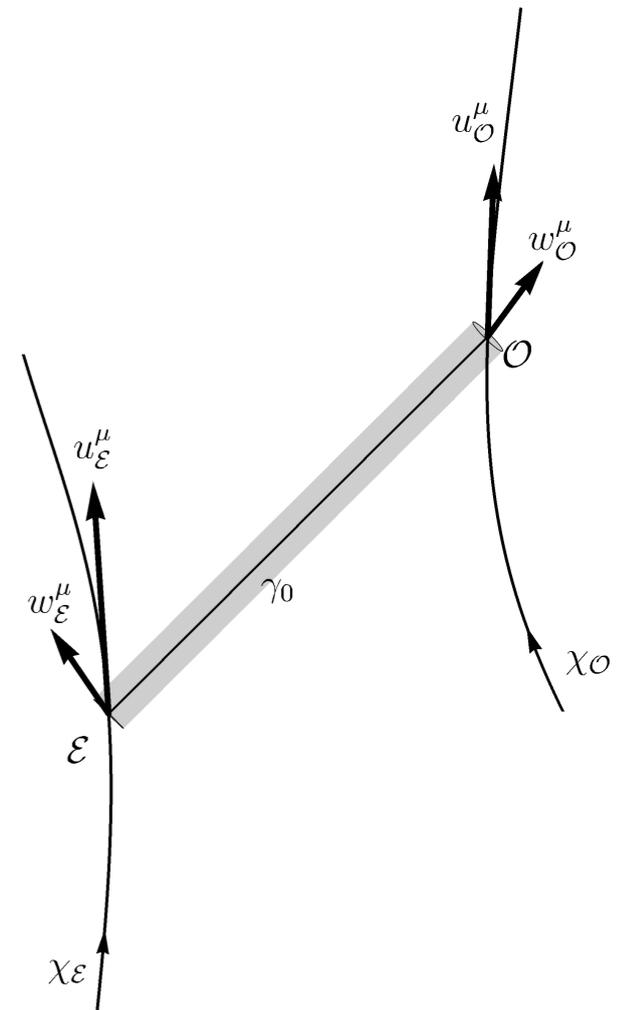
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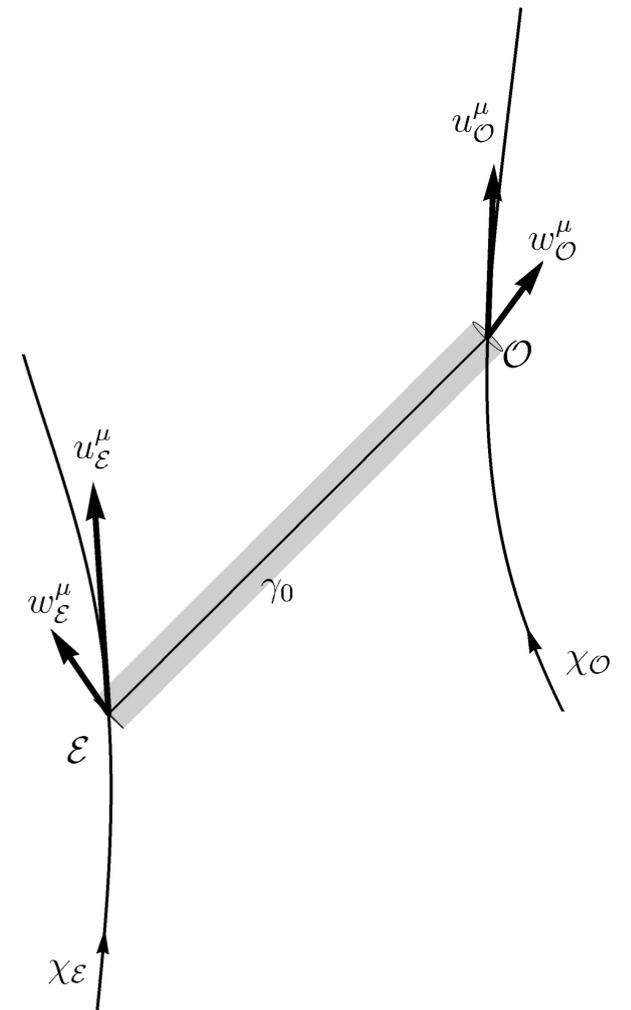
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- Non-geodesic observer: calculate the Fermi-Walker derivative of the position vector on the celestial sphere r^{μ}



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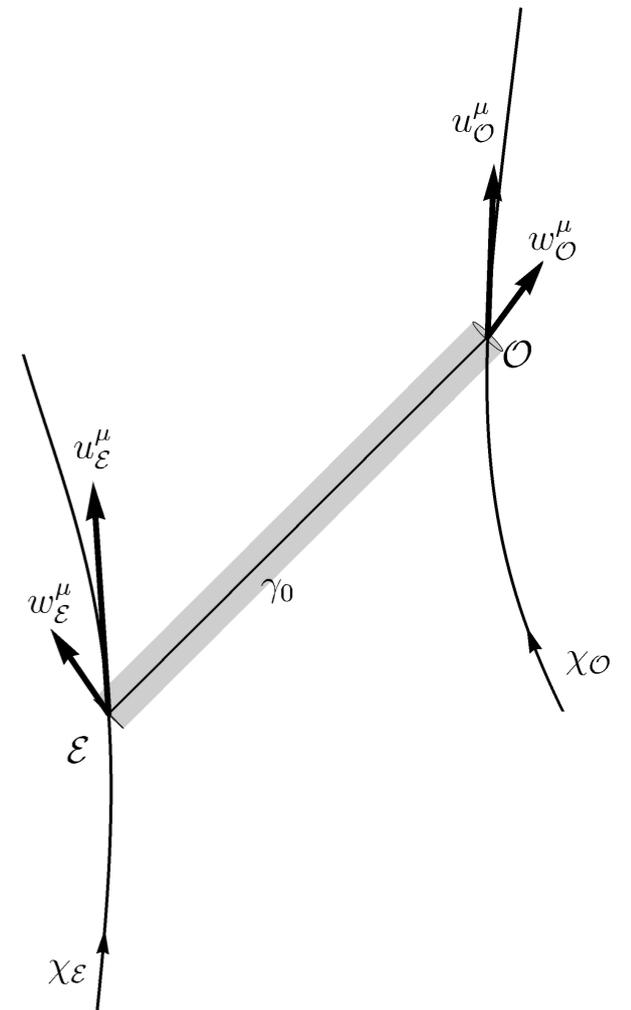
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$$\delta_{\mathcal{O}} r^A = w_{\mathcal{O}}^A + \frac{1}{l_{\sigma} u_{\mathcal{O}}^{\sigma}} \mathcal{D}^{-1A}{}_B \left(\left(\frac{1}{1+z} u_{\mathcal{E}} - \hat{u}_{\mathcal{O}} \right)^B - m^B{}_{\mu} u_{\mathcal{O}}^{\mu} \right)$$

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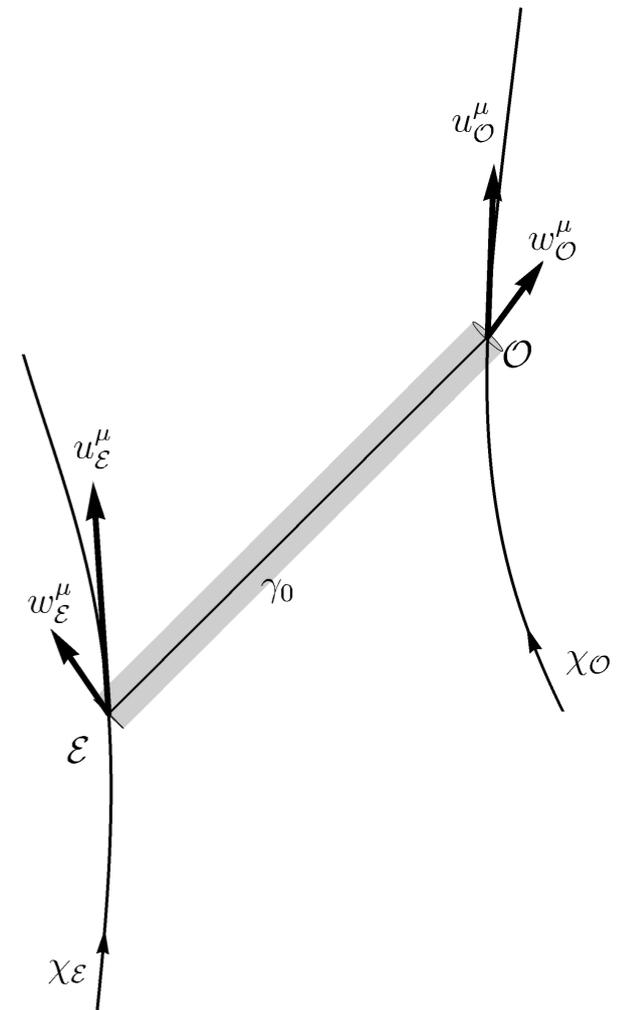
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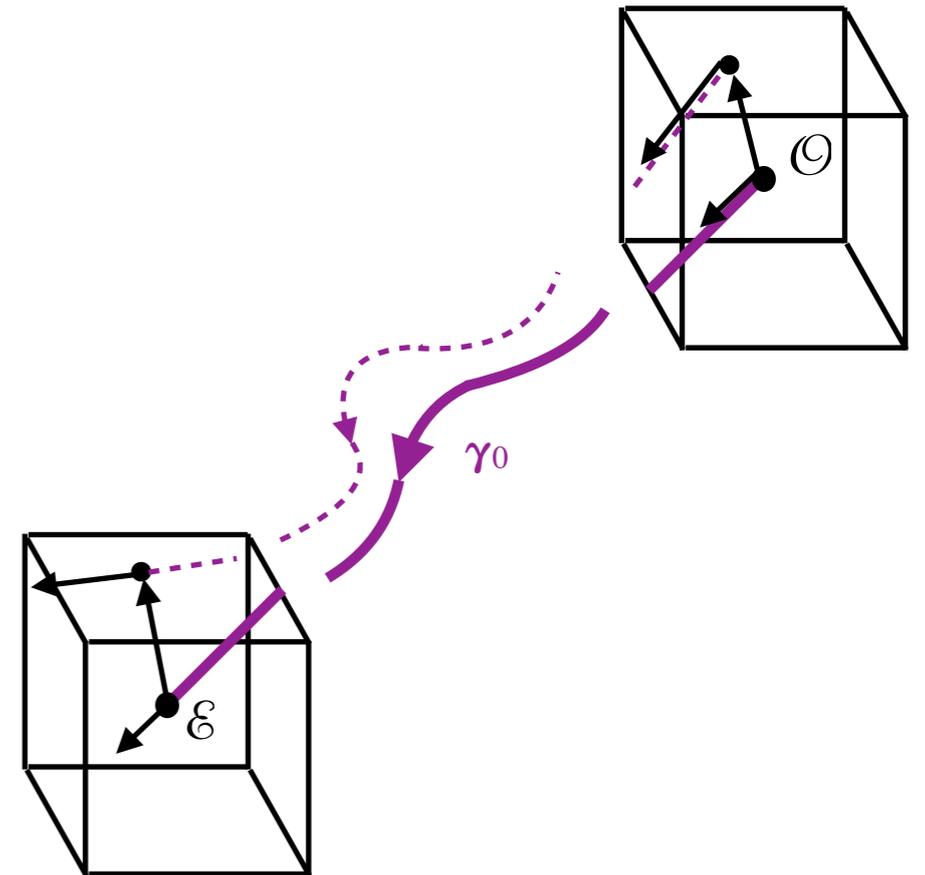
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aberration drift

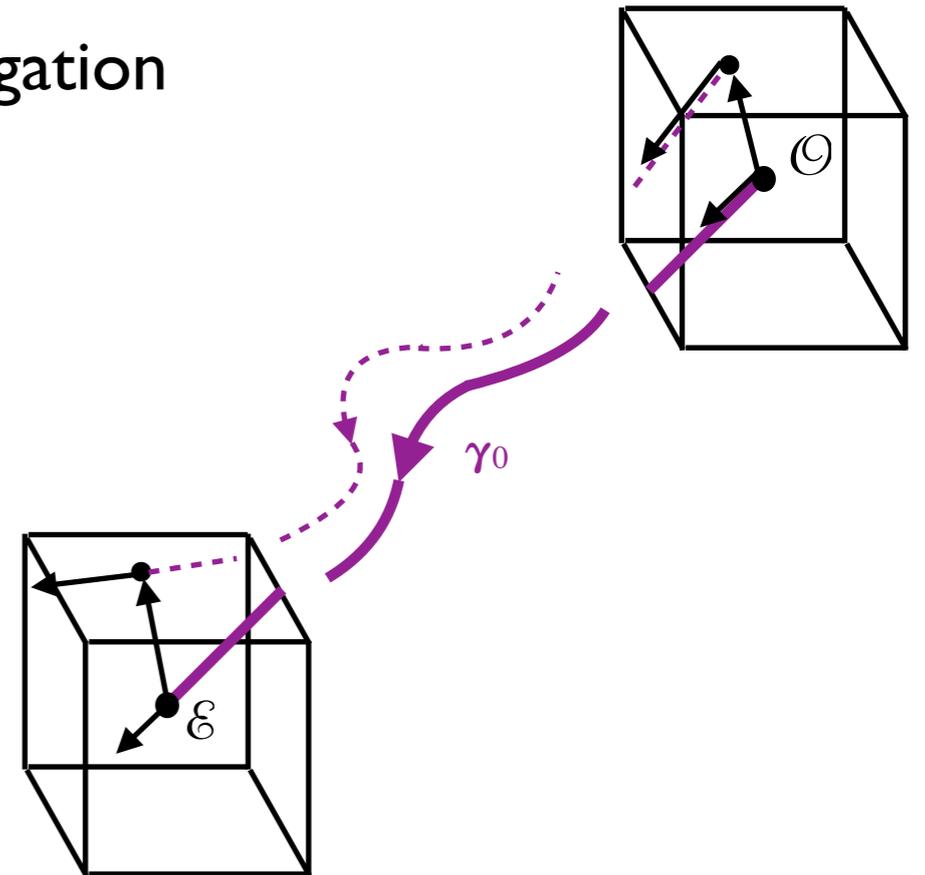
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Summary



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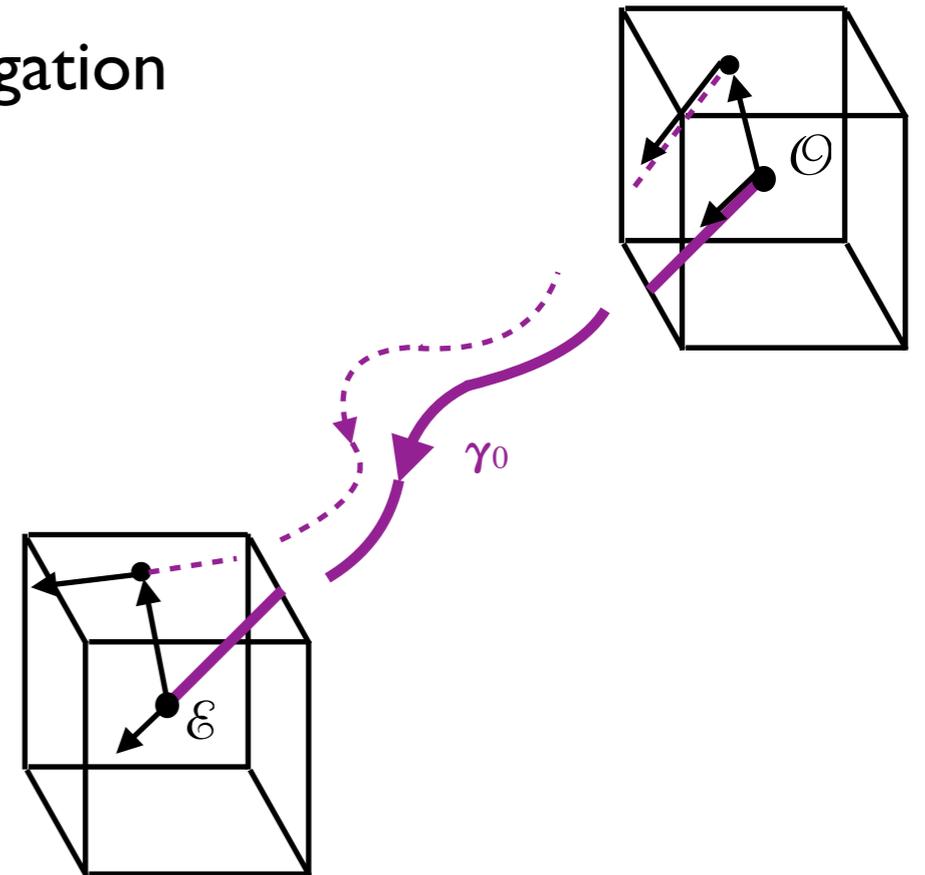
Covariant description of geometric optics and light propagation
between small, distant regions



Summary

Covariant description of geometric optics and light propagation
between small, distant regions

Valid in any spacetime (distant observer approximation)

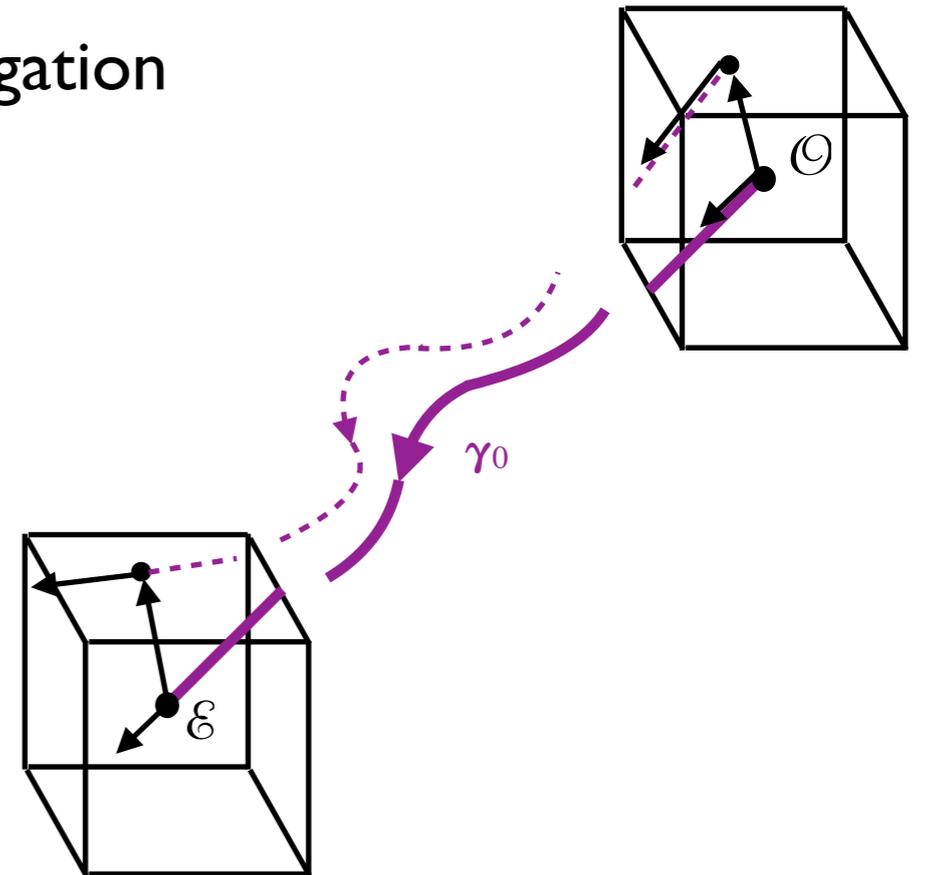


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Frame- (observer-) and coordinates-independent



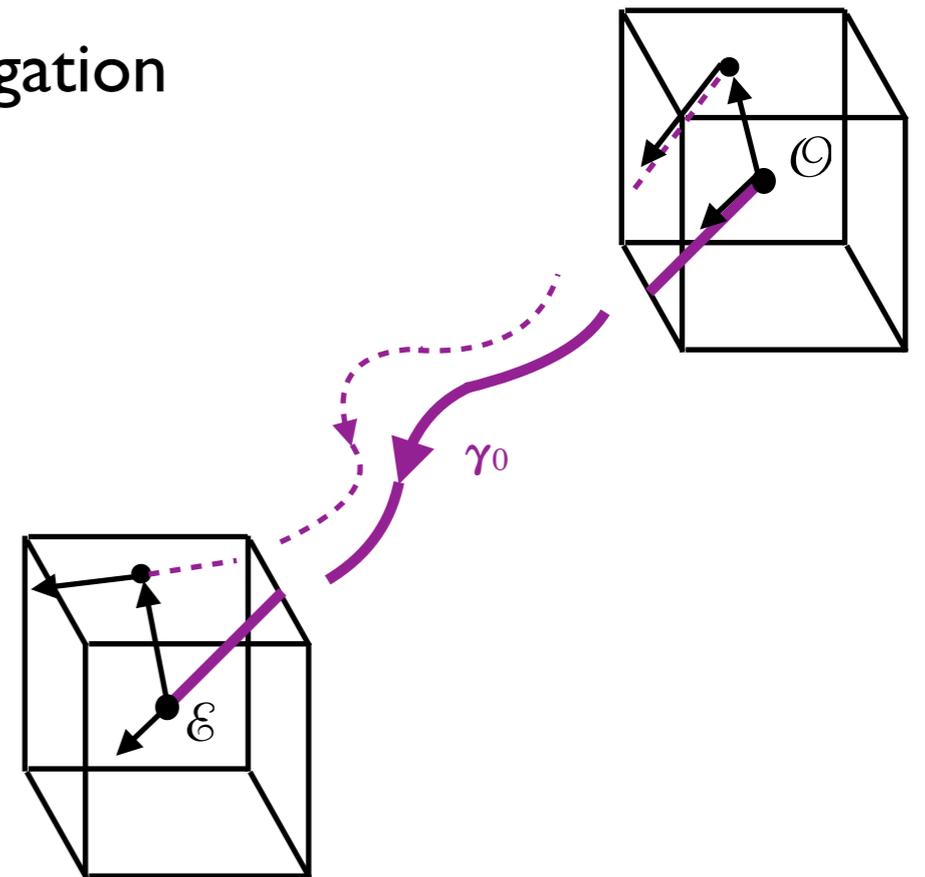
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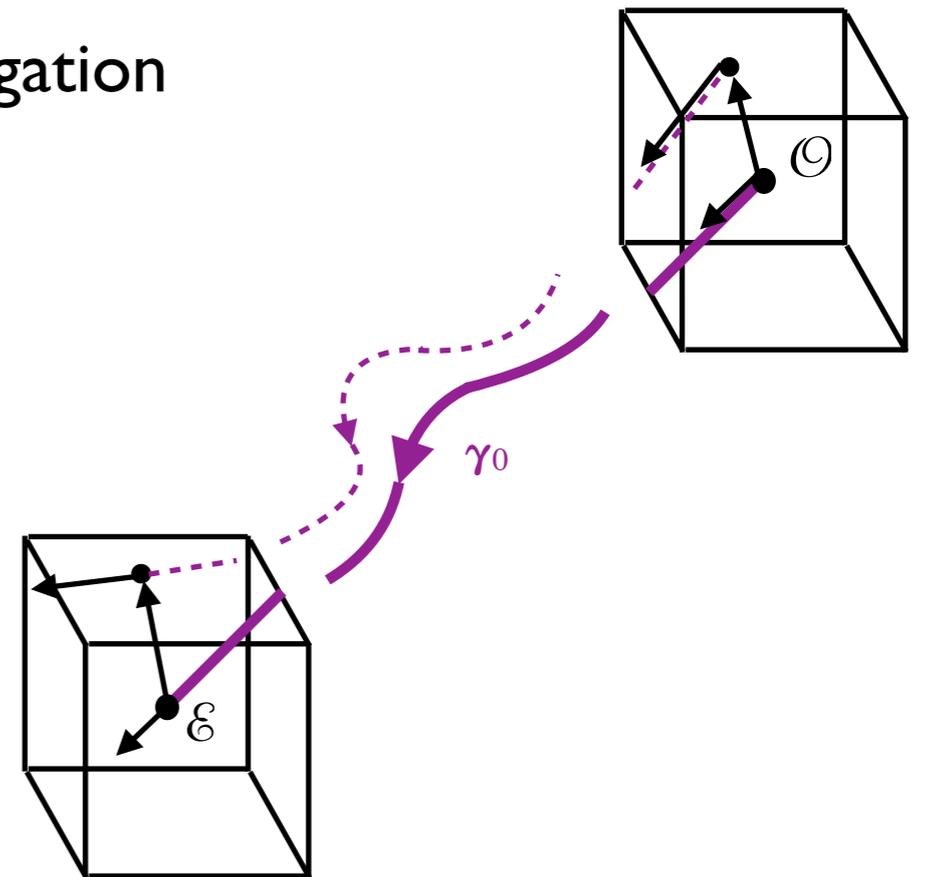
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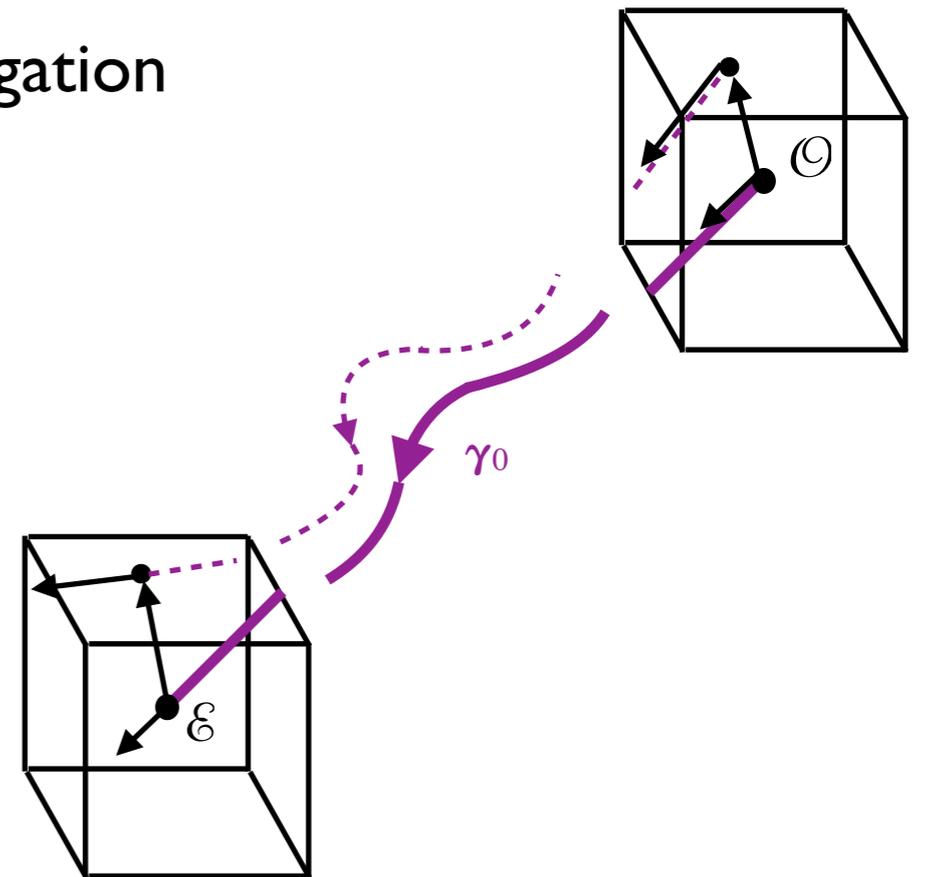
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Operators are linear functionals of the Riemann tensor along the line of sight

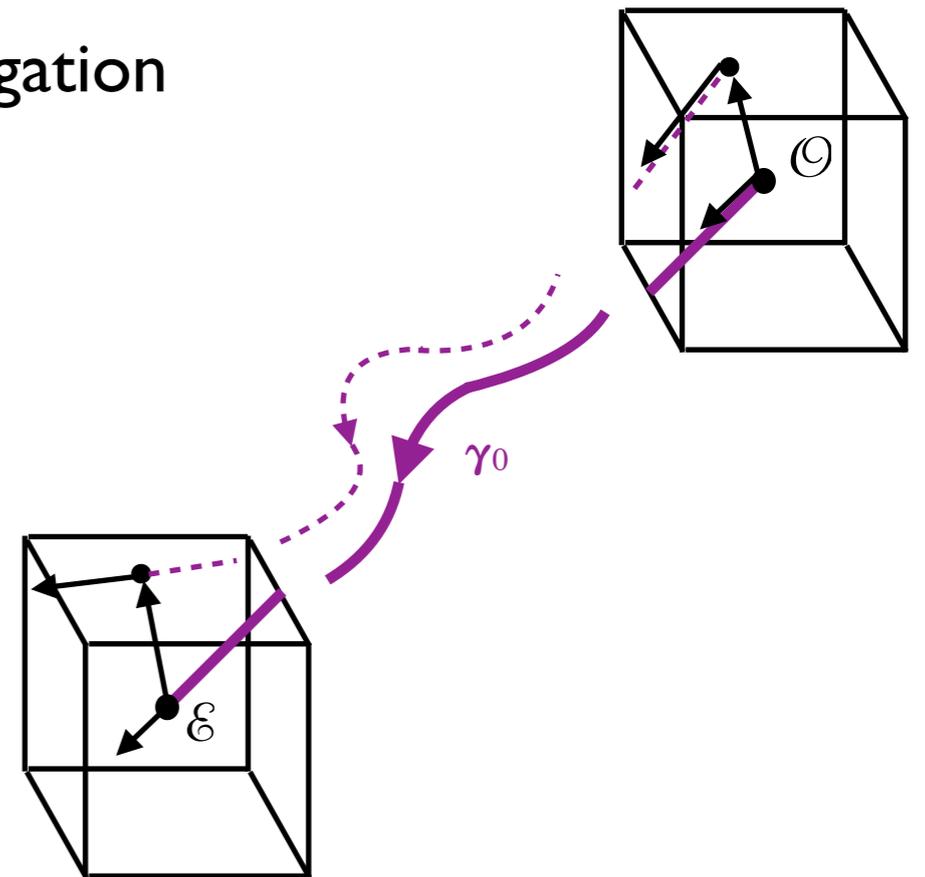
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Operators are linear functionals of the Riemann tensor along the line of sight

Image distortion, parallax, drift effects, redshift drift effects, Jacobi matrix drift, angular and luminosity distance drift

Applications + future research

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Cosmology: “real-time cosmology”, comparison of positions and redshifts of distant objects after ≈ 10 ys

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Relativistic astrophysics - corrections to parallax

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Light propagation through small-scale inhomogeneities in the Universe

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Thank you!